

Chemistry in Plasma Treated Liquids

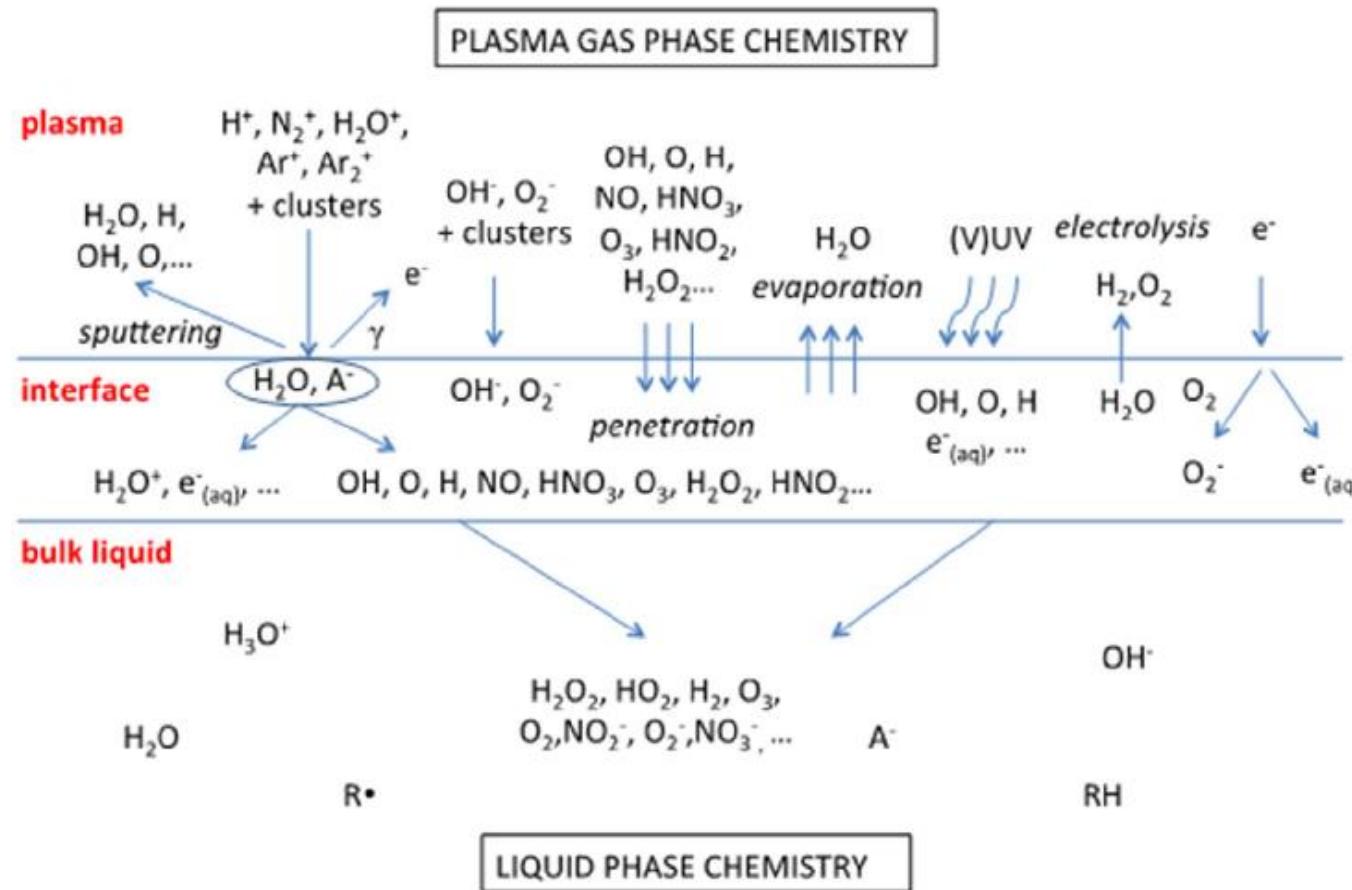
A. von Keudell

Experimental Physics II, Ruhr University Bochum



Transient Atmospheric Plasmas
from plasmas to liquids to solids

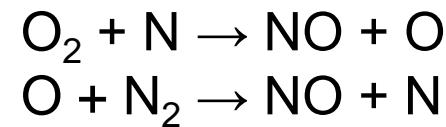
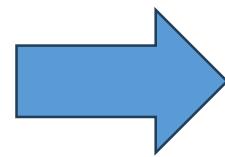
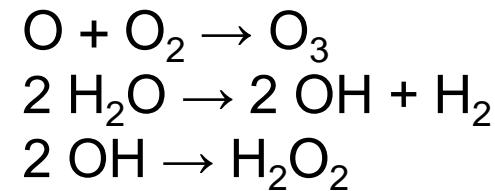
The plasma liquid interface



: P J Bruggeman et al 2016 *Plasma Sources Sci. Technol.* **25** 053002

Plasma in liquid chemistry in a nutshell

primary species in the plasma

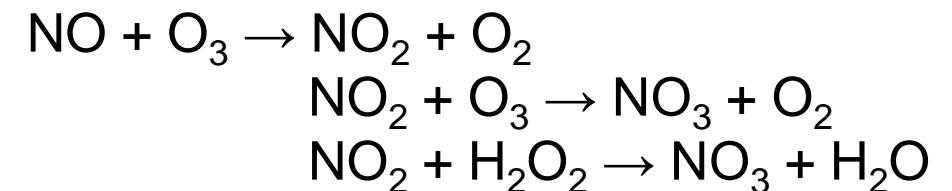
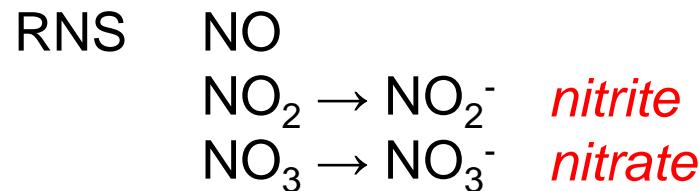
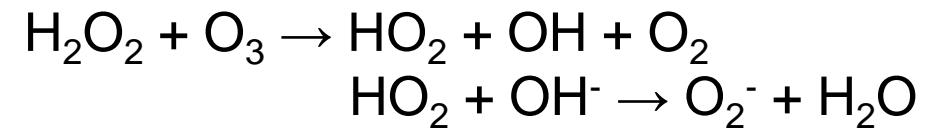
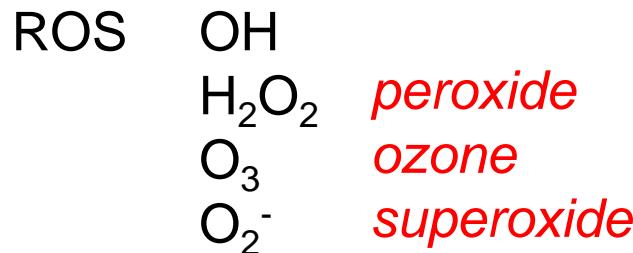


Zeldovich mechanism

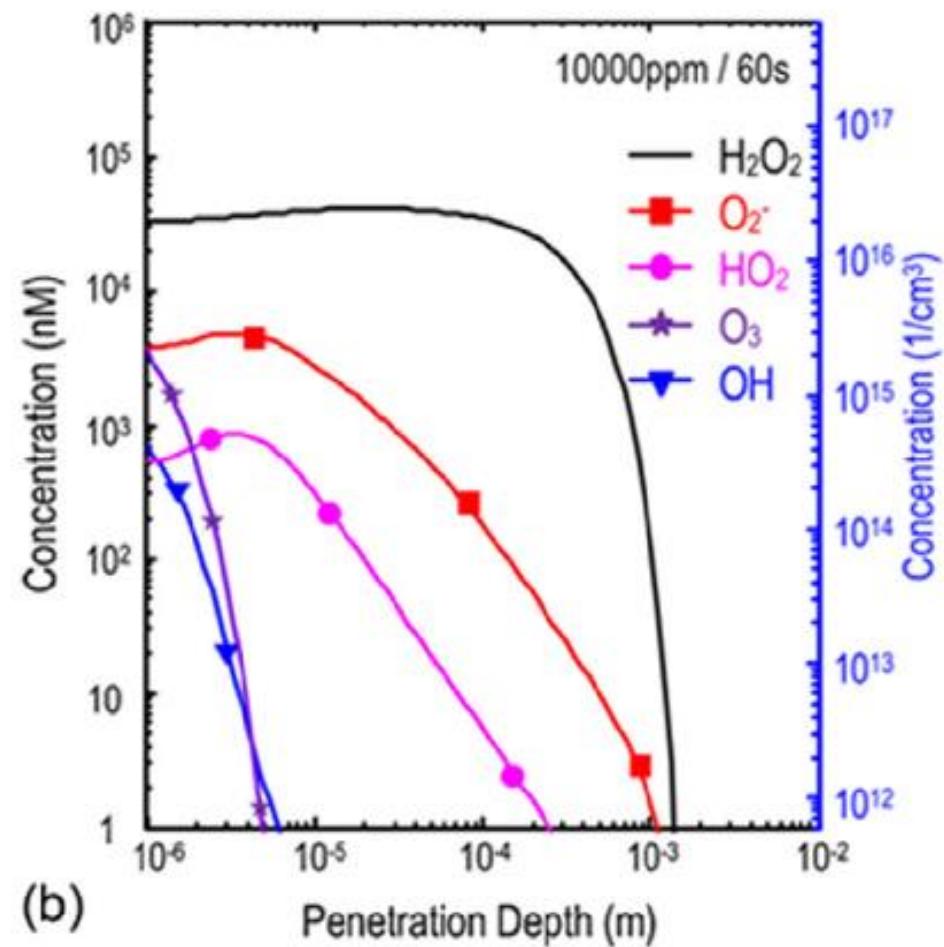
Henry
constants
determine
transfer

RONS

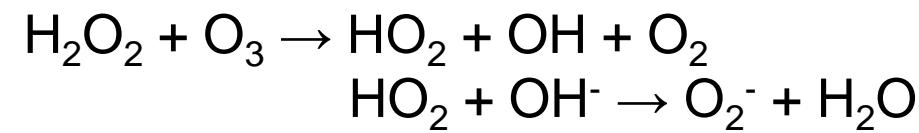
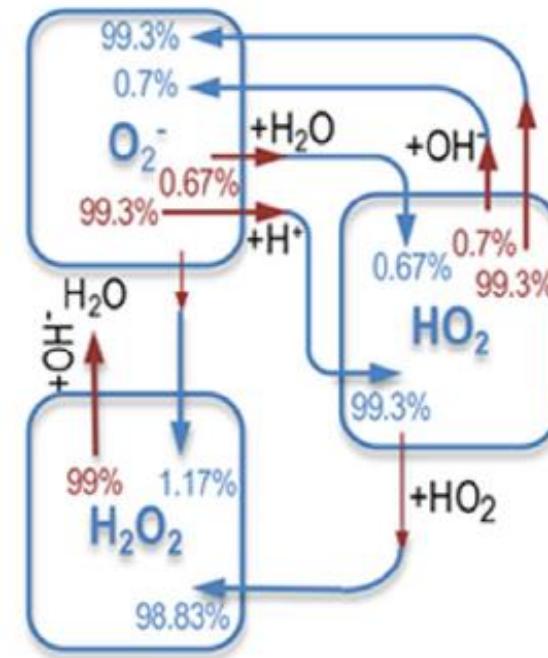
primary species in the liquid



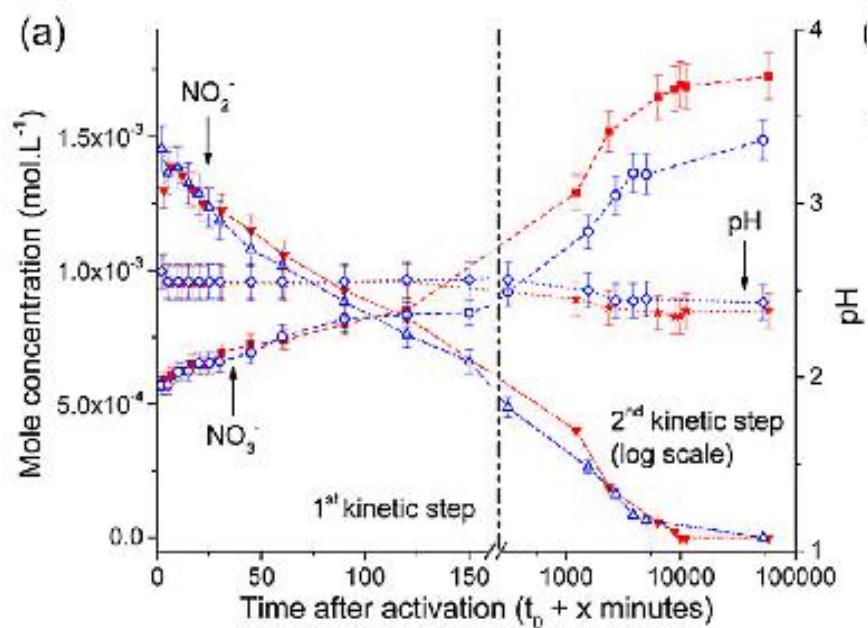
Plasma in liquid chemistry in a nutshell – chemistry modeling



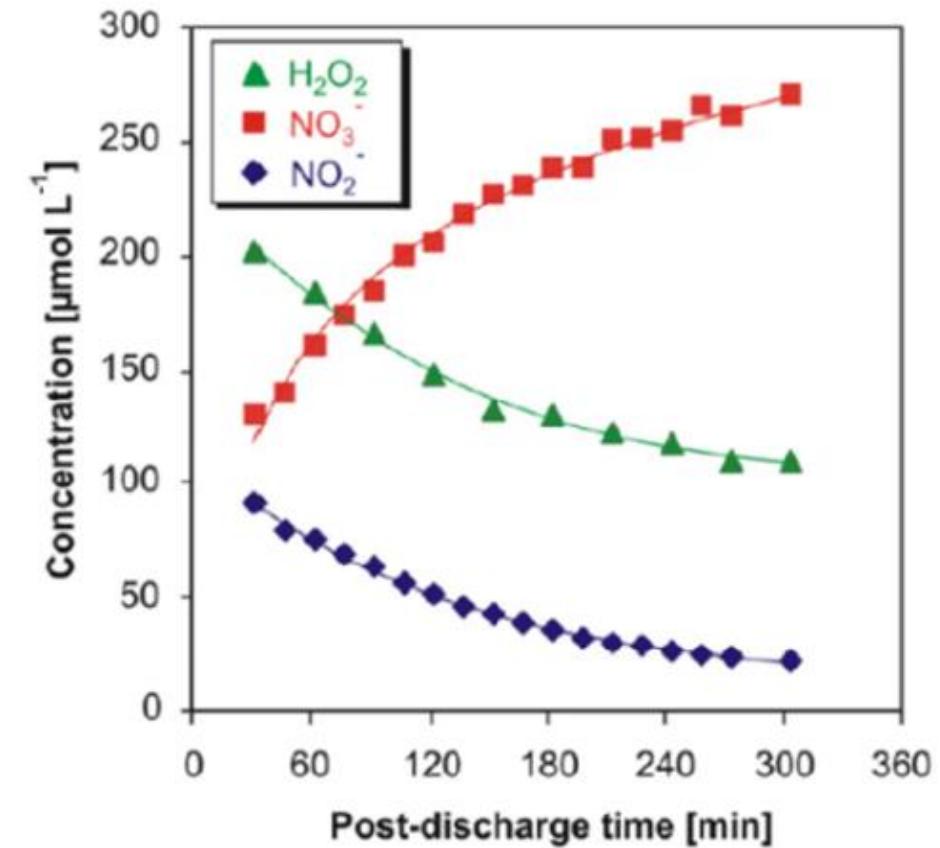
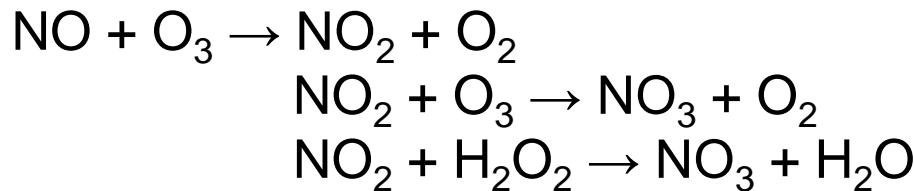
[188] Chen C, Liu D X, Liu Z C, Yang A J, Chen H L, Shama G and Kong M G 2014 A model of plasma-biofilm and plasma-tissue interactions at ambient pressure *Plasma Chem. Plasma Process.* **34** 403–41



Plasma activated water – PAW, change in reactivity



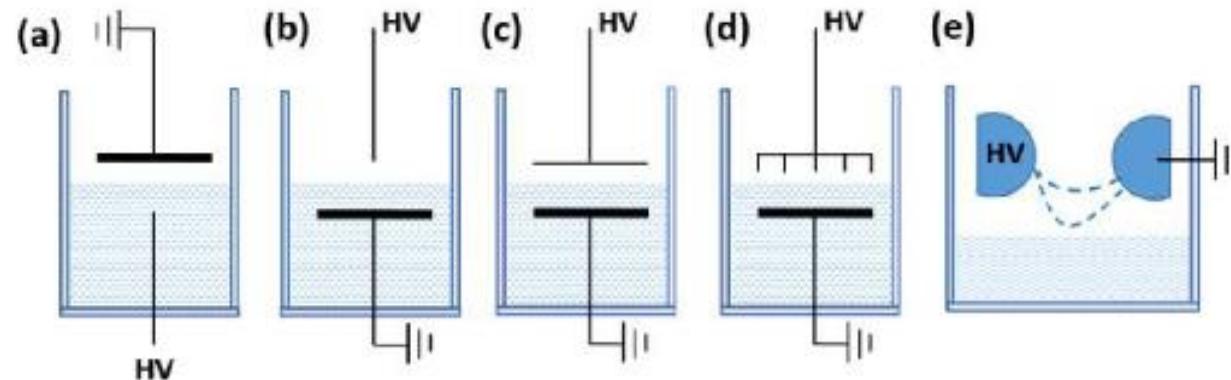
Wartel, M., Faubert, F., Dirlau, I. D., Rudz, S., Pellerin, N., Astanei, D., Burlica, R., Hnatiuc, B., & Pellerin, S. (2021). Analysis of plasma activated water by gliding arc at atmospheric pressure: Effect of the chemical composition of water on the activation. *Journal of Applied Physics*, 129(23), 233301. <https://doi.org/10.1063/5.0040035>



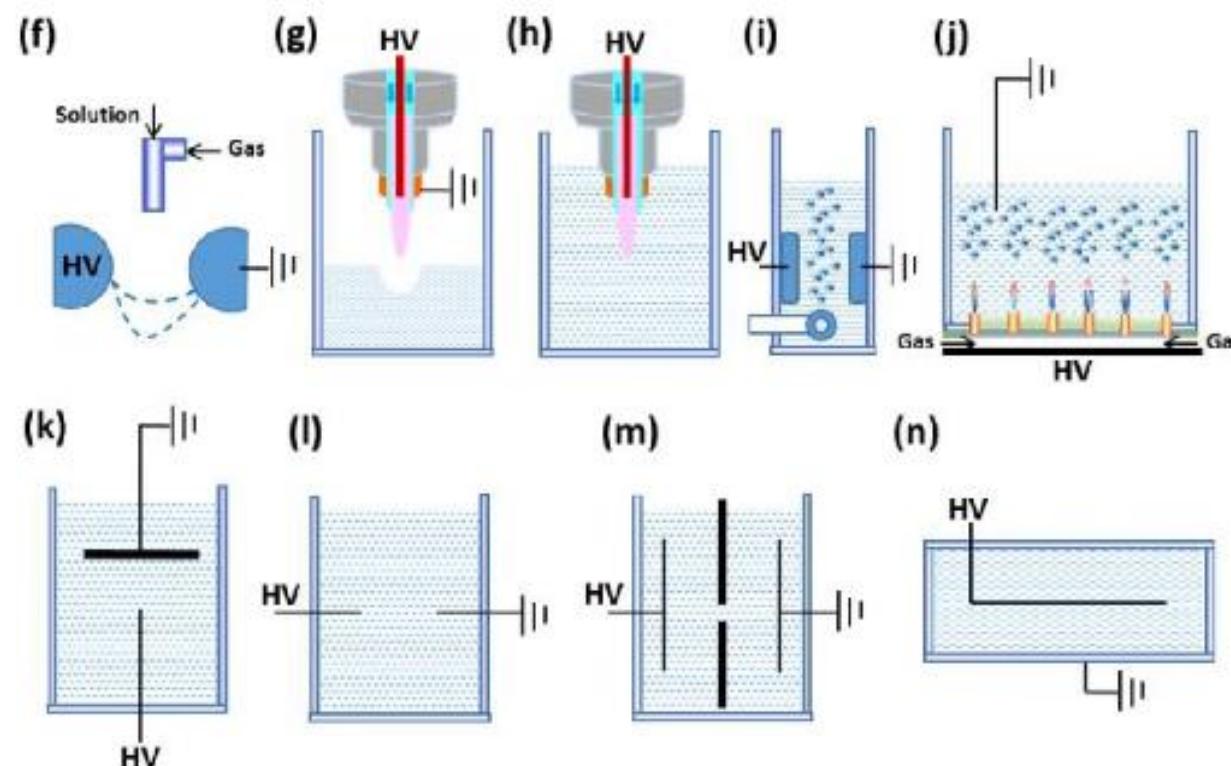
[488] Lukes P, Dolezalova E, Sisrova I and Clupek M 2014
Aqueous-phase chemistry and bactericidal effects from an air discharge plasma in contact with water: evidence for the formation of peroxy nitrite through a pseudo-second-order post-discharge reaction of H_2O_2 and HNO_2 *Plasma Source Sci. Technol.* **23** 015019

Possible plasma arrangements of plasma in liquids

Plasma above the liquid
(liquid electrode)
→ RONS

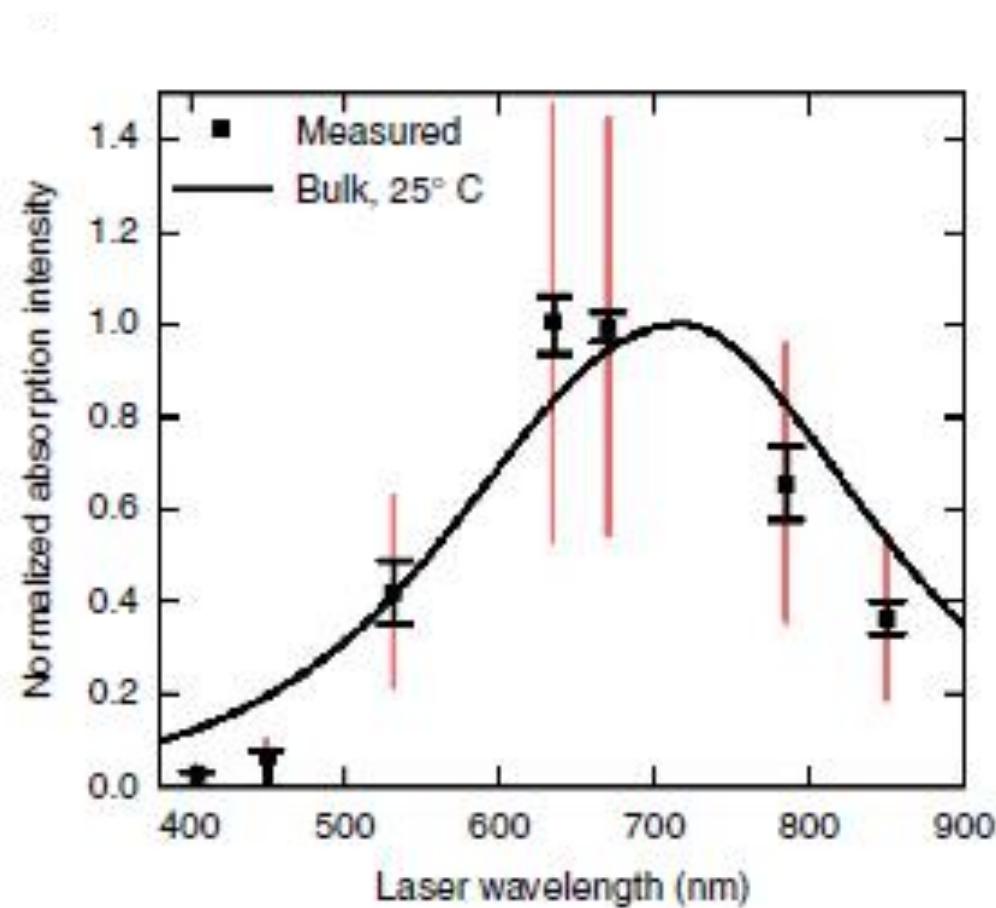
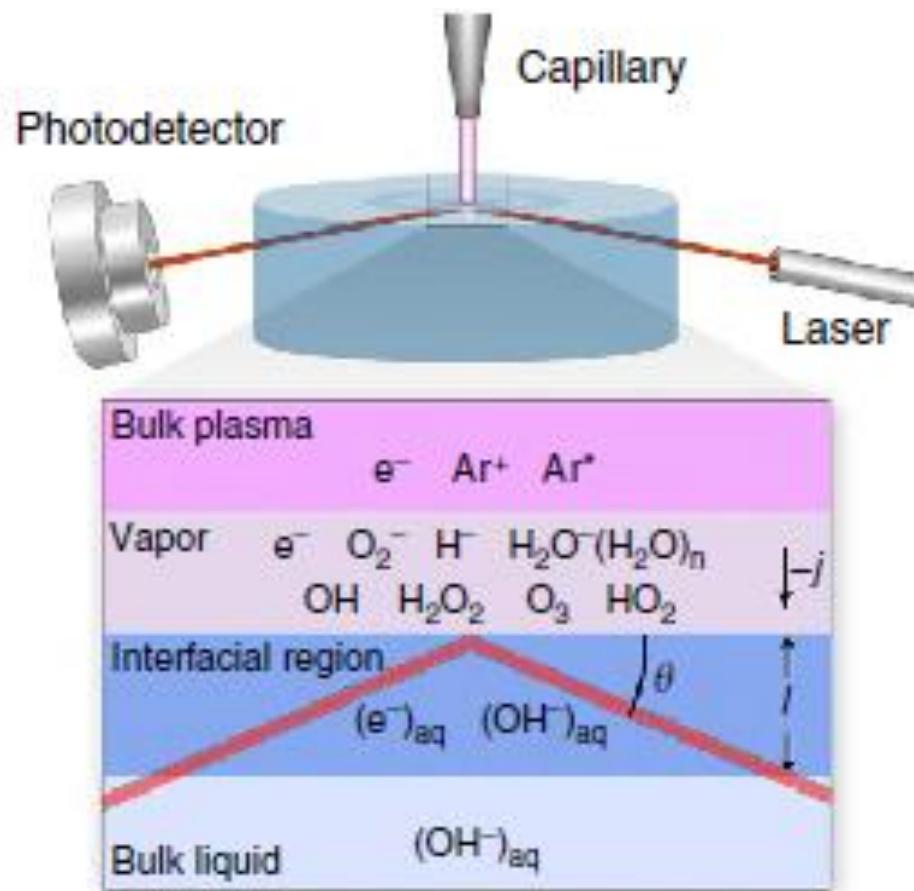


Plasma inside the liquid
→ ROS



: Renwu Zhou et al 2020 J. Phys. D: Appl. Phys. 53 303001

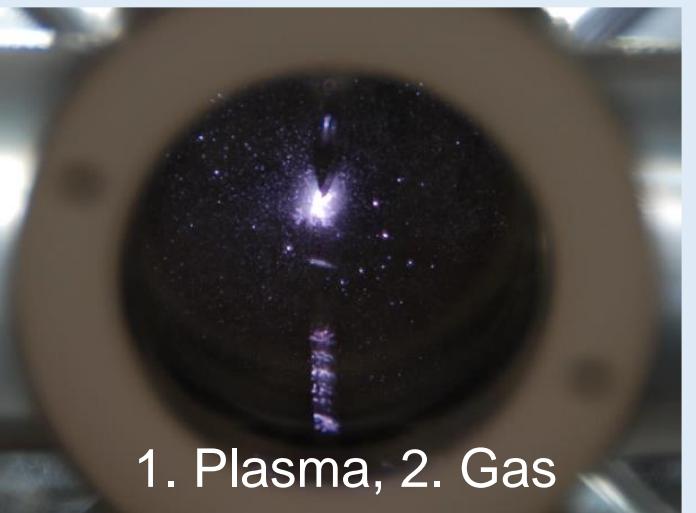
Solvated electrons in liquids induced by plasmas



Nature Communications (2015)
P. Rumbach, D. Bartels, M. Sankaran, D. Go,
doi 6:7248 | DOI: 10.1038/ncomms8248

Plasma Chemistry in liquids – example H₂O₂ generation

Nanosecond plasmas

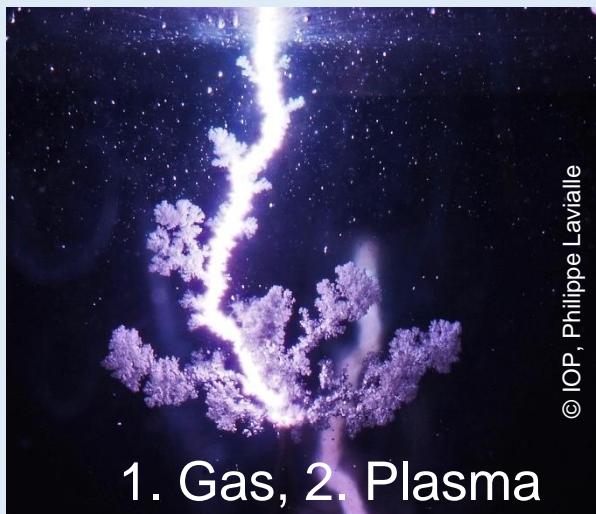


1. Plasma, 2. Gas

- Short nanosecond pulses
- Pulses too short for formation of a gas bubble
- Plasma ignition inside the liquid

Extremely high pressures and temperatures
10000 bar, 50000 K

Microsecond plasmas

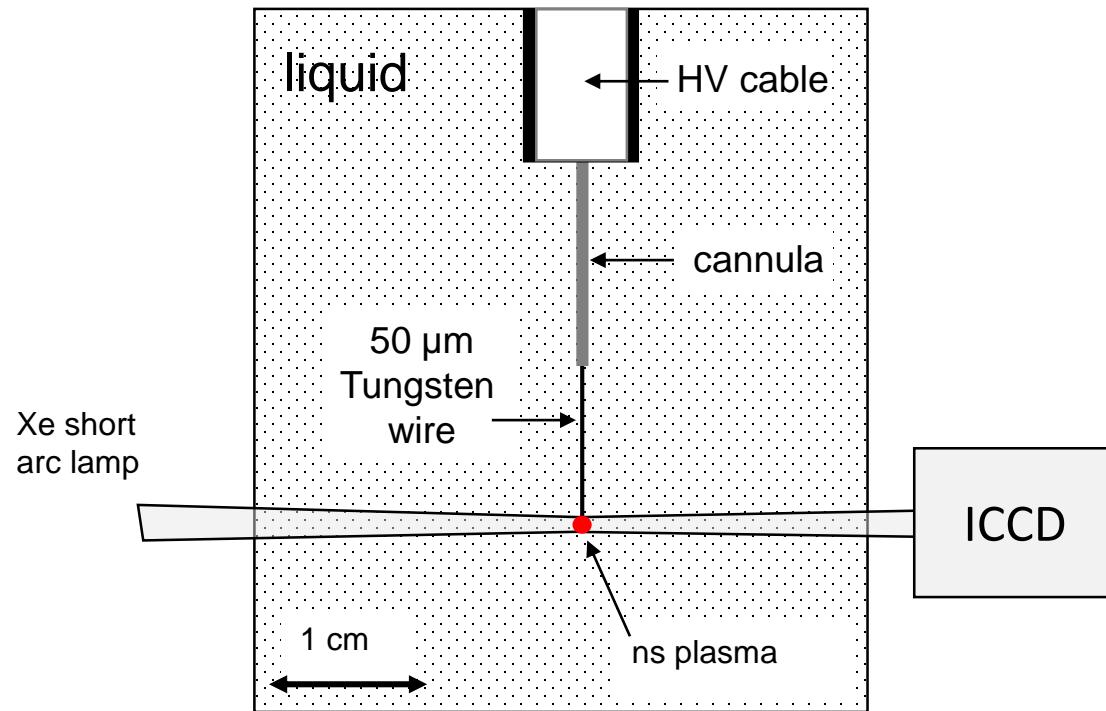
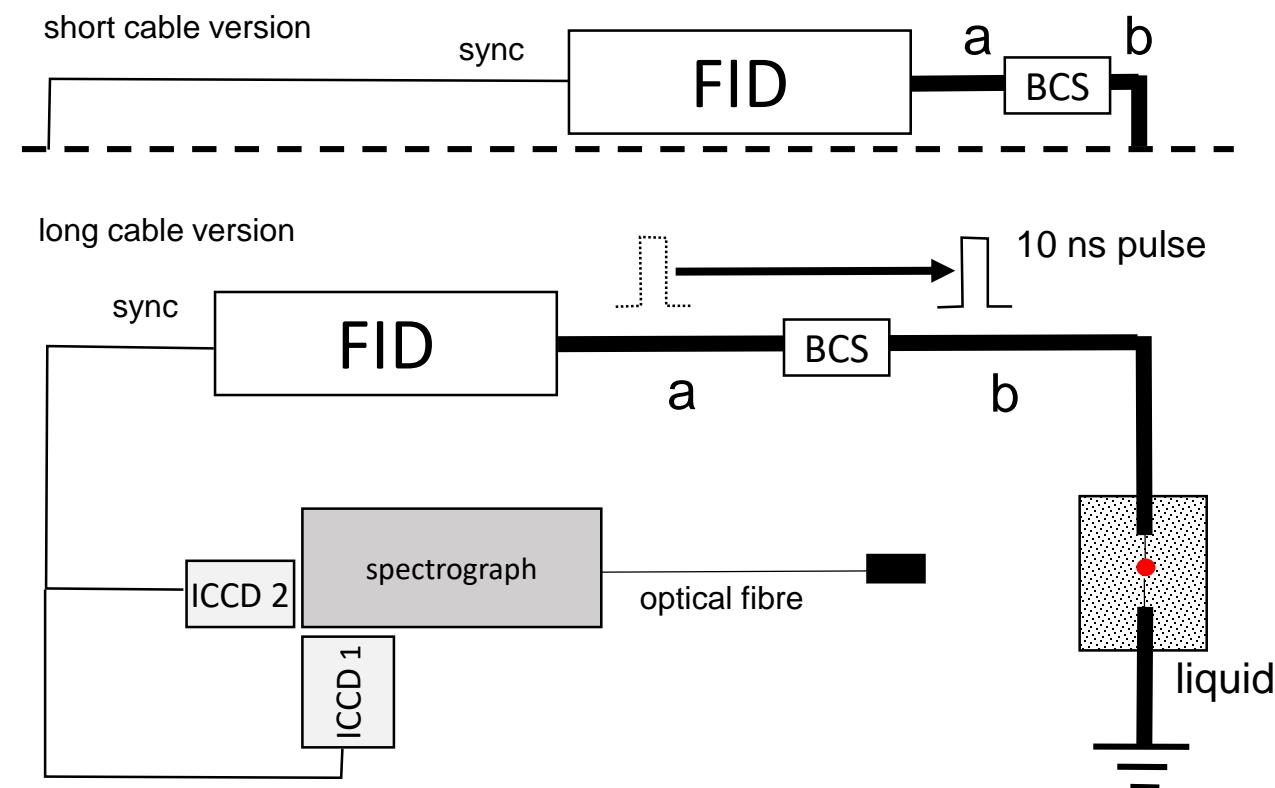


1. Gas, 2. Plasma

- Long microsecond pulses
- Initial formation of a gas bubble due to Ohmic heating or E-field rupture
- Plasma ignition inside the water vapor in the bubble

Moderate high pressures and temperatures
10 bar, 20000 K

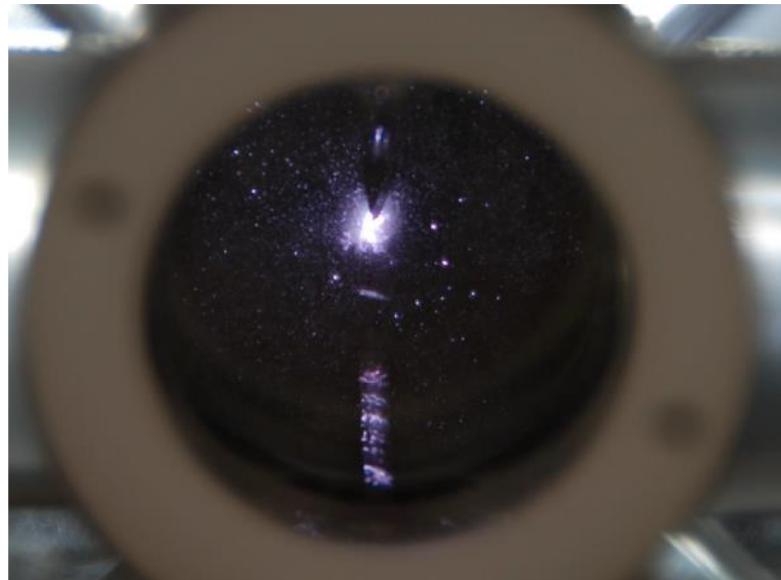
Plasma Chemistry in liquids – example H₂O₂ generation



- 10 ns pulses, 14...30 kV, distilled water
- Shadowgraphy 2 ns ... 70 ns gate widths
- Optical emission spectroscopy with 2ns and 30 ns gate time

Goals:

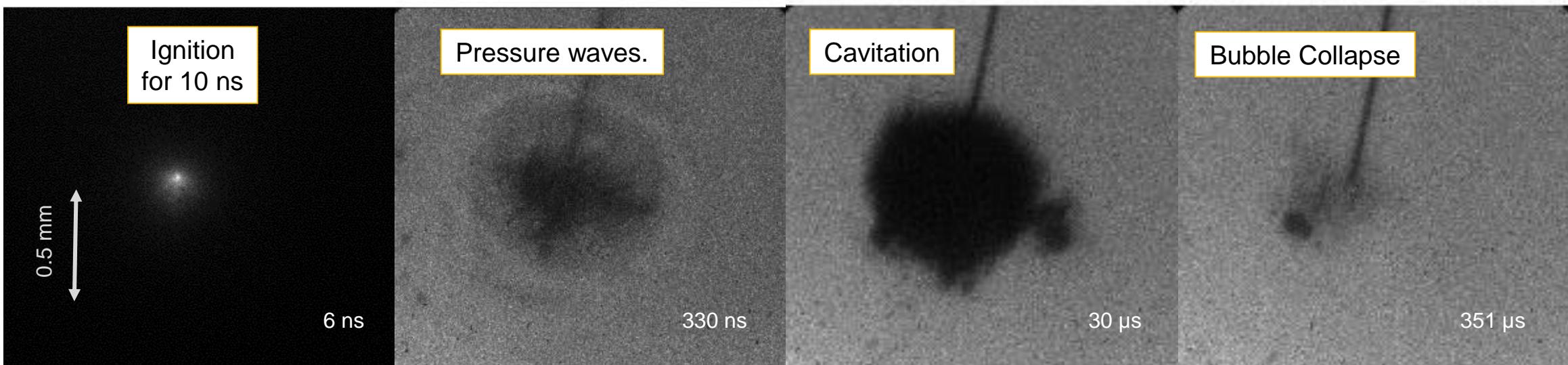
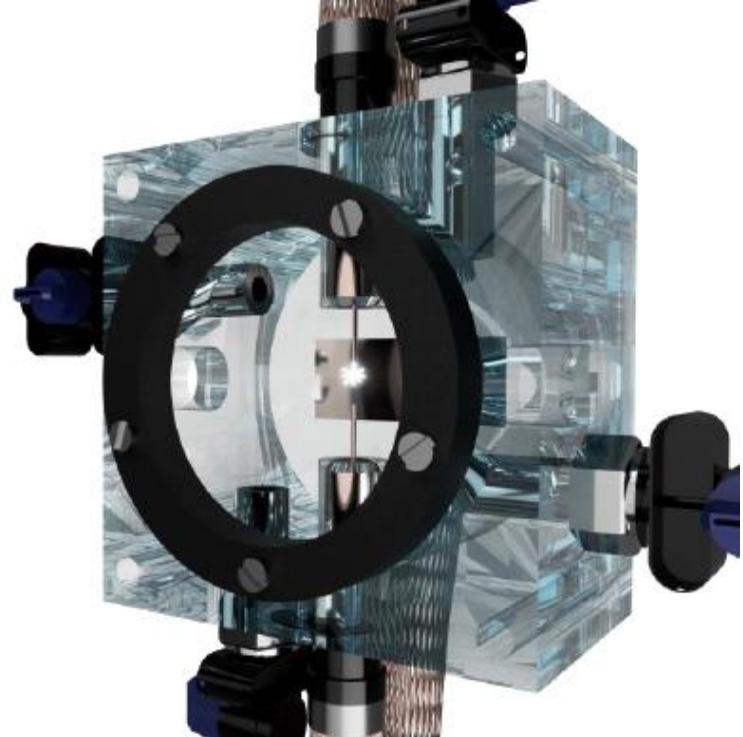
- reactive liquid production
- regeneration of catalysts for plasma enhanced electrolysis



Temporal evolution of the discharge

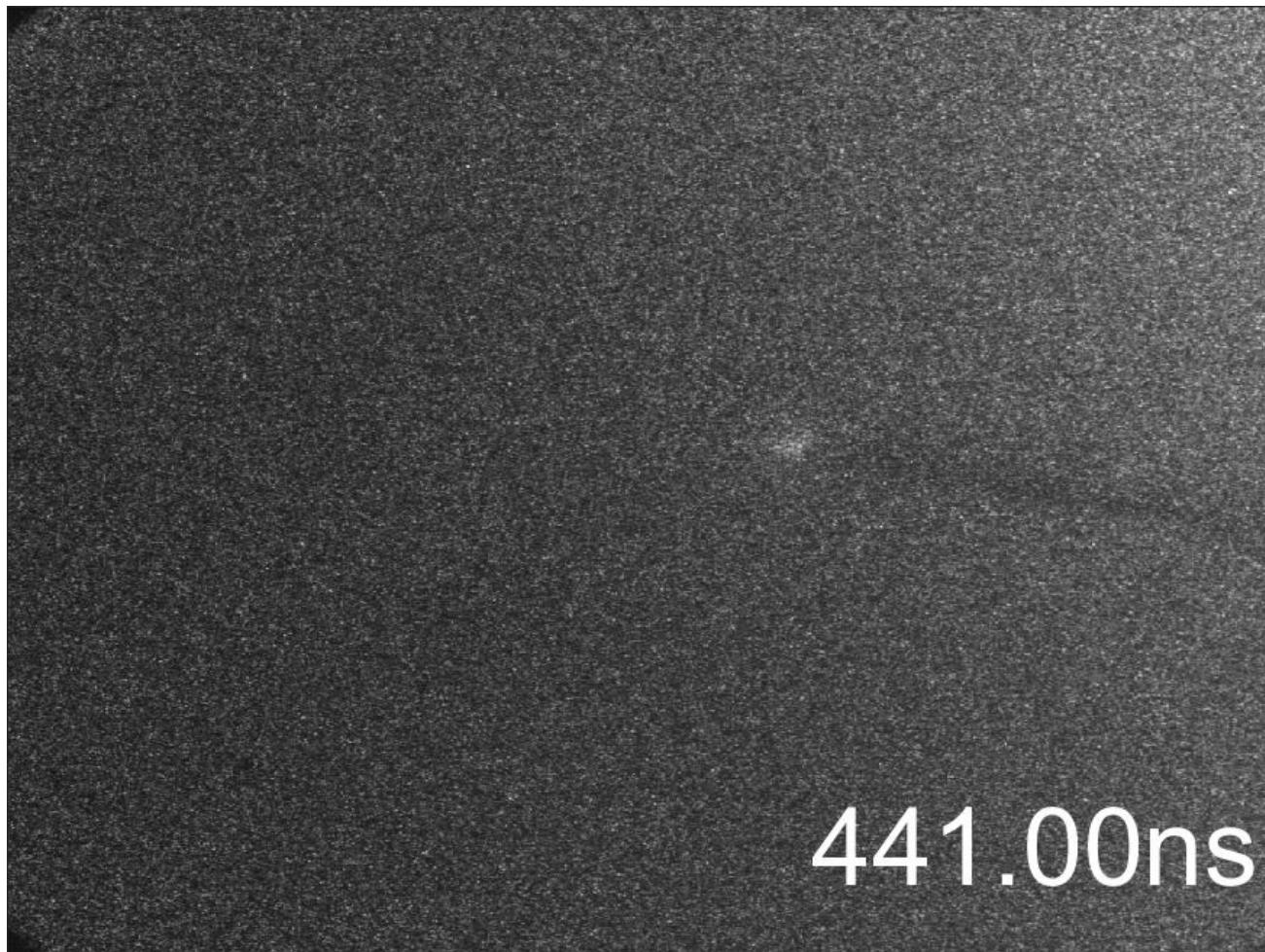
4 phases

- Ignition (Streamer)
- Pressure waves propagate
- Conversion of water into plasma and vapor
- Expansion of a cavitation bubble

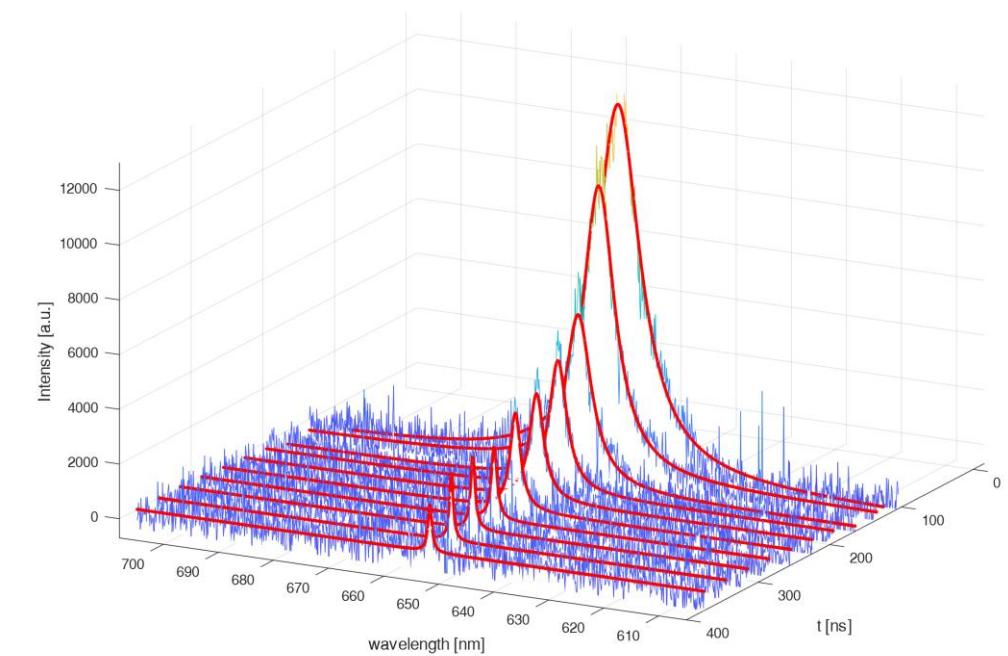


Plasma Chemistry in liquids – example H₂O₂ generation

Temporal evolution of the discharge

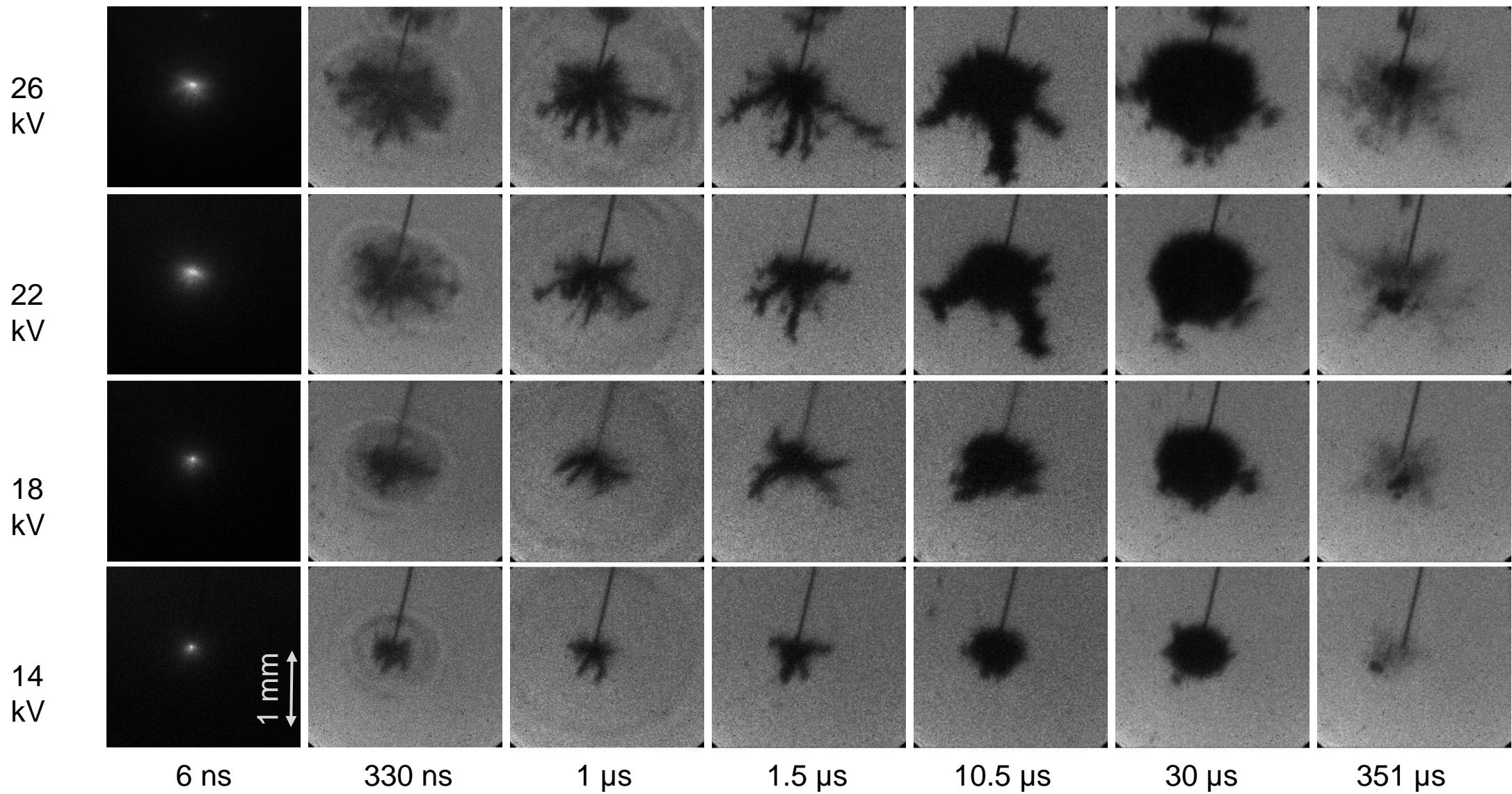


H alpha emission after the plasma pulse



60 ns gate, 30 ns steps

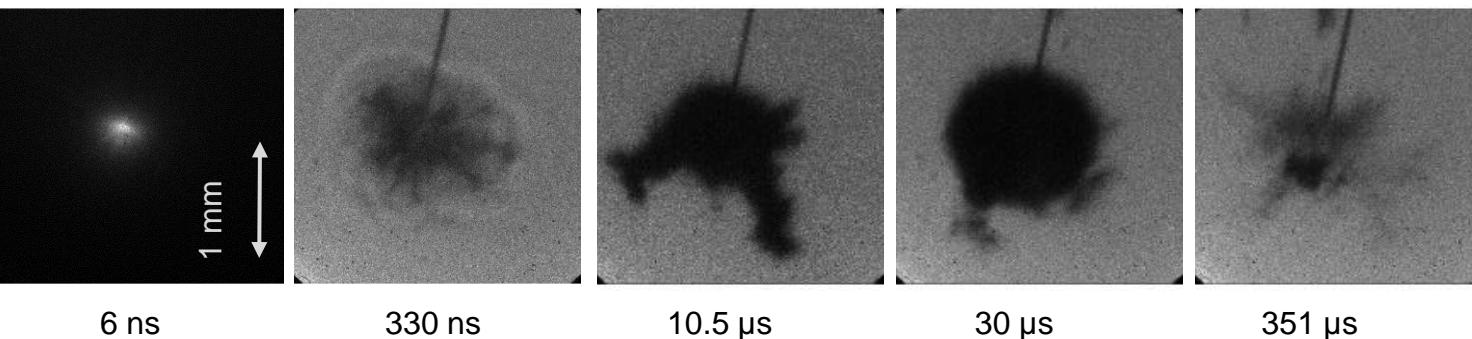
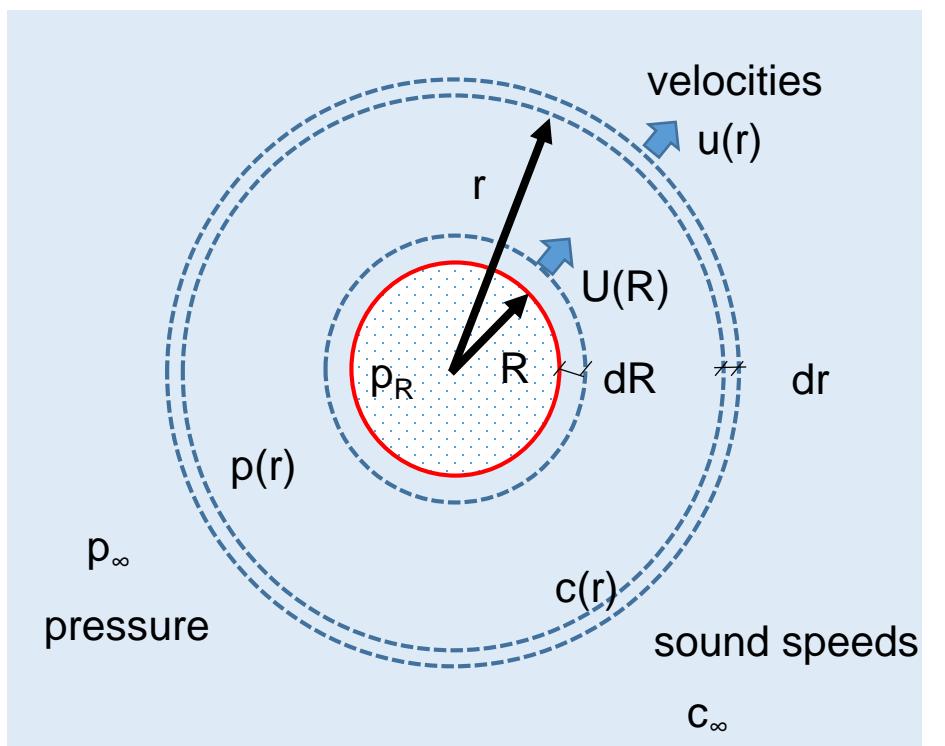
Monitoring Bubble Expansion with Shadowgraphy – monitoring sound waves and cavitation



Comparison of Bubble Expansion with cavitation theory

Rayleigh Plesset equation:

$$R \ddot{R} \left(1 - \frac{\dot{R}}{c}\right) + \frac{3}{2} \dot{R}^2 \left(1 - \frac{\dot{R}}{3c}\right) = h \left(1 + \frac{\dot{R}}{c}\right) + \left(1 - \frac{\dot{R}}{c}\right) \frac{R}{c} \frac{\partial h}{\partial t}$$



Solution for R(t) assuming a compressible liquid and condensation of vapor species

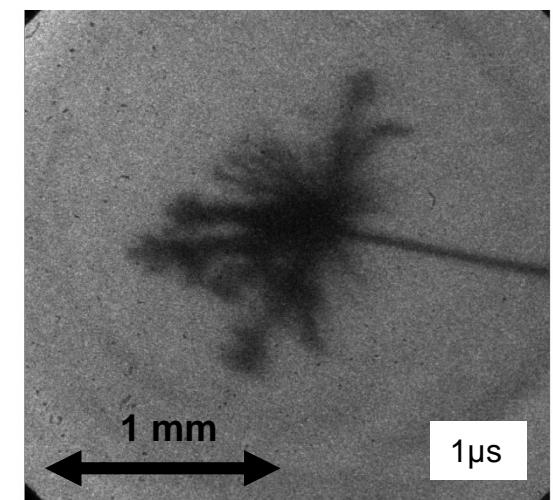
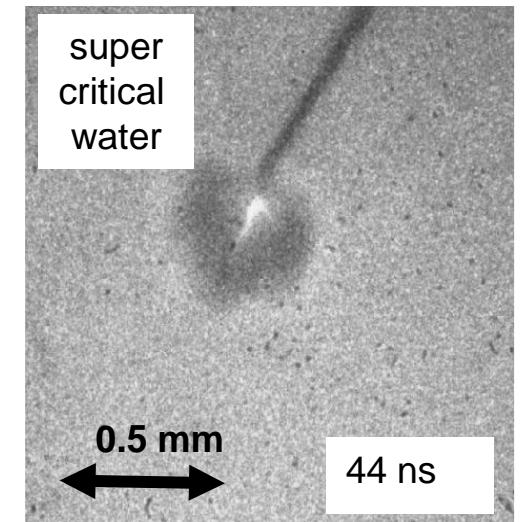
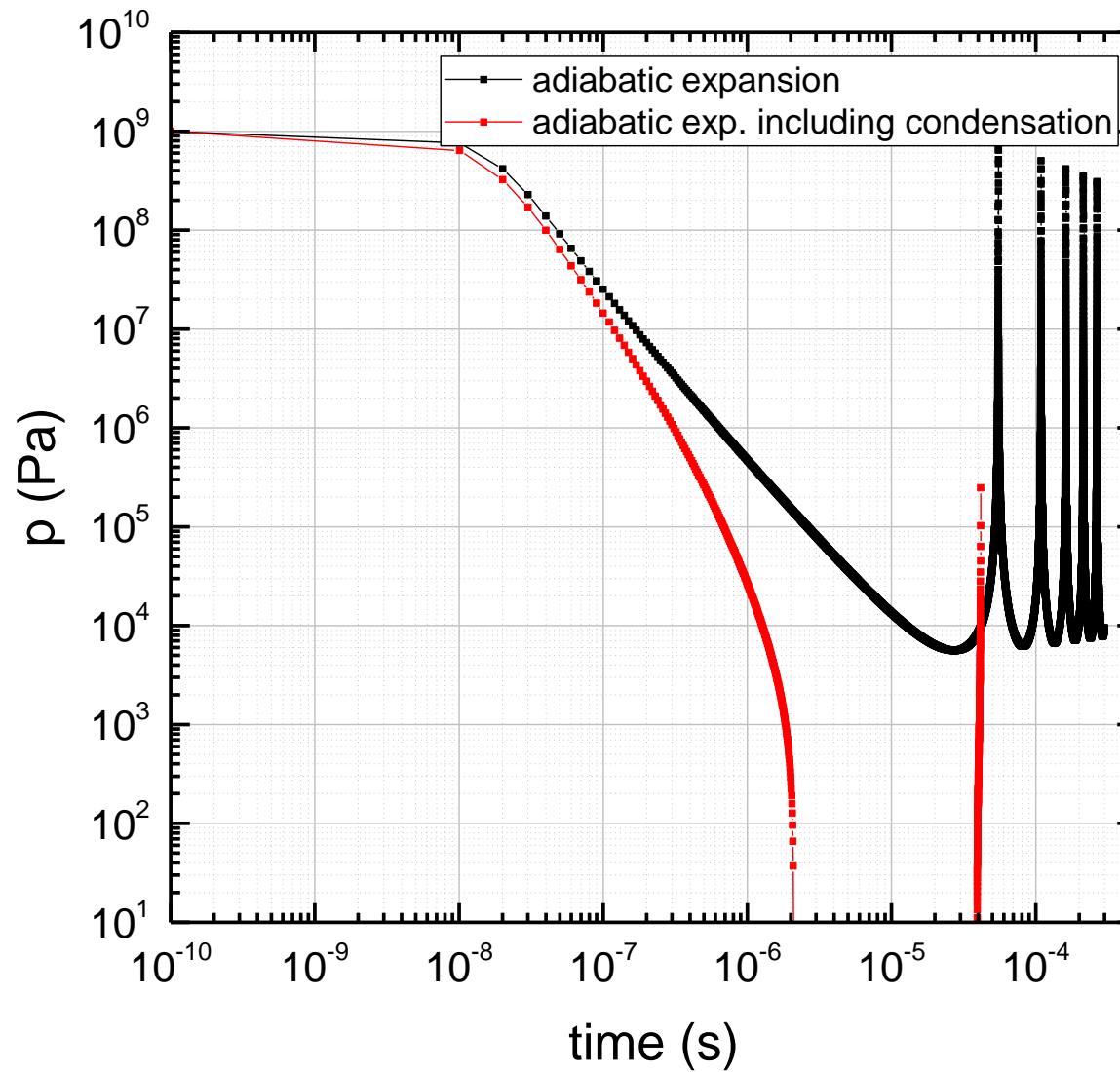
Balance between surface flux and volume loss depending on R,
 β =sticking coeff. = 1

$$\frac{1}{4} n v_{therm} 4\pi R^2 \beta = \frac{4\pi}{3} R^3 n \frac{1}{\tau}$$

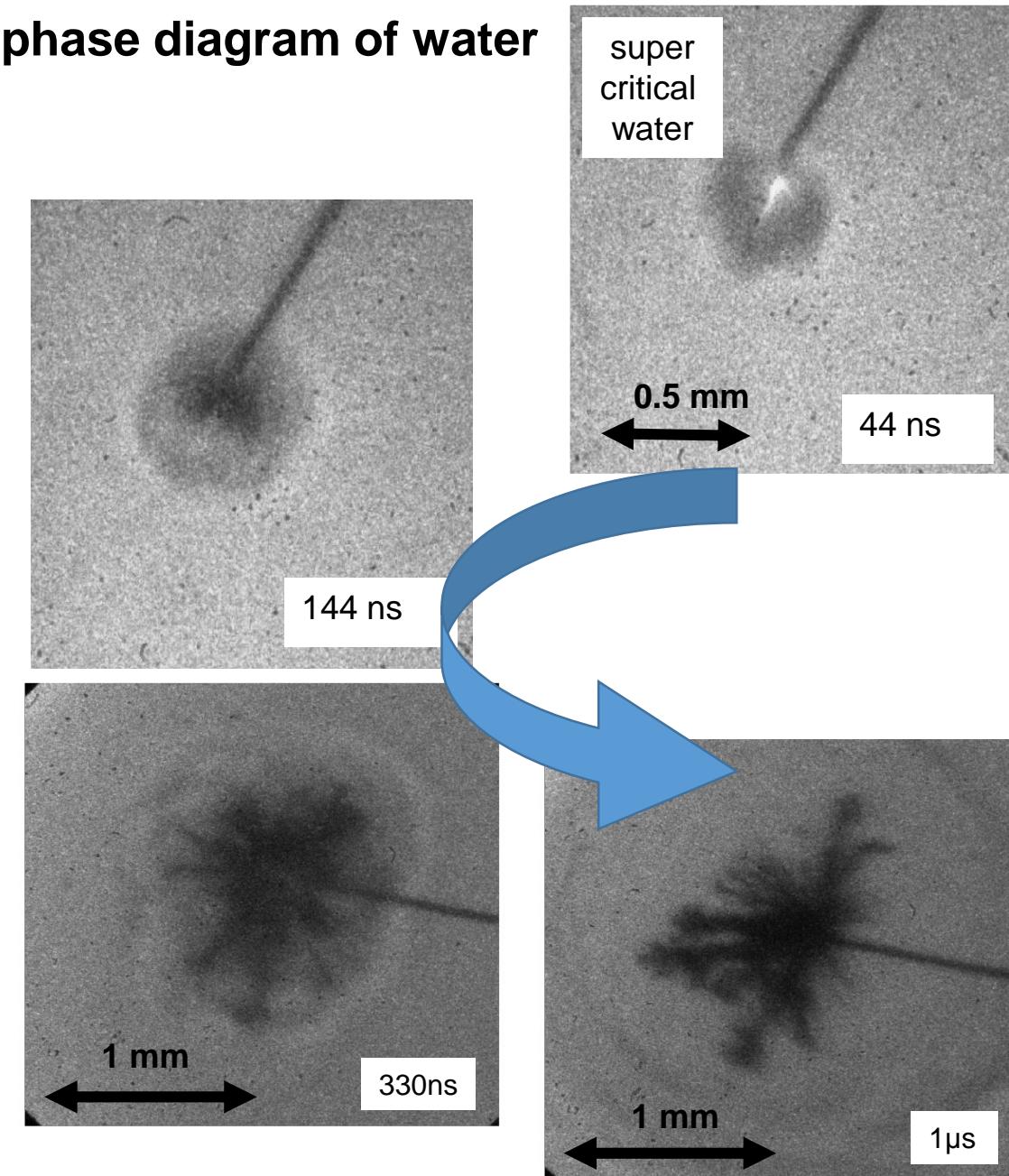
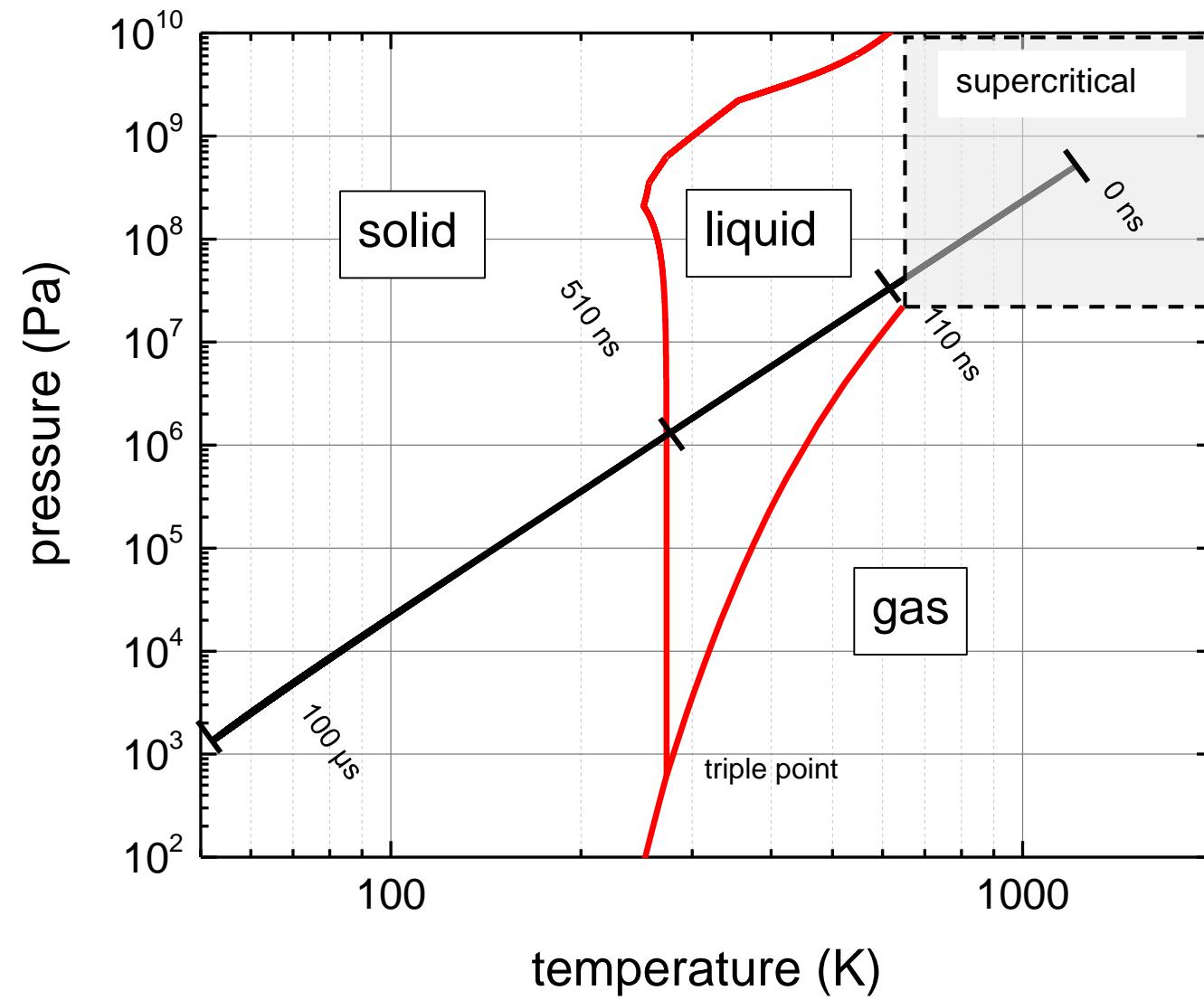
$$\tau = \frac{4R}{3v_{therm}\beta}$$

$$p_0 \rightarrow p_0 \exp\left(-\frac{t}{\tau}\right)$$

$$p(t) \rightarrow p_0 \exp\left(-\frac{t}{\tau}\right) \left(\frac{R_0}{R}\right)^{3\gamma}$$

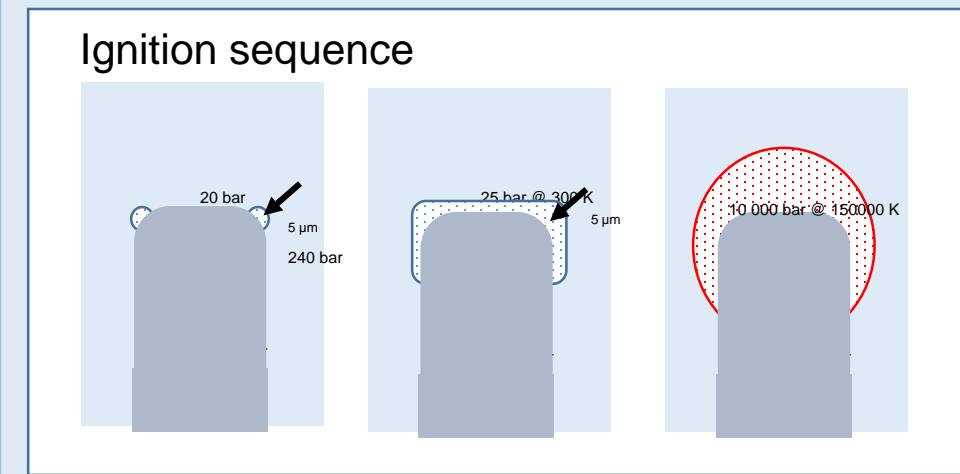


Transition of the water vapor in the bubble through the phase diagram of water



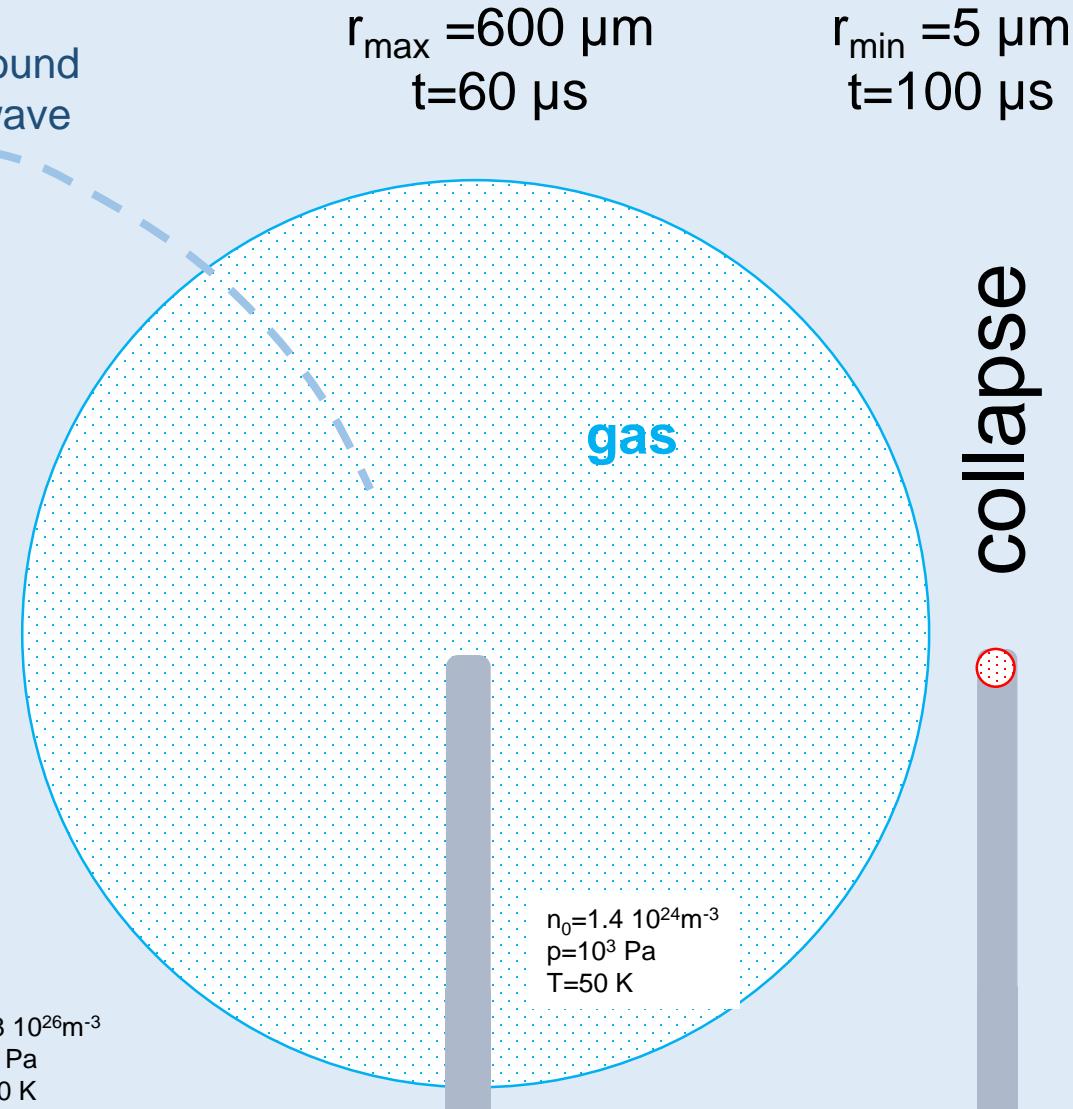
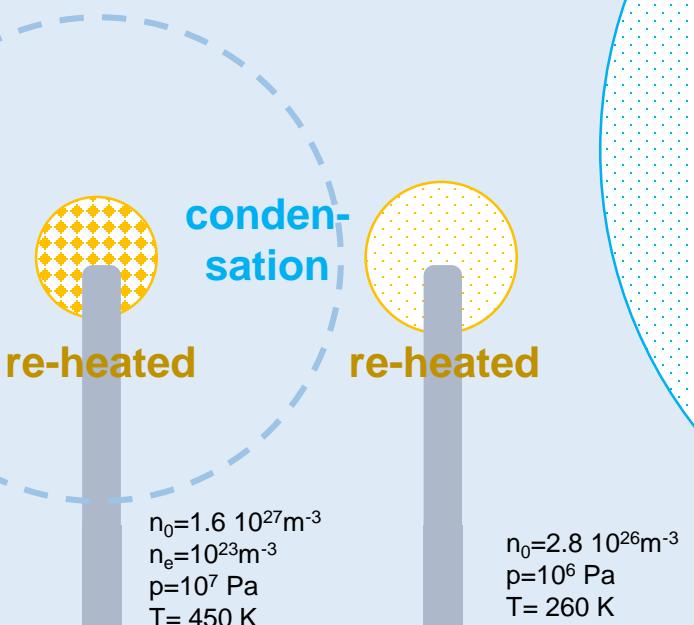
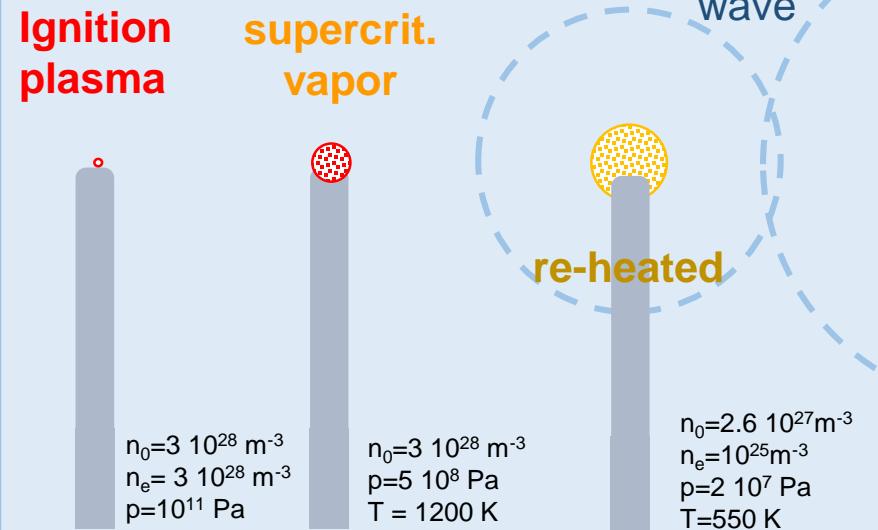
Plasma in a bubble expansion

everything to scale



$r = 80 \mu\text{m}$ $t = 200\text{ ns}$

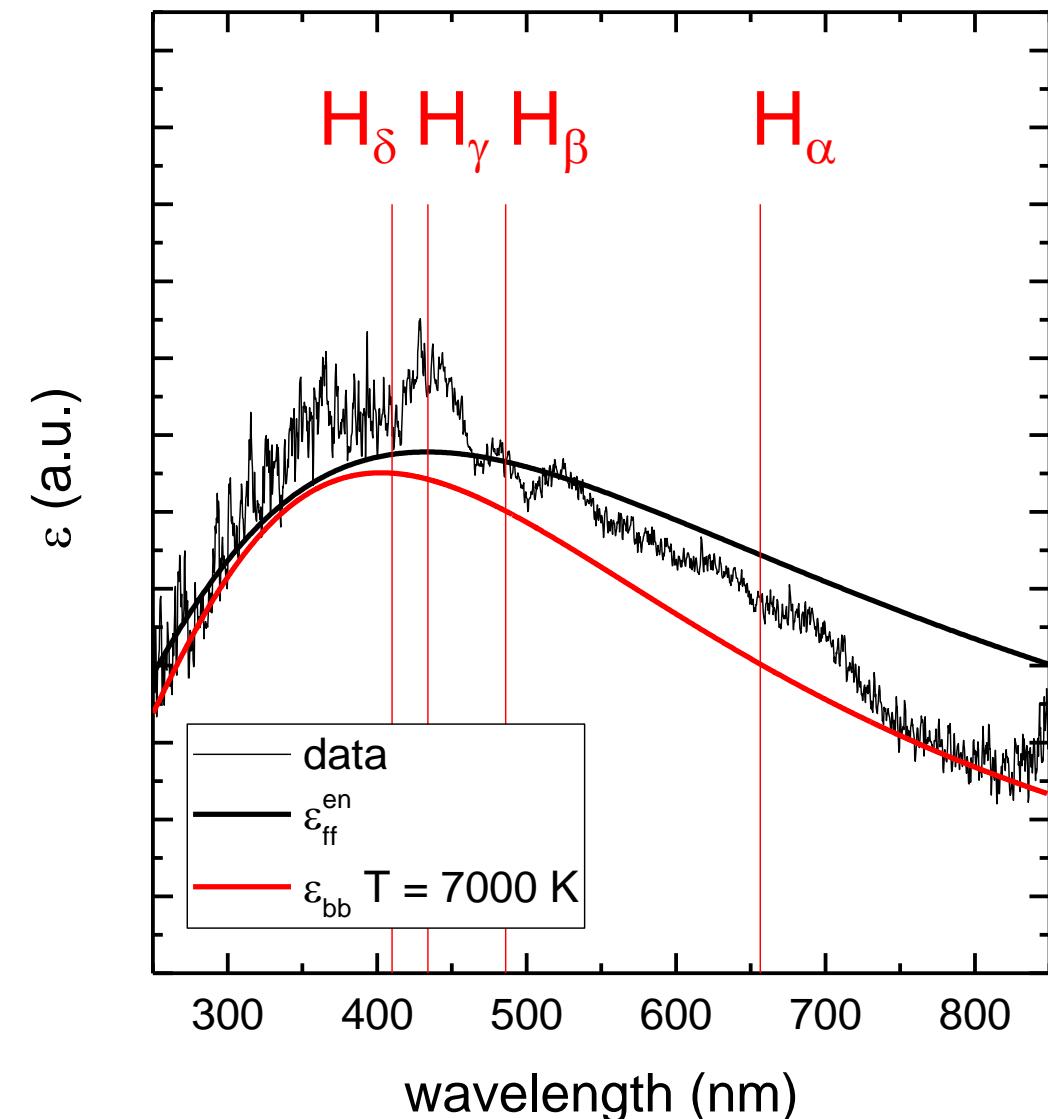
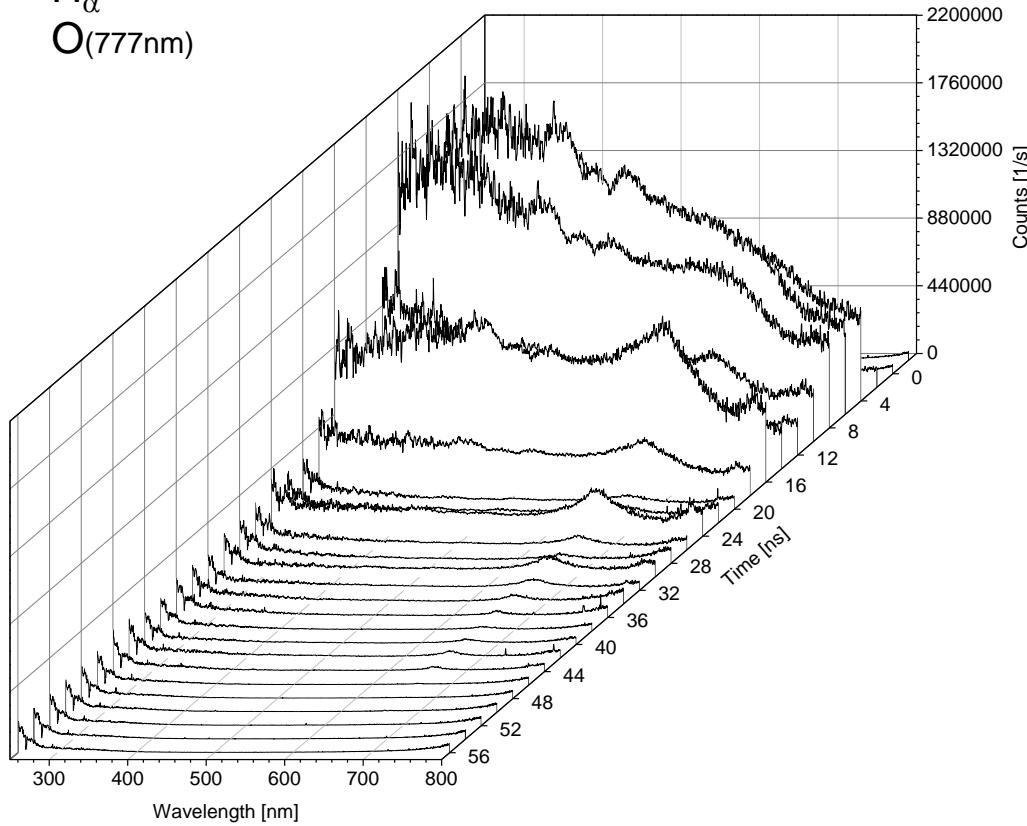
$r = 100 \mu\text{m}$ $t = 400\text{ ns}$



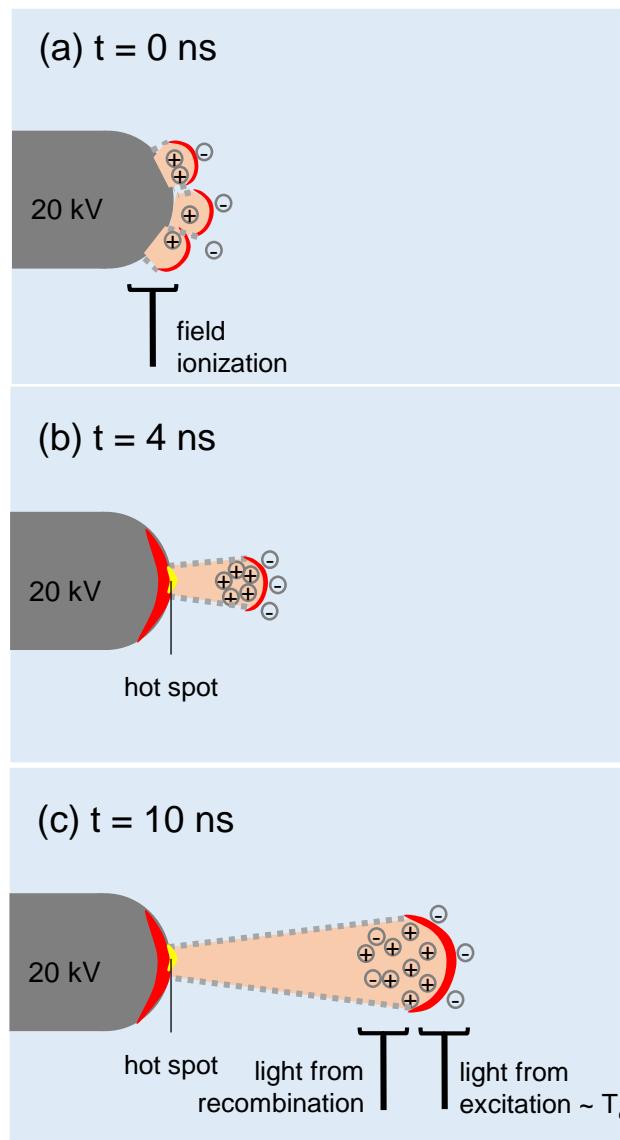
Temperatures from emission spectra of the plasma into the UV

Spectrum (first 50 ns, time resolution 2 ns)

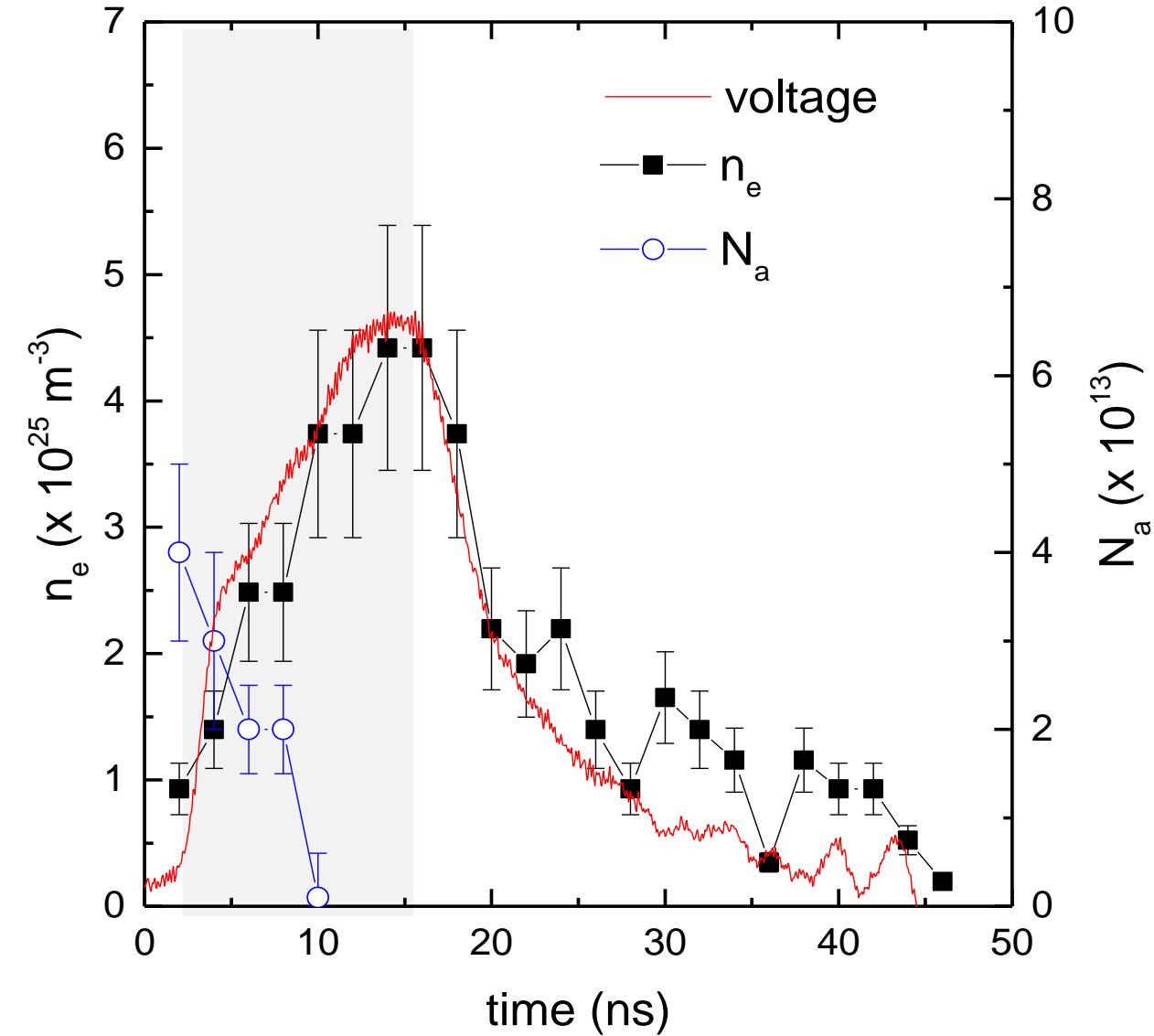
- Recombination
- H_2 , H_2O recomb. continuum
- OH, OH^+ , H_2O , H_2O^+ bands
- Black body radiation (W)
- H_α
- O(777nm)



Line emission – Ionization front propagates Hydrogen Balmer Series



self absorption of
Halpha from the
ionization region
disappears,
ionization front
steepens



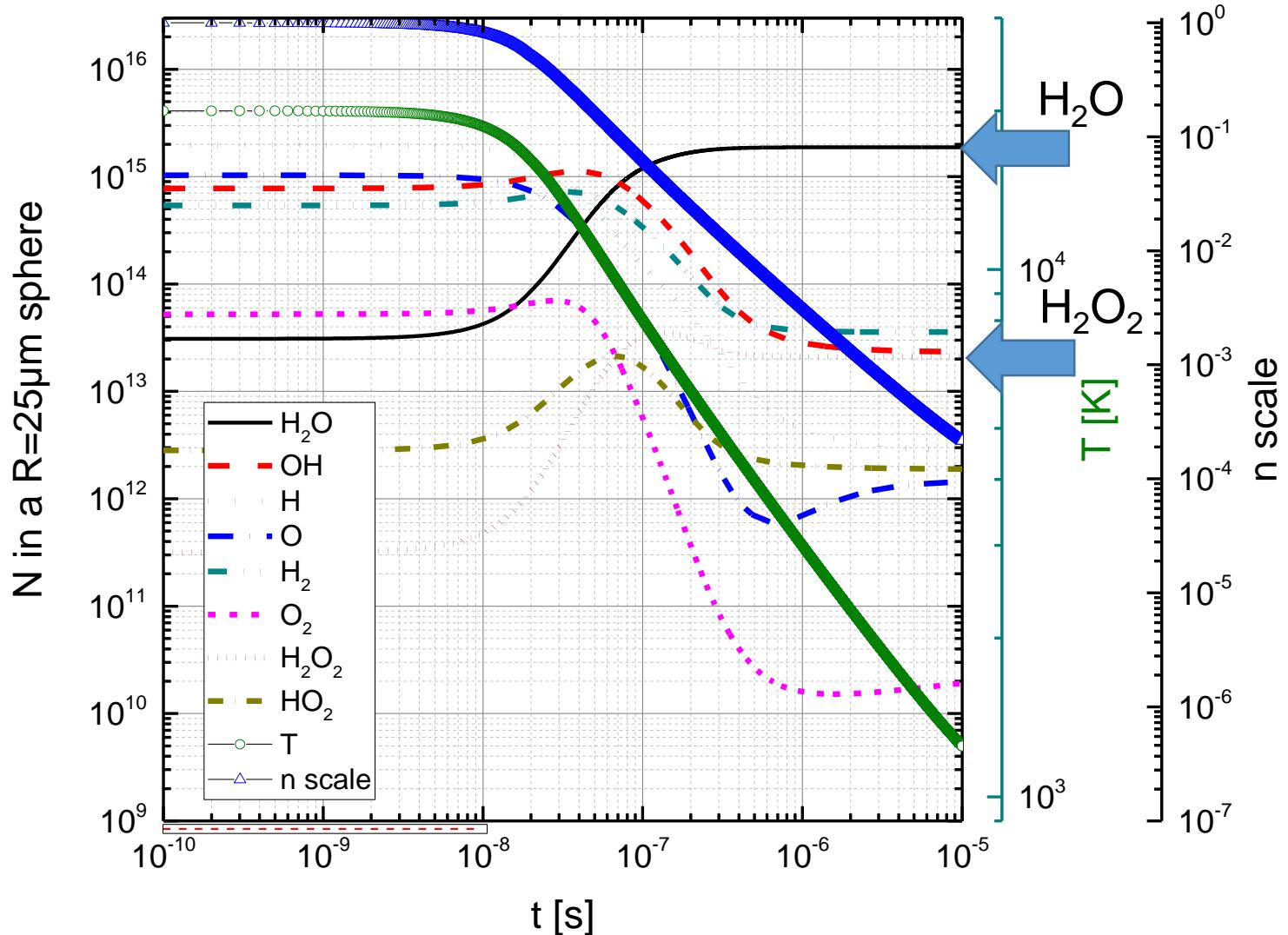
Efficiency to create H_2O_2 with these plasmas

Modeling the chemistry

- Starting point
gas density $3 \times 10^{22} \text{ cm}^{-3}$
- Starting Temperature $T = 20000 \text{ K}$
- $\text{H}_2\text{O}_2/\text{H}_2\text{O}$ ratio at $10 \mu\text{s} \sim 10^{-2}$

Hot phase:

- $\text{H}_2\text{O} \rightarrow \text{OH} + \text{H}$
- recombination phase
- $\text{OH} + \text{OH} \rightarrow \text{H}_2\text{O}_2$
 - $\text{H} + \text{H} \rightarrow \text{H}_2$
 - $\text{OH} + \text{H} \rightarrow \text{H}_2 + \text{O}$



In the expanding phase, chemistry freezes out, $\text{H}_2\text{O} \rightarrow \text{H}_2\text{O}_2 + \text{H}_2 + \text{O} + \text{H}$

*S. Mededovic, B. Locke
JPD 40, 7734 (2007)
PCPP 32, 875 (2012)

Efficiency to create H_2O_2 with these plasmas

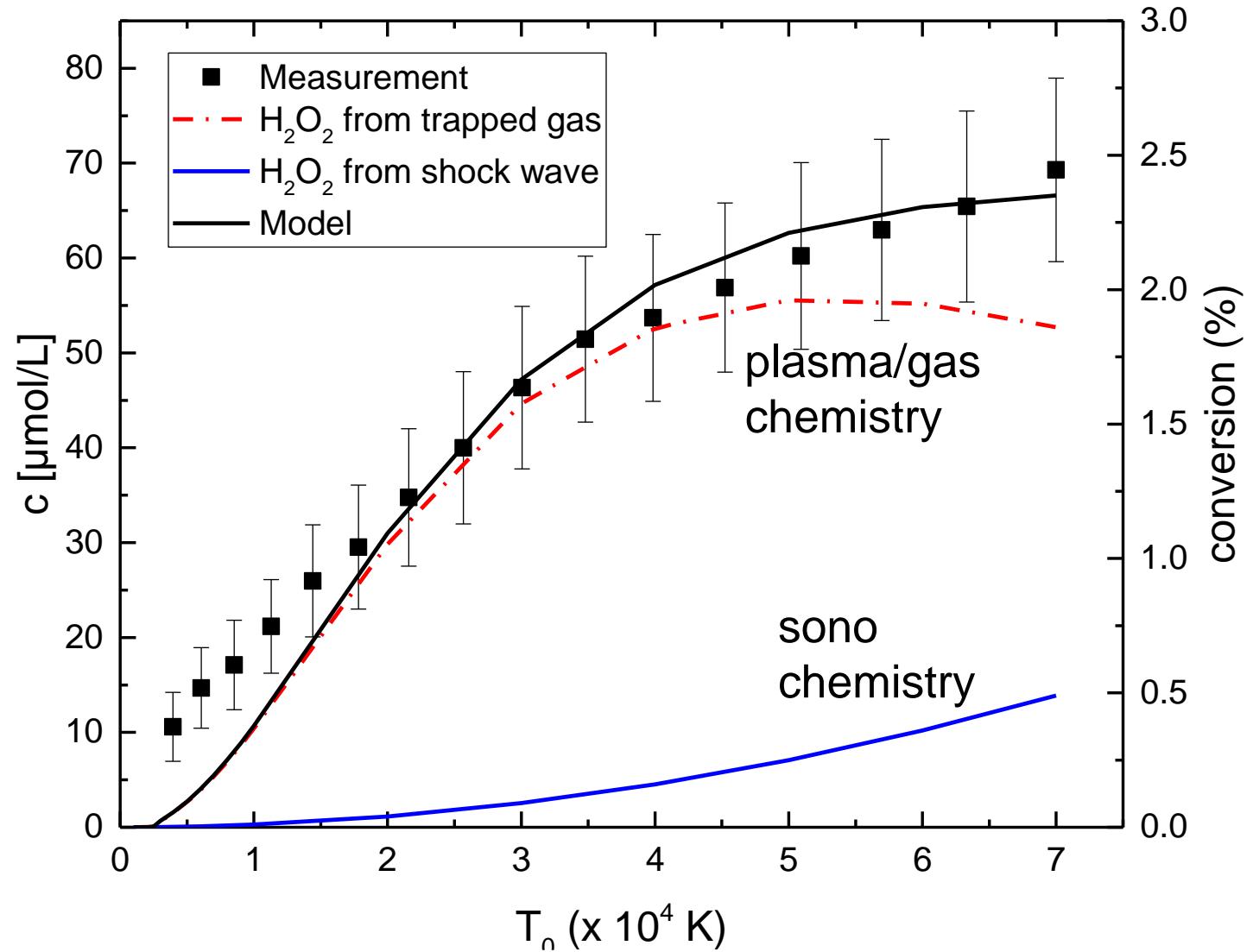
Modeling the chemistry

- Take T_0 from cavitation
- calculate time development of the chemistry until 10 μs
- determine H_2O_2 concentration after 10 μs in the model

efficiency 1.1%
@ 20 kV (20000 K)

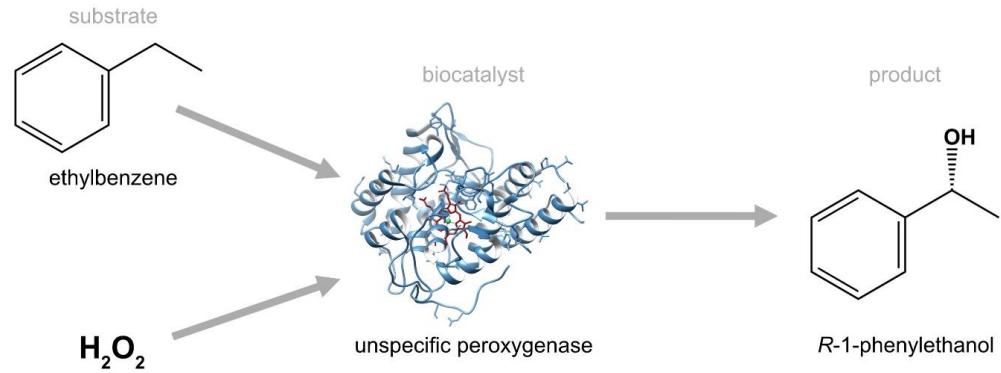
- At high voltages/ T_0 , contribution due to sonochemistry becomes significant

46 g $\text{H}_2\text{O}_2/\text{kWh}$
(most other plasma methods only a few g/kWh)



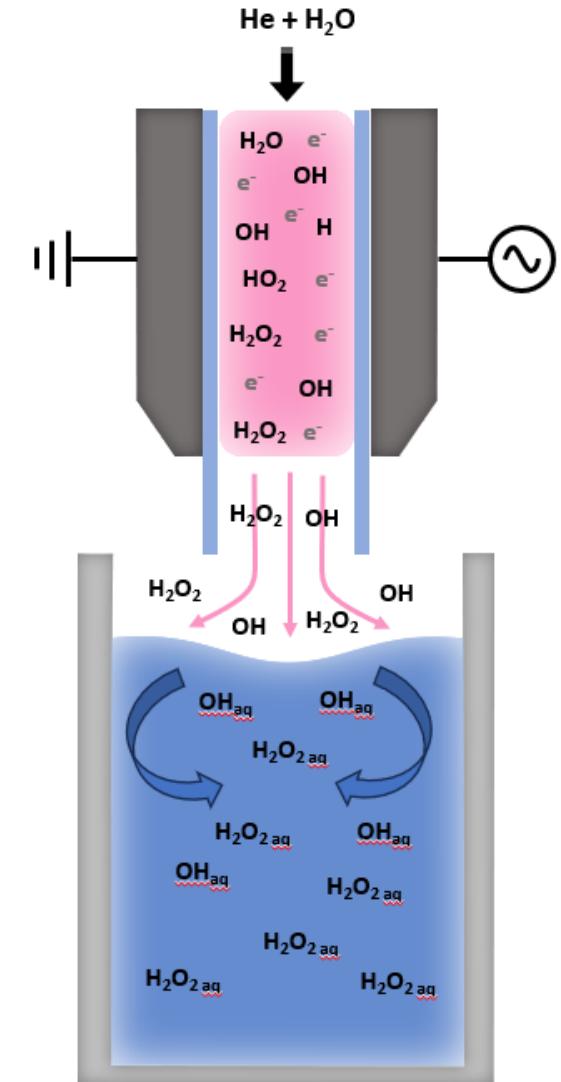
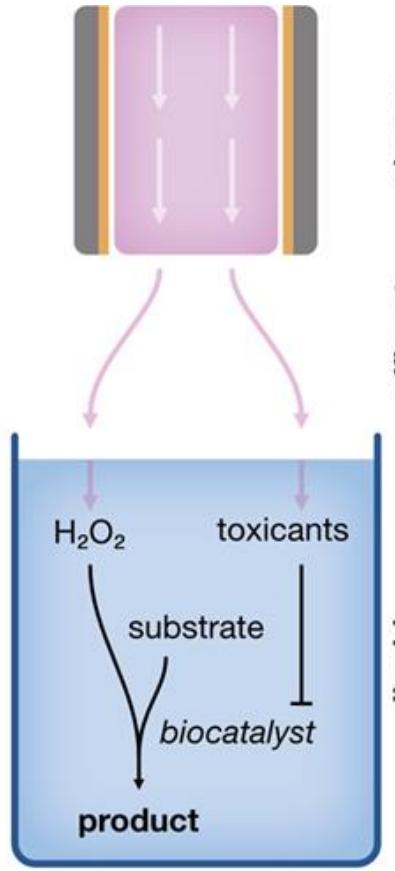
advantages of enzymatic reactions

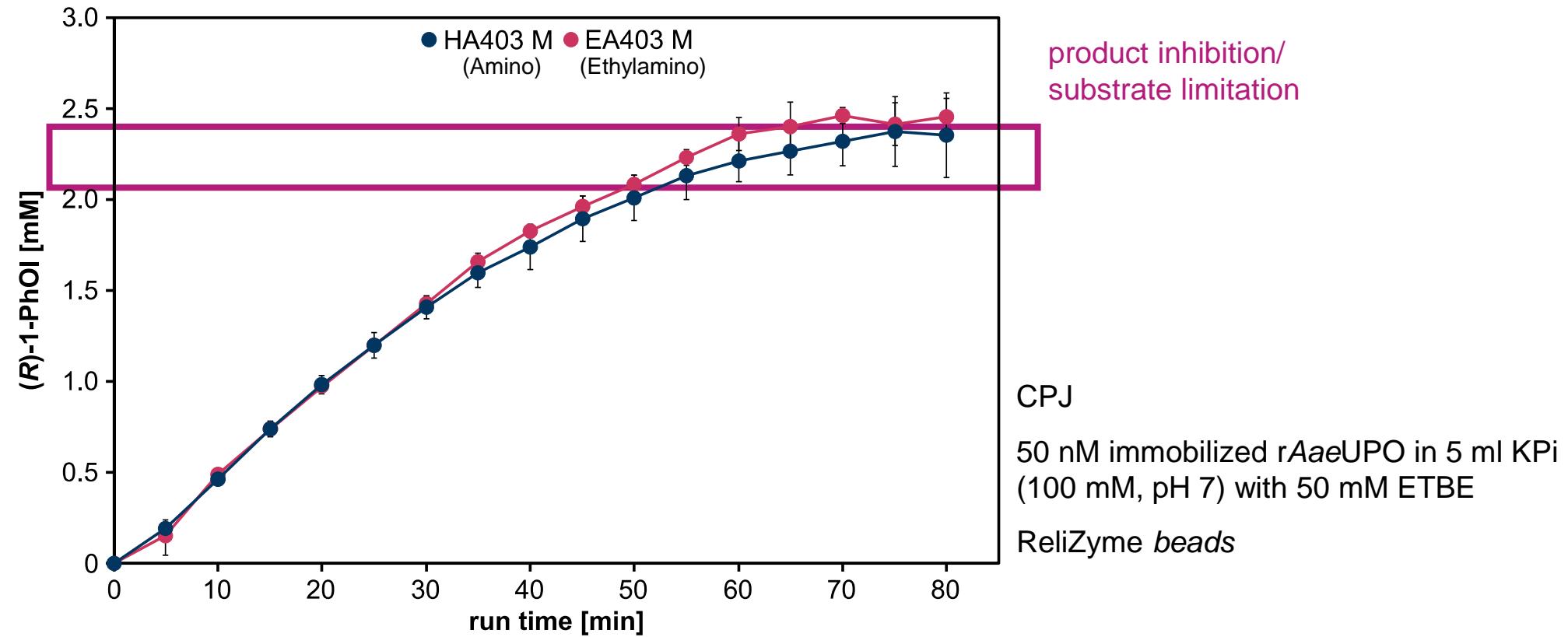
- mild conditions
- less by-products / waste
- chemo- and stereo-selectivity



advantages of plasma

- easily turned on/off
 - H_2O_2 is also toxic for enzymes
- non-invasive
- not increasing the reaction volume

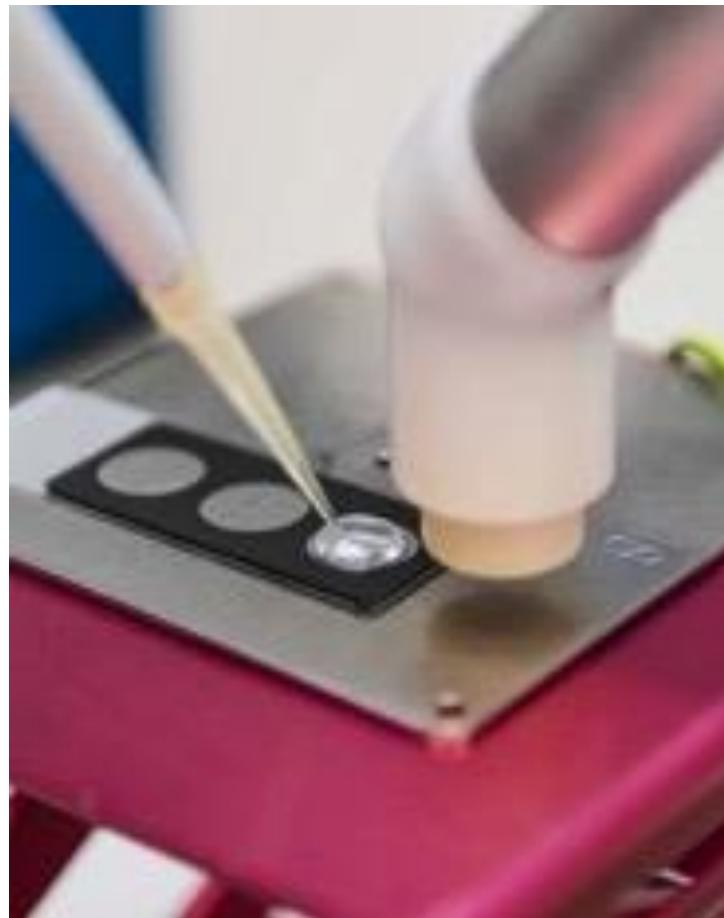




- HA403M, EA403 M on the same level of product formation → more bead types usable in biocatalysis (further tested beads not on the same level)
- plateau reached (probably due to product inhibition and substrate limitation)

Different plasma sources (Bandow, Golda, Gibson group @ RUB)

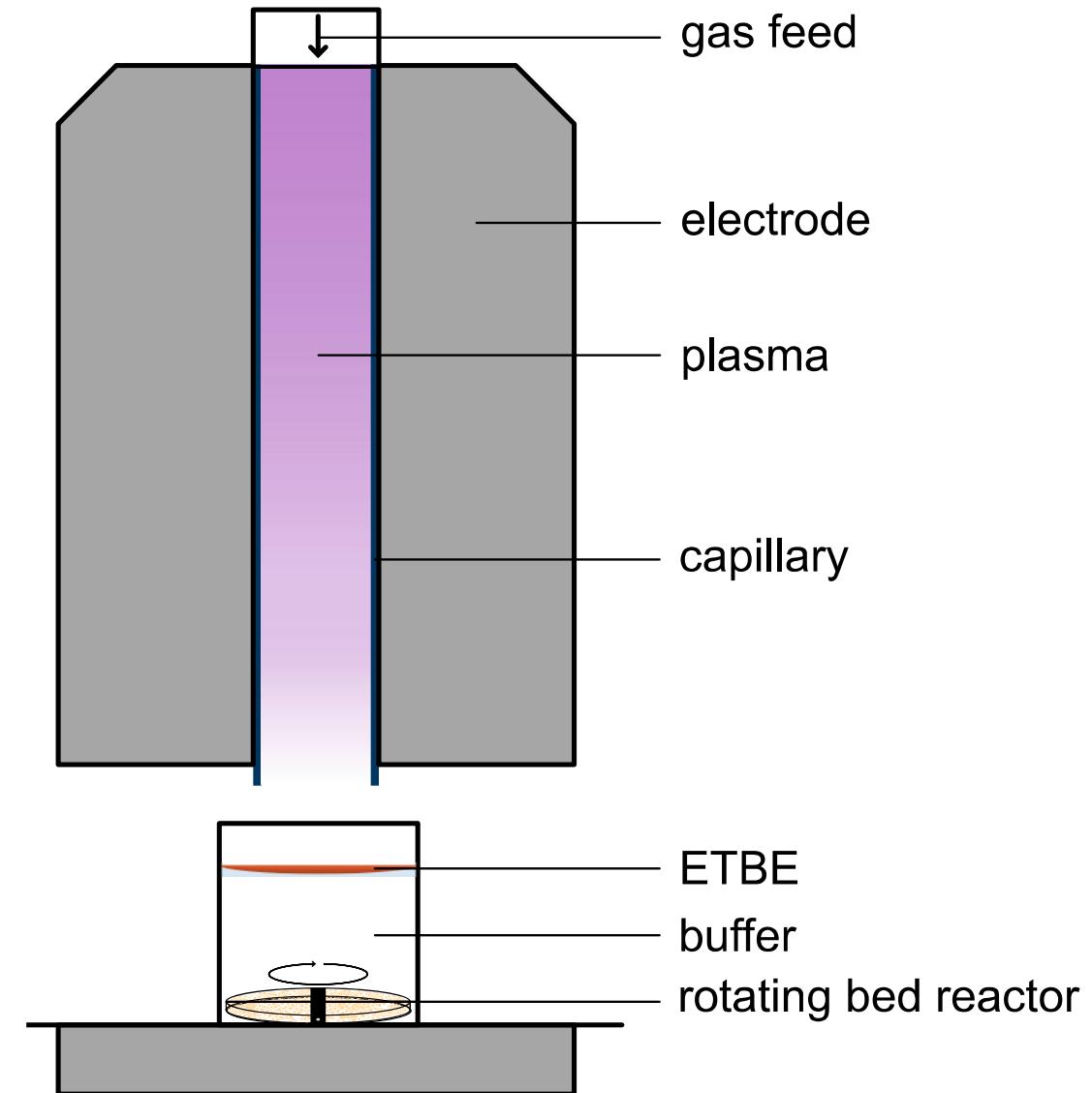
DBD



Capillary jet

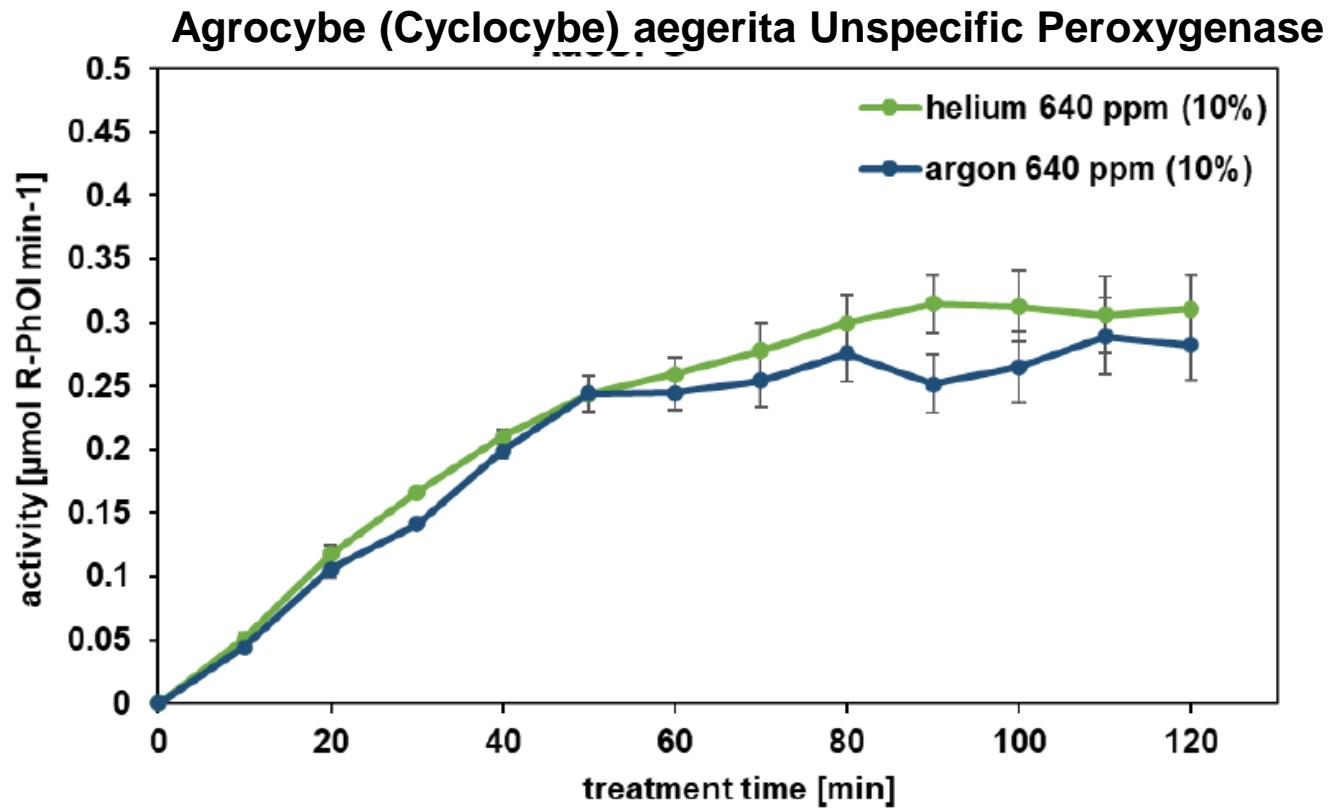
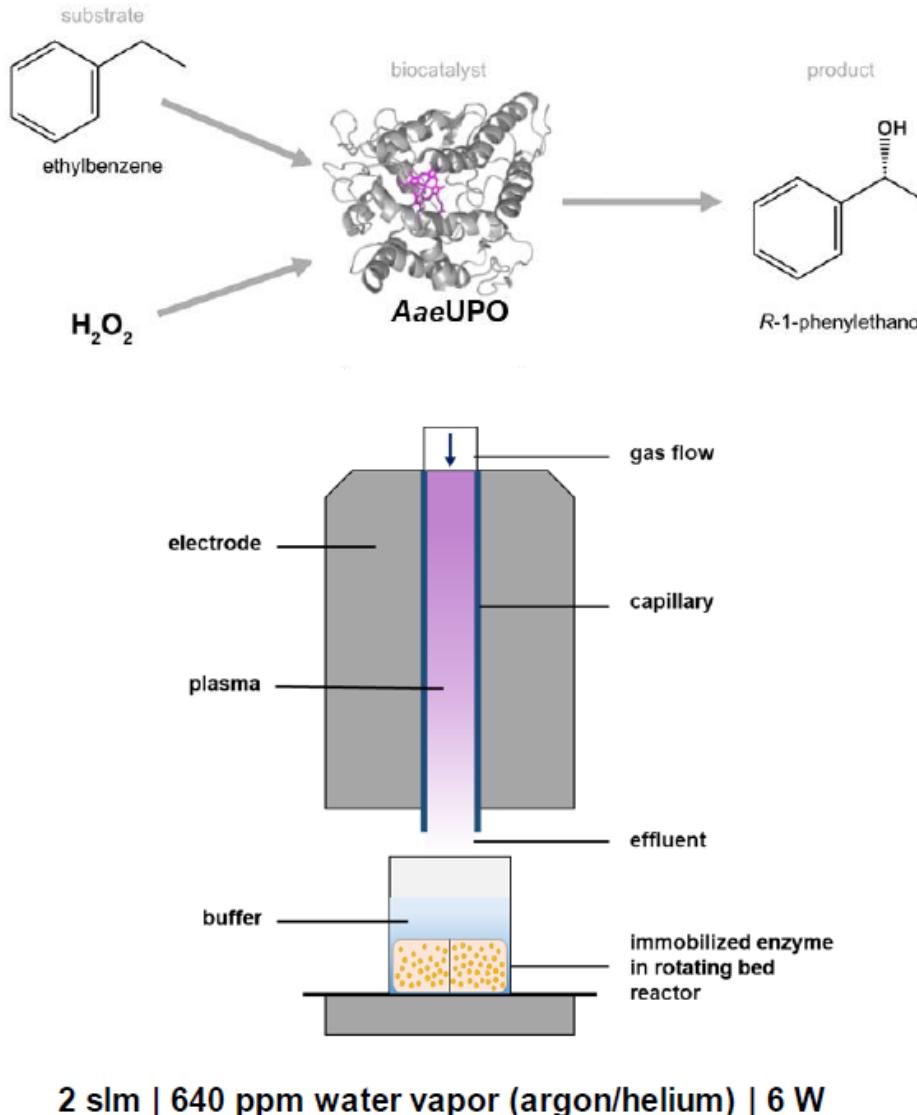


Capillary jet, rotating bead reactor



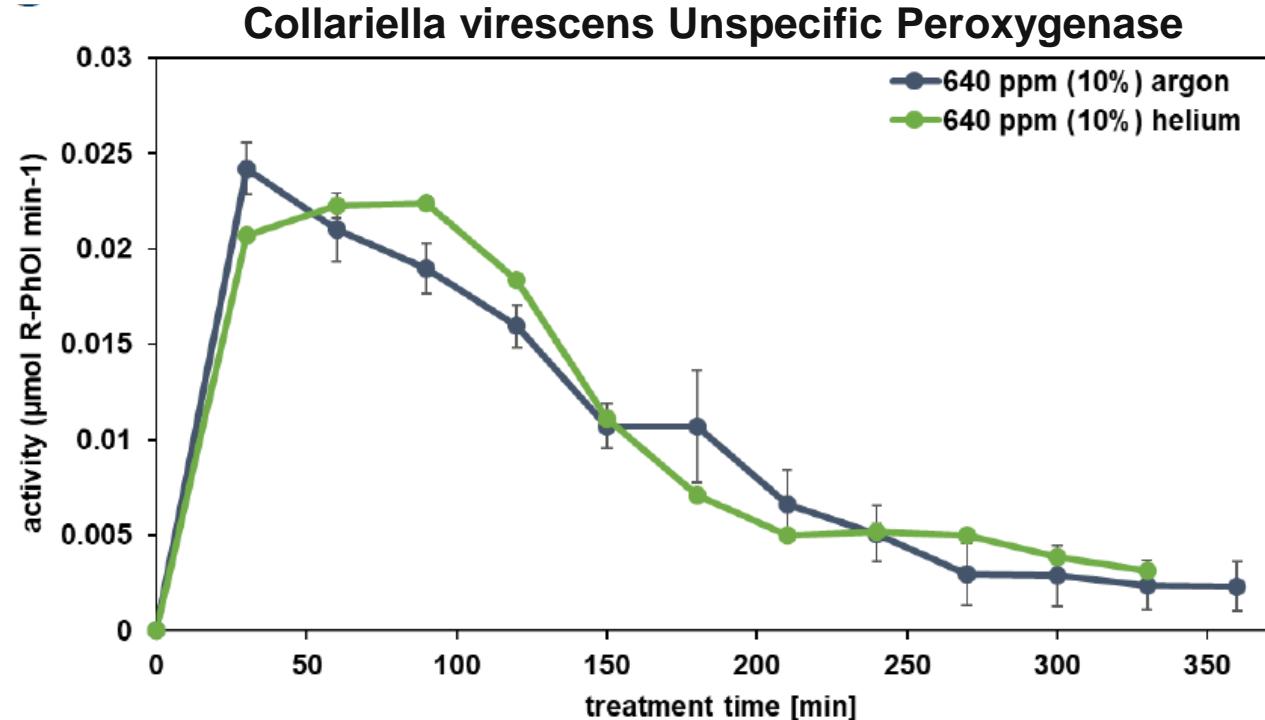
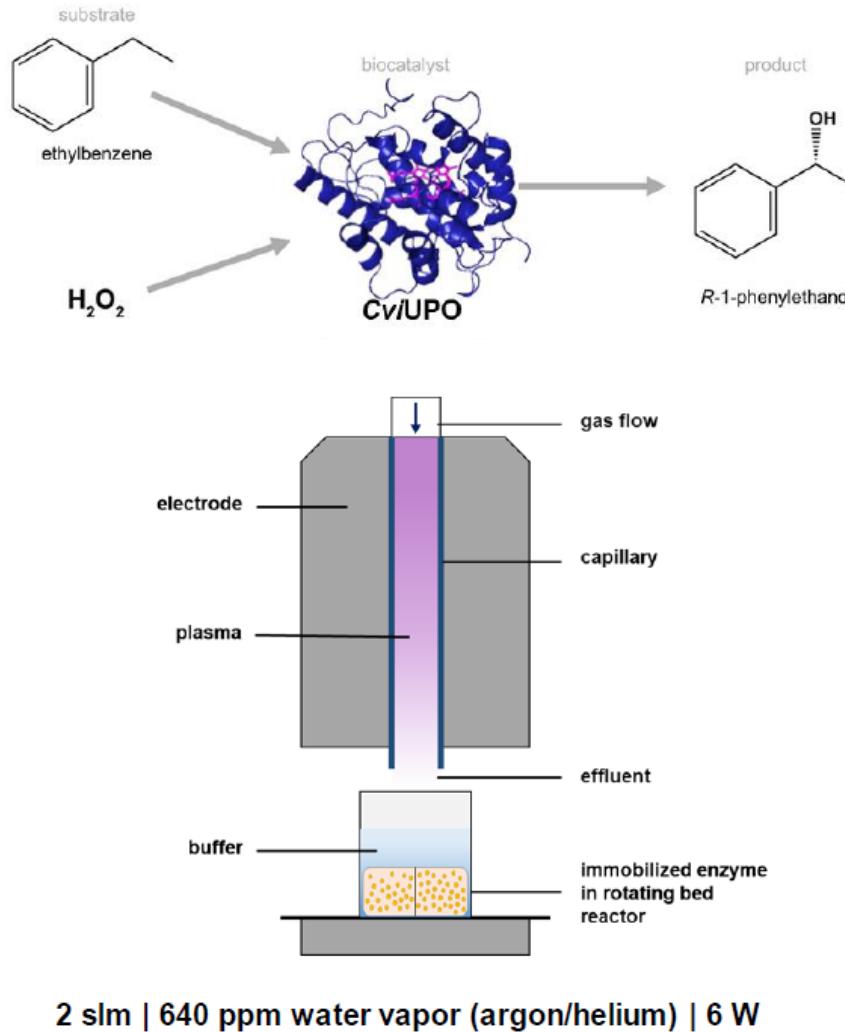
Yayci et al. (2020) *ChemSusChem*.

Plasma driven biocatalysis (group Bandow @ RUB)



Argon and Helium | 640 ppm water vapor (10%)
→ No significant difference in product concentration after plasma-driven biocatalysis with argon and helium as feed gas

Plasma driven biocatalysis (group Bandow @ RUB)

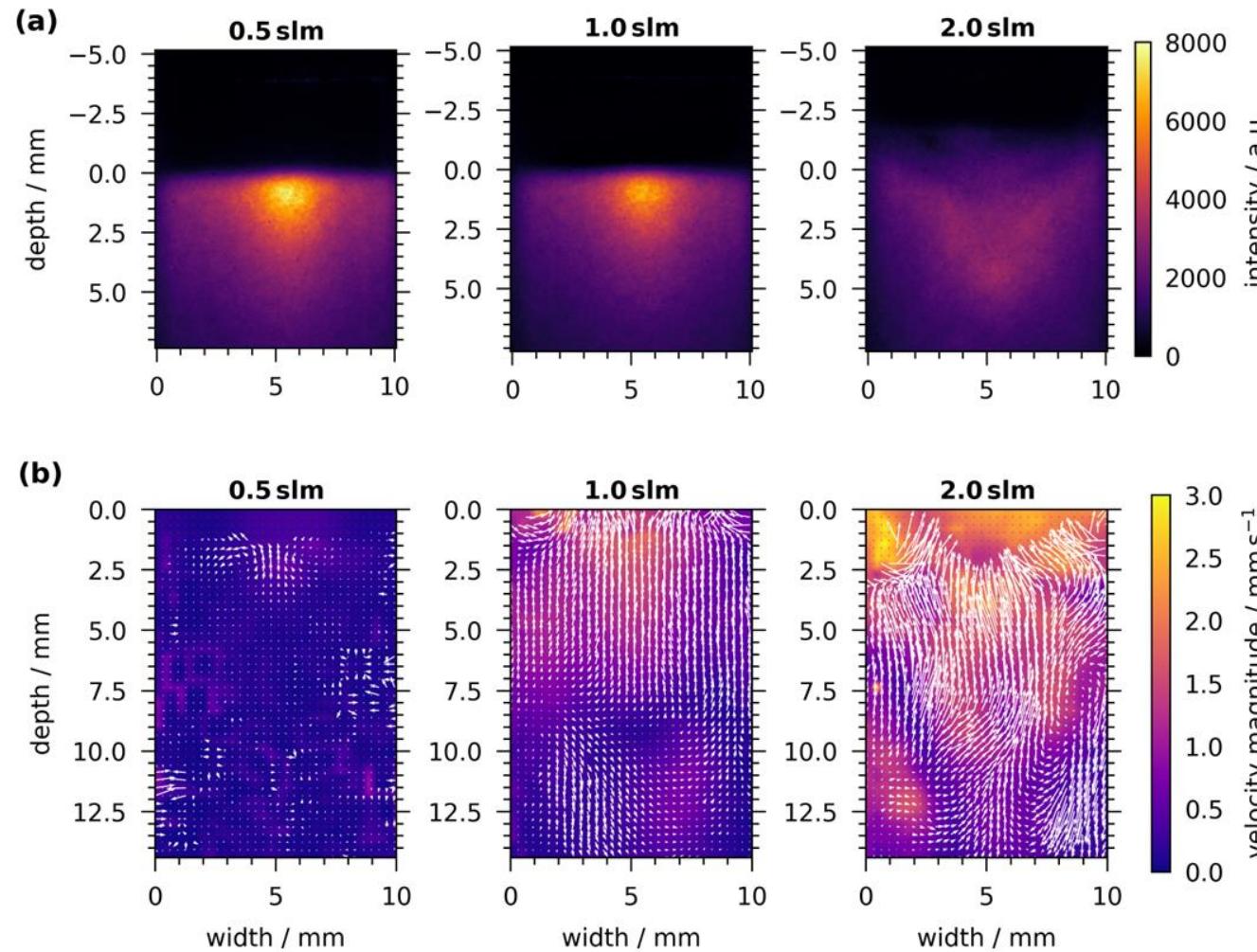


The use of **argon** as feed gas in plasma-driven biocatalysis is possible and leads to similar product formation

Helium → $0.004 \mu\text{mol R-Phol min}^{-1}$ after 350 min

Argon → $0.003 \mu\text{mol R-Phol min}^{-1}$ after 360 min

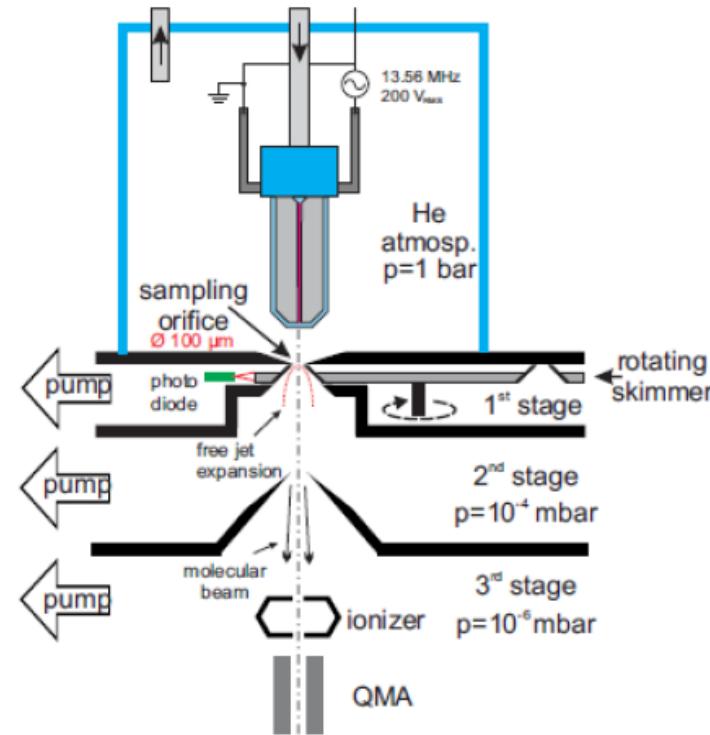
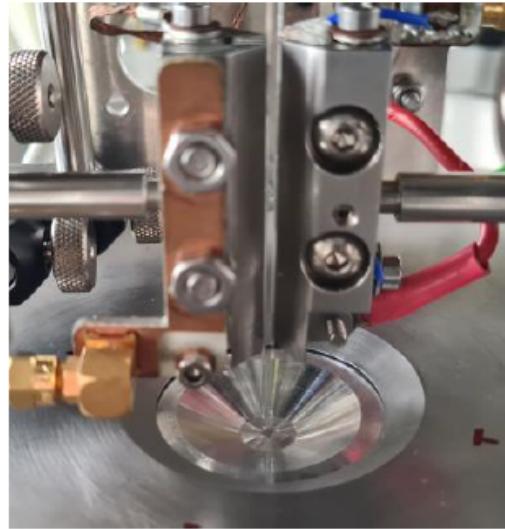
Diffusion of OH and convective transport (Golda group @ RUB)



- CL of luminol
→ Main signal at the liquid surface
- Decrease in CL signal deeper than 2.5 mm
- OH is mainly distributed at the liquid surface

- PIV measurements → Visualization of convective transport
- Enhanced convective transport at higher gas flow rates
- At 0.5 sLM and lower → Almost no convective transport

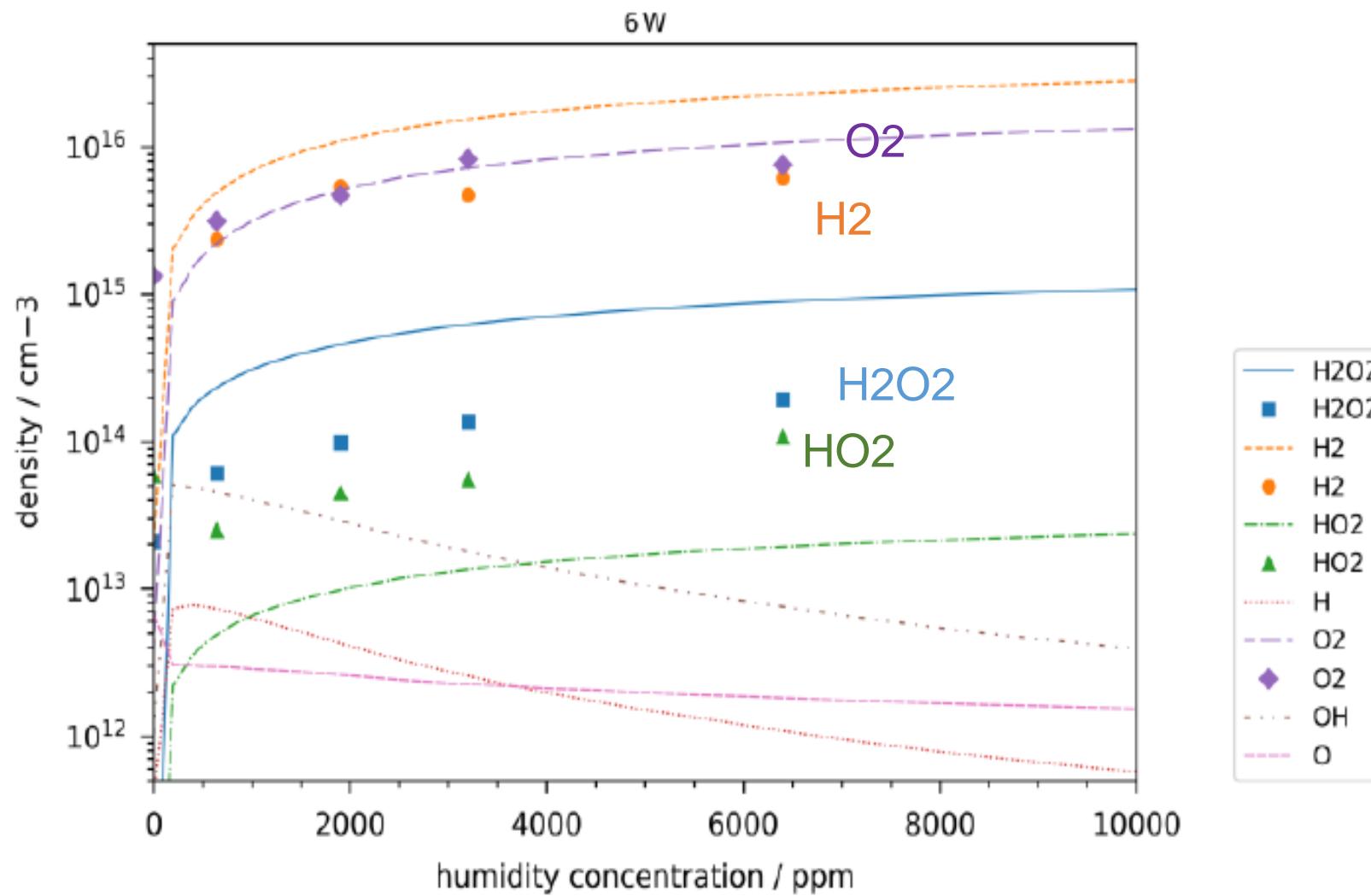
Plasma driven biocatalysis (group Benedikt @ RUB)



Dissertation Gert Willems, Ruhr-University Bochum (2016)

Species	Calibration
H ₂ O	H ₂ O
H ₂ O ₂	Ar
OH	-
O	-
O ₂	Ar
H	-
H ₂	H ₂
HO ₂	Ar

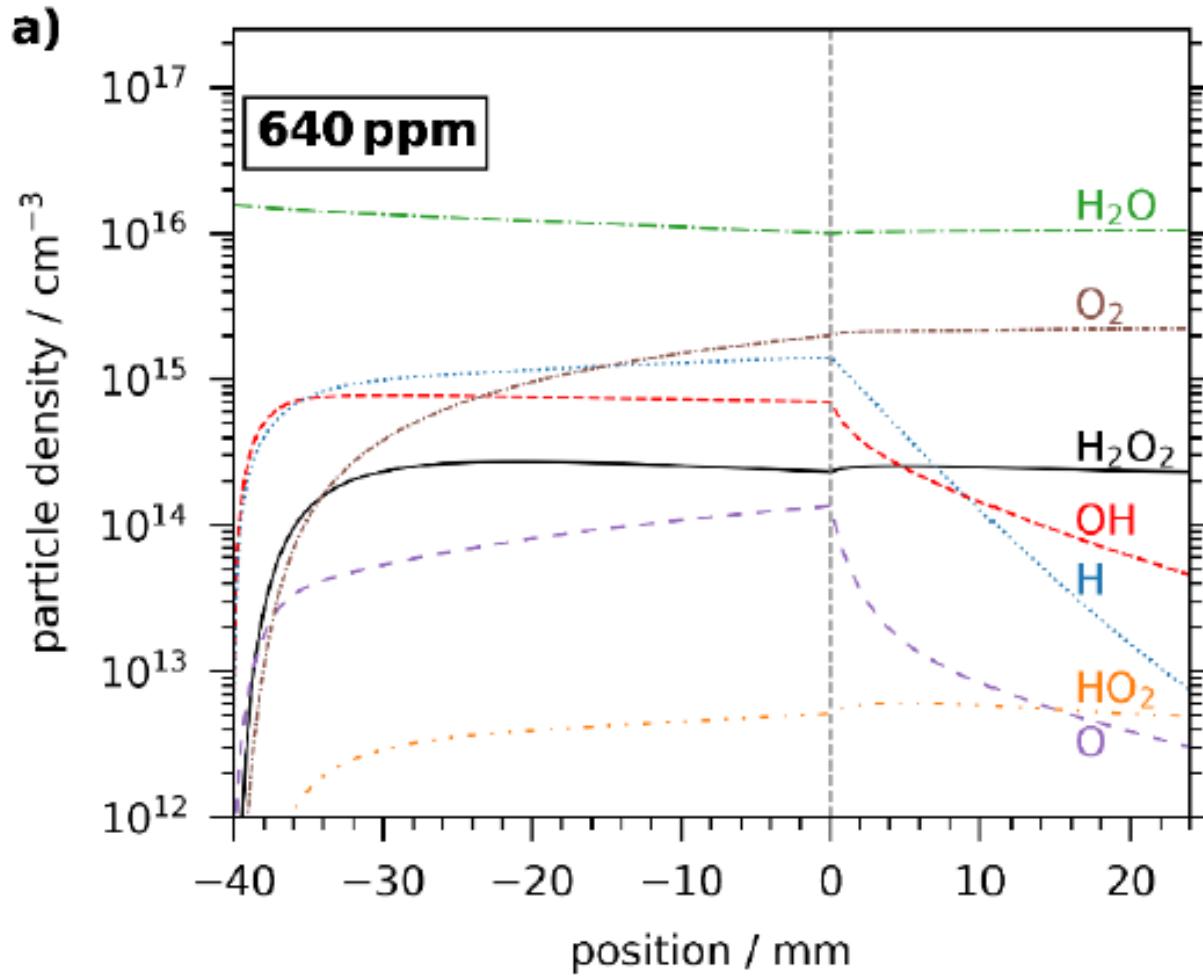
Plasma driven biocatalysis (group Golda, Gibson @ RUB)



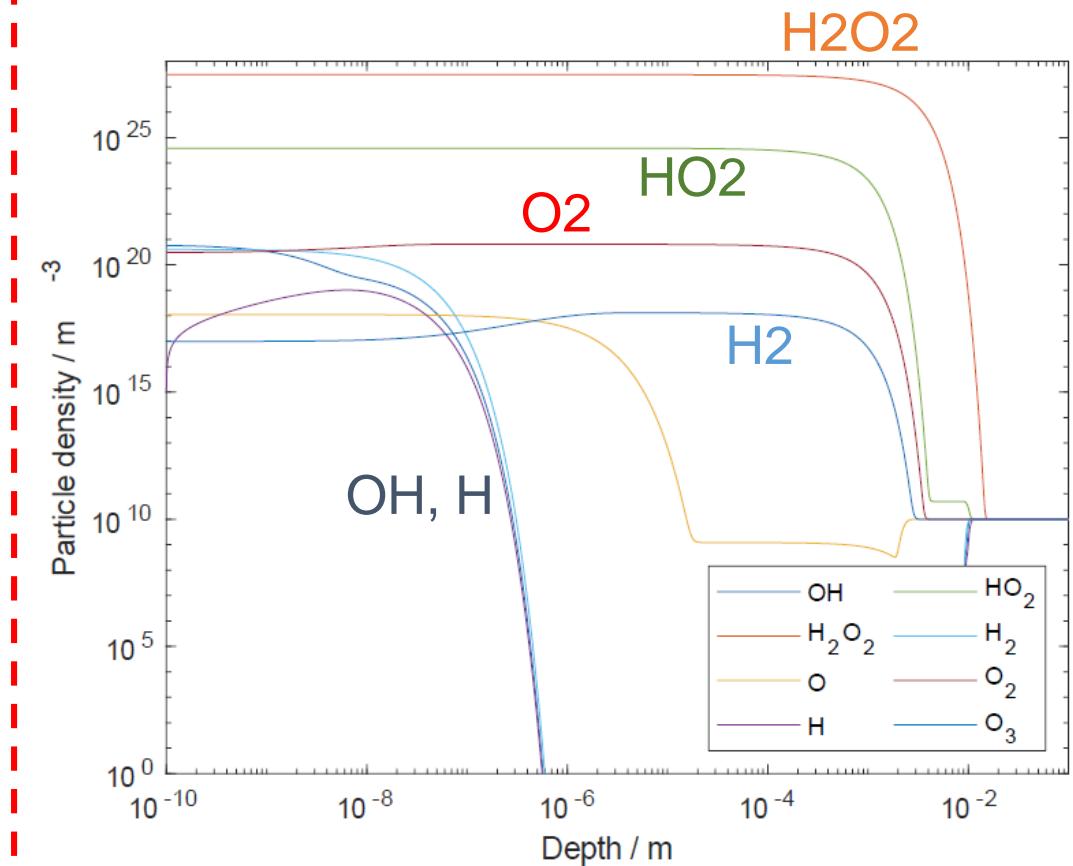
Plasma effluent modeling with 100s of reactants and several 100s of reactions

Agreement only within a factor of 10

In the gas phase



In the liquid phase

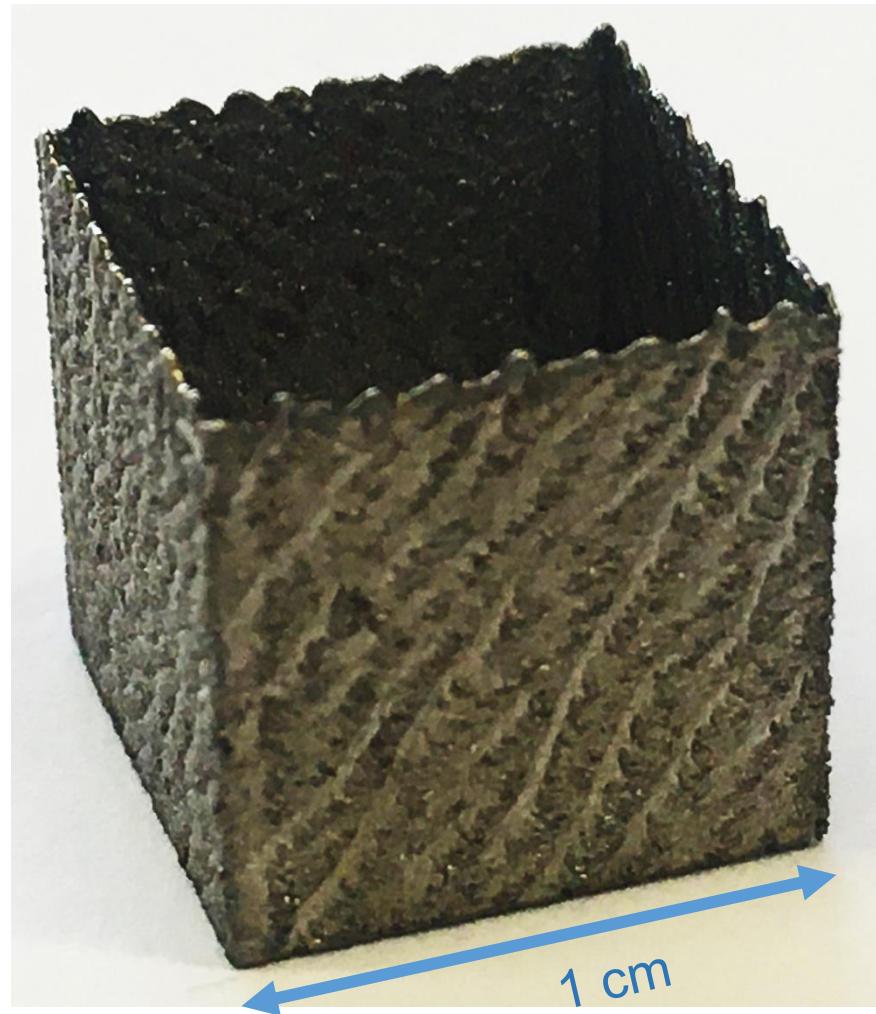


Conclusion

plasmas in liquids for life science applications

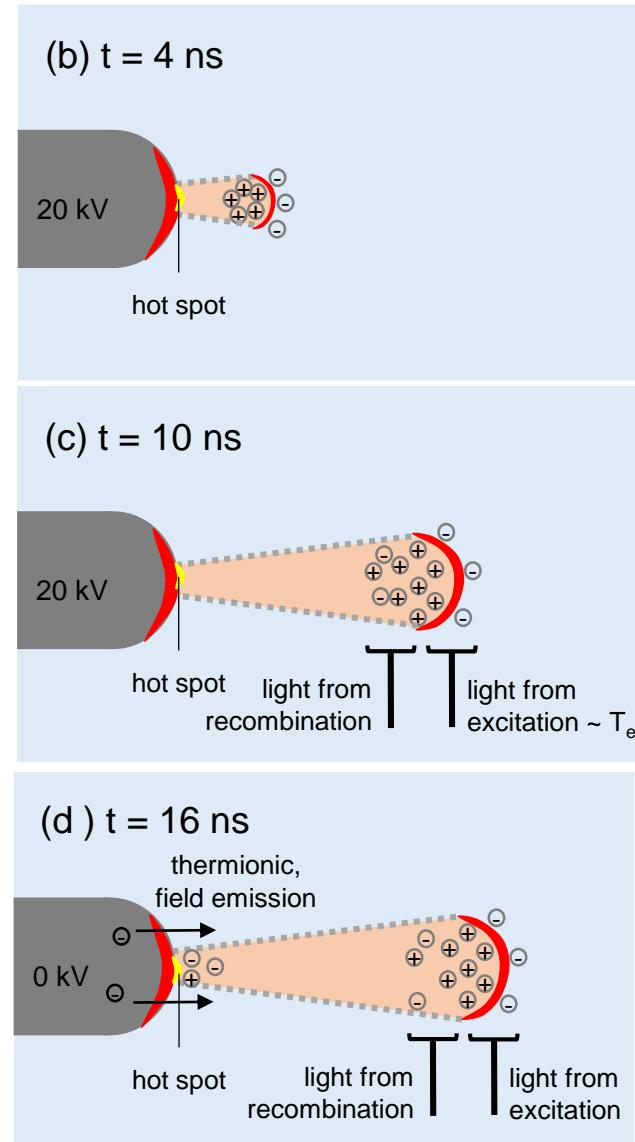
- Plasma-in-liquids are very efficient, avoiding RONS and H₂O₂ decay.
Physics is extremely complex
- Plasma above the liquid suffer from transfer (Henry constant). More applications.
- Plasmas act as a on-demand source for H₂O₂
- Design of the perfect plasma source for bio application is crucial, right ‘dose’ of reactive species

Backup Slides

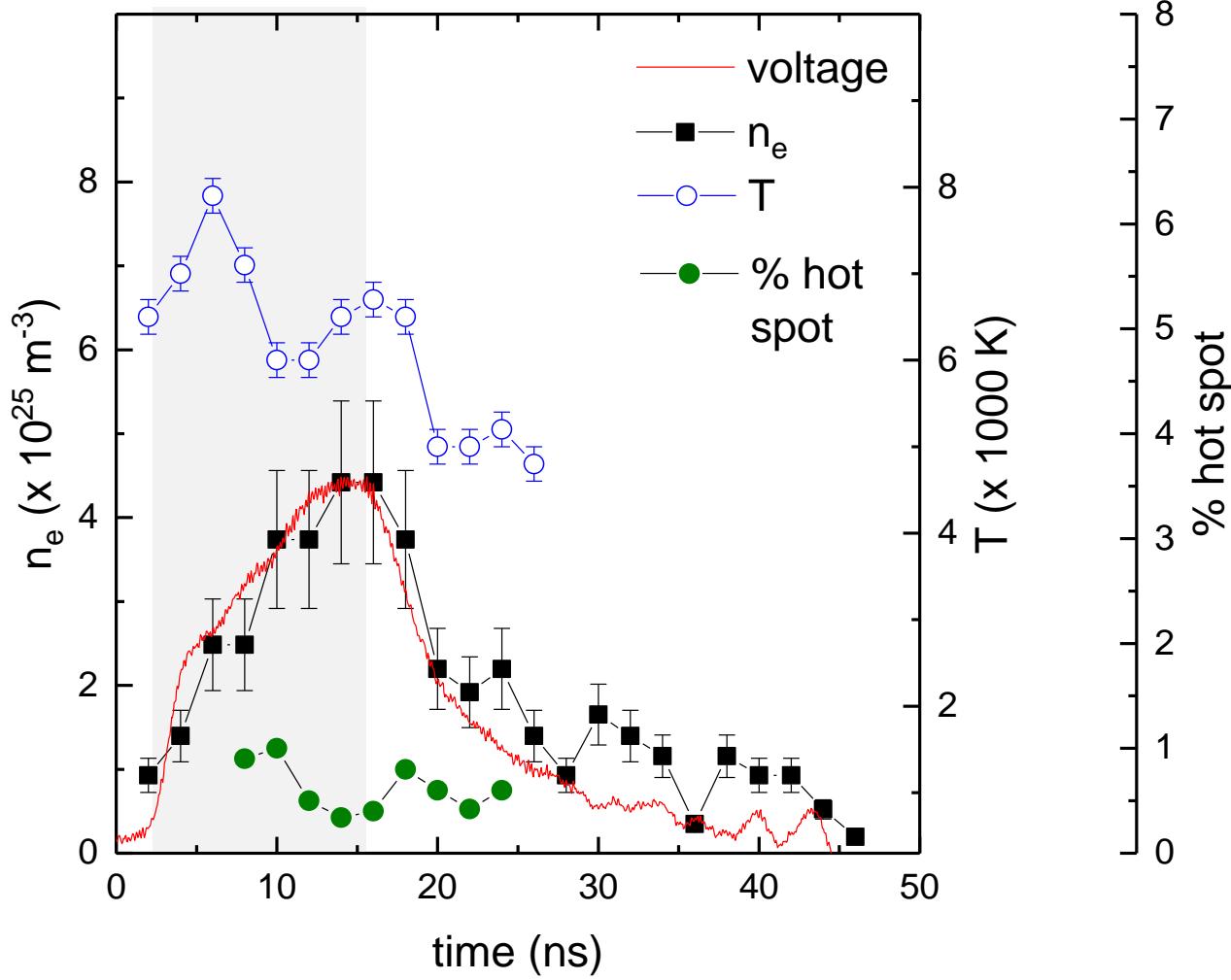


1 cm

Line emission – H alpha from ionization region of Hydrogen Balmer Series



Temperature variation

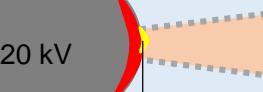


Line emission – H alpha from ionization region of Hydrogen Balmer Series

(b) $t = 4 \text{ ns}$



(c) $t = 10 \text{ ns}$



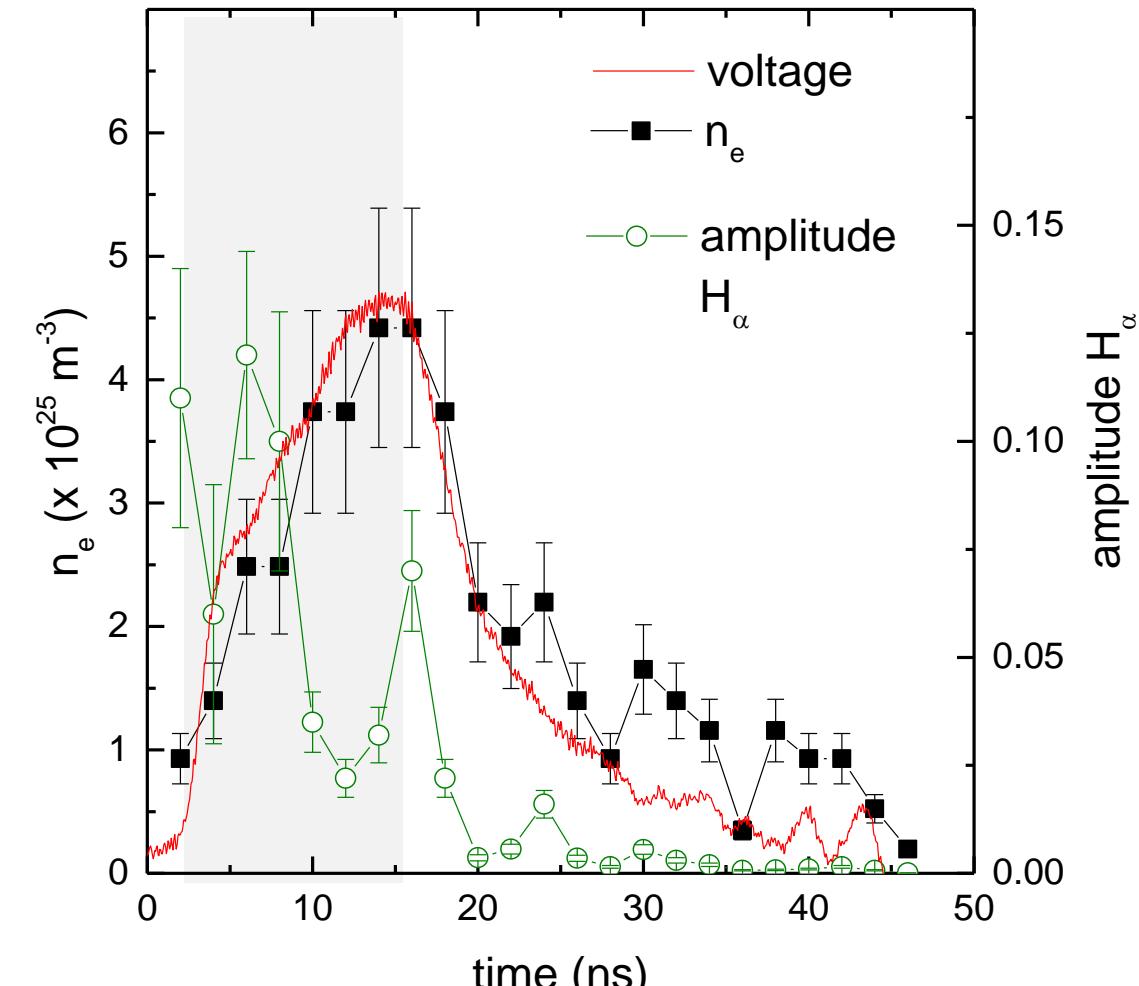
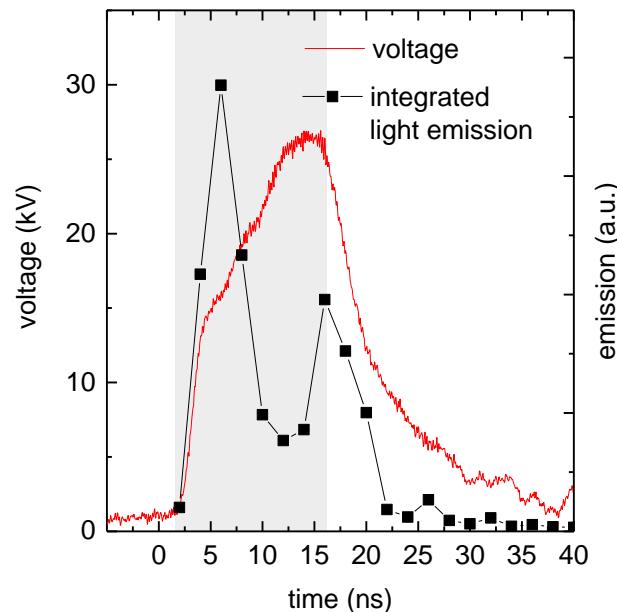
(d) $t = 16 \text{ ns}$

thermionic,
field emission

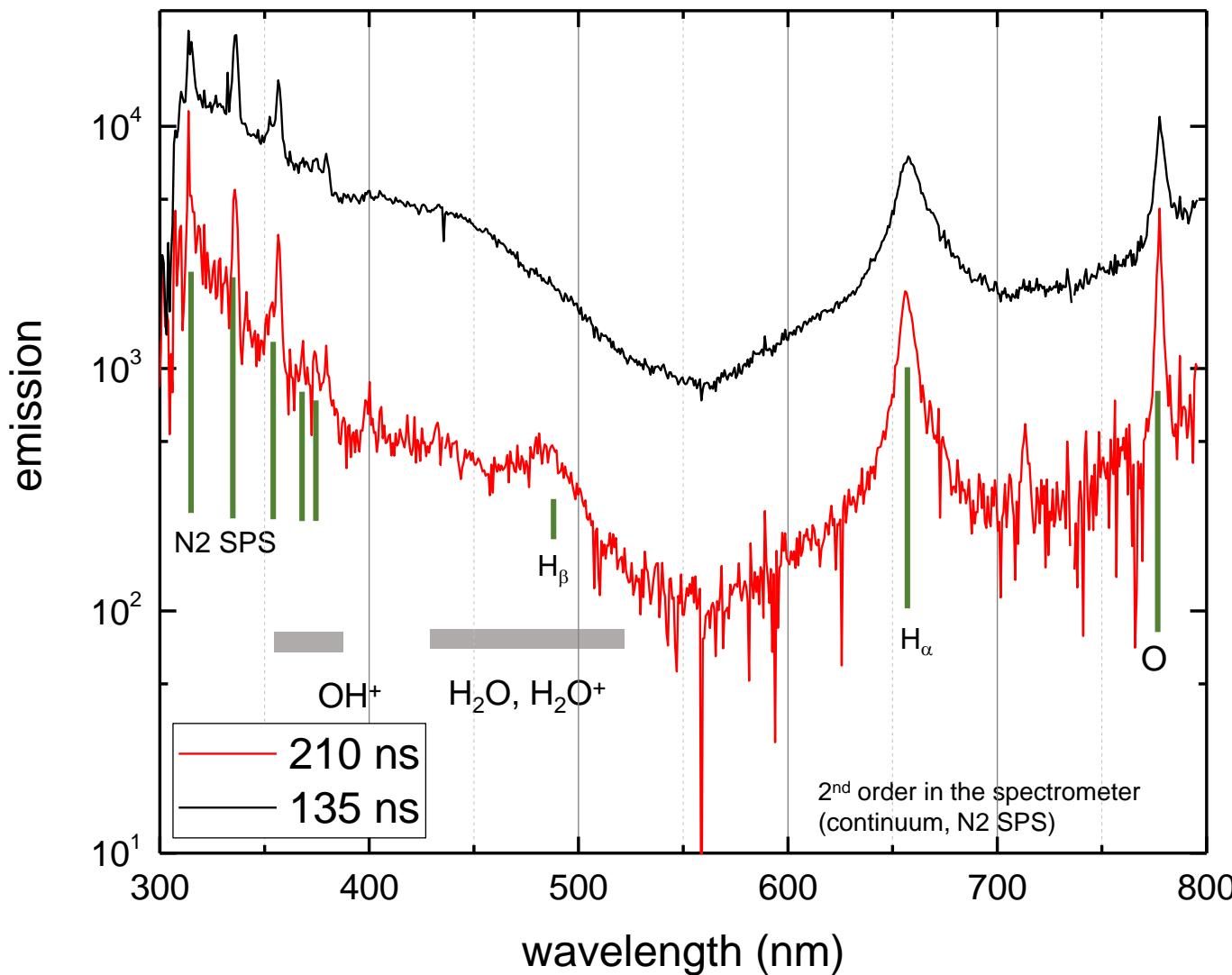


H alpha from ionization region dominates when electrons are being accelerated

- Strong gradients in the beginning
- Field emission at the end of the pulse

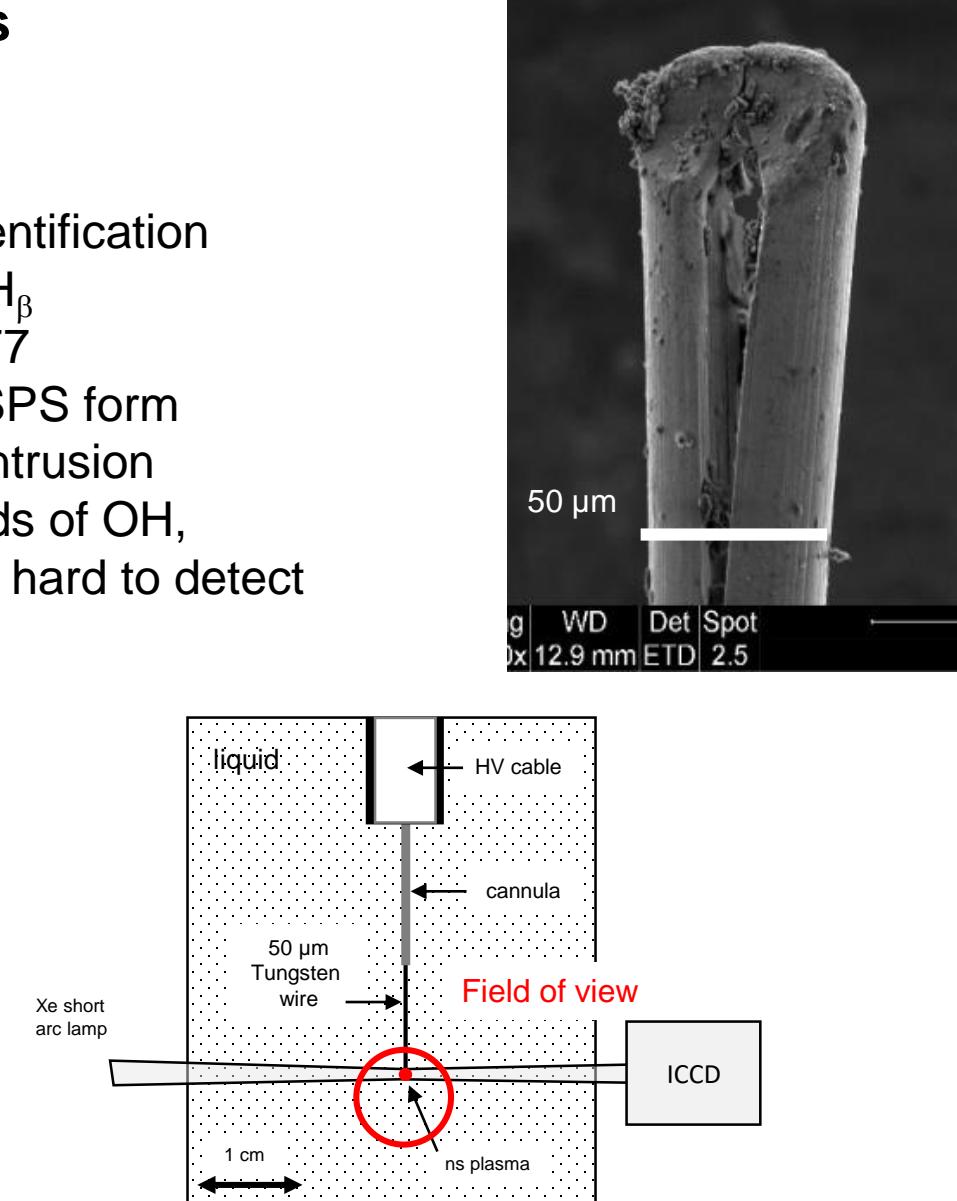


Temperatures from N₂ SPS as thermometer for the later stages



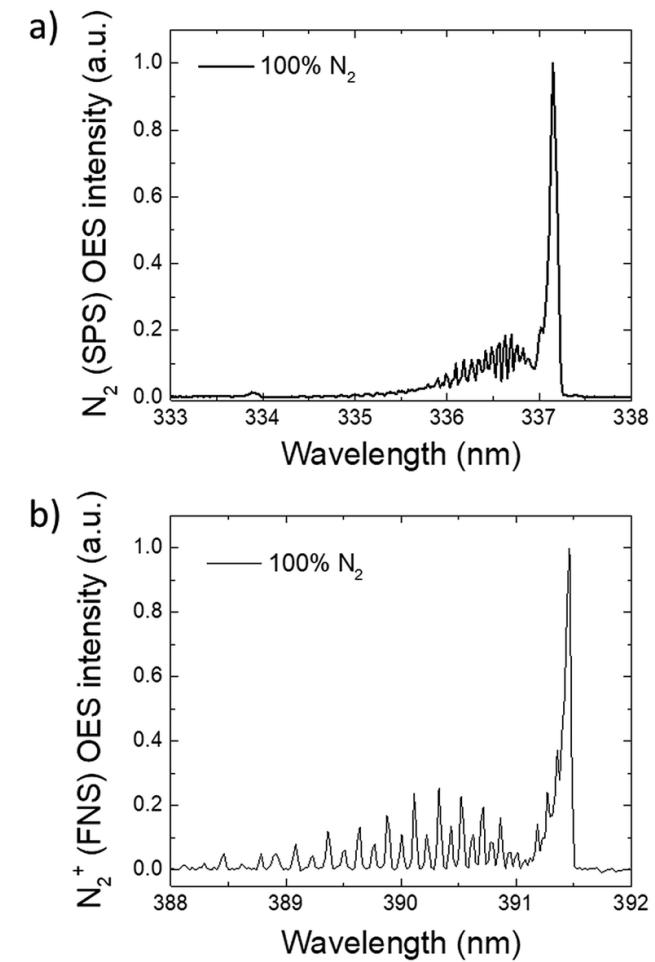
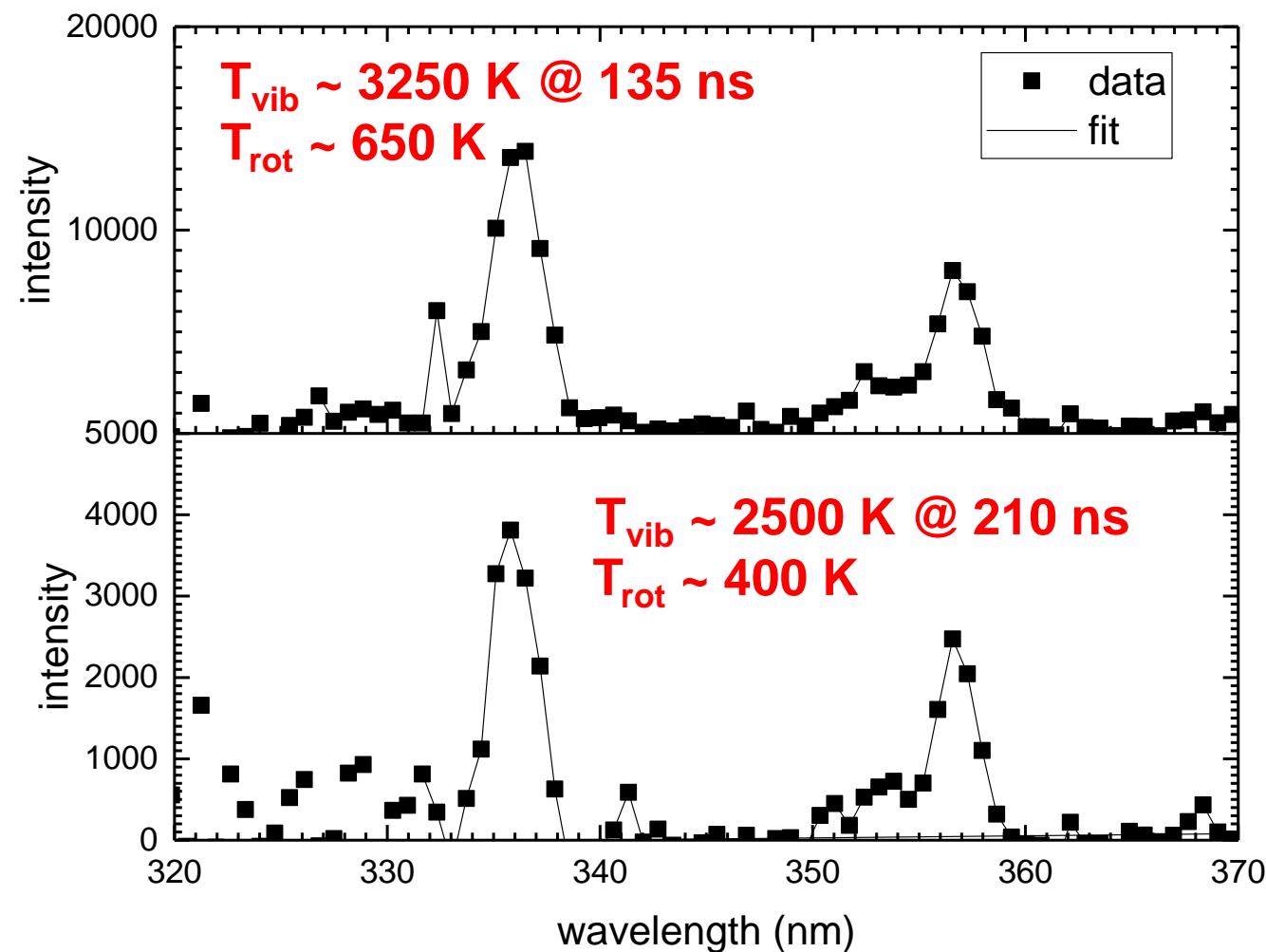
Line identification

- H_α, H_β
- O 777
- N₂ SPS form
- N₂ intrusion
- Bands of OH,
H₂O hard to detect



Temperatures from N₂ SPS as thermometer for the later stages

For 20 kV Experiment



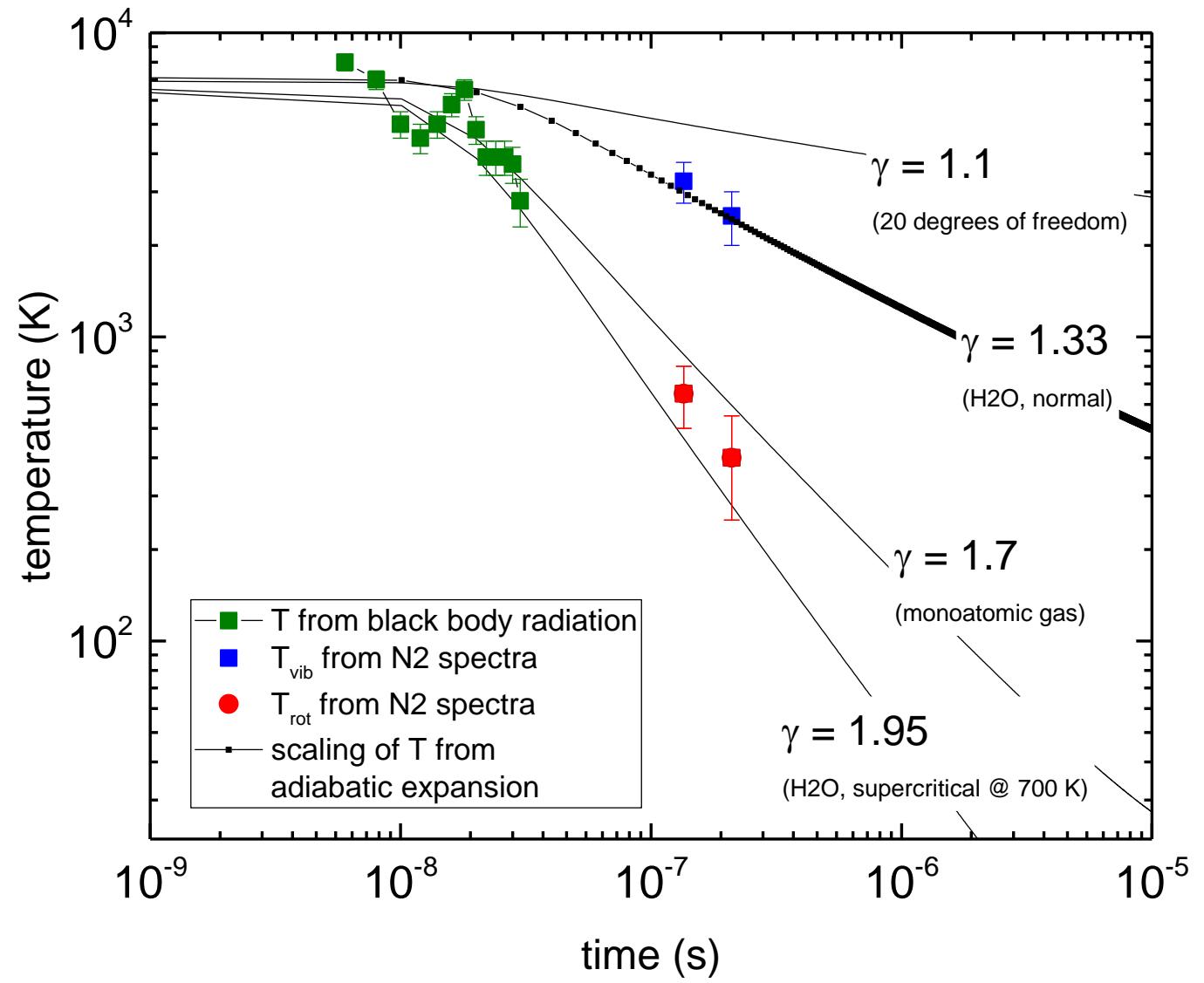
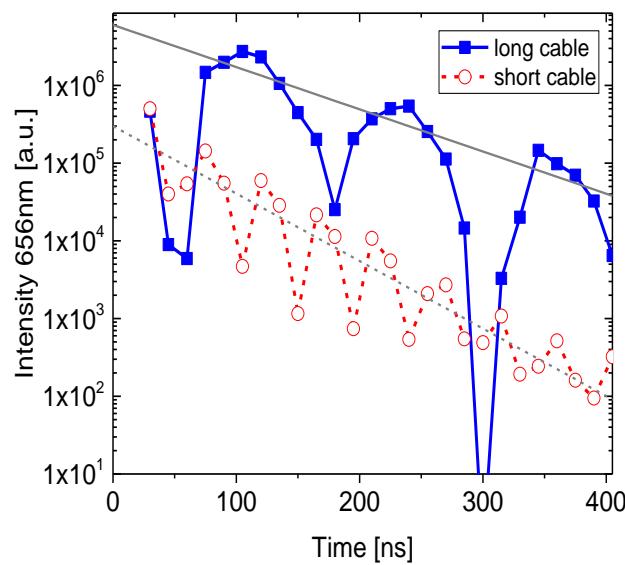
From Girard et al. PCCP 20, 2198 (2018)

Temperatures from spectroscopies

T from cavitation theory

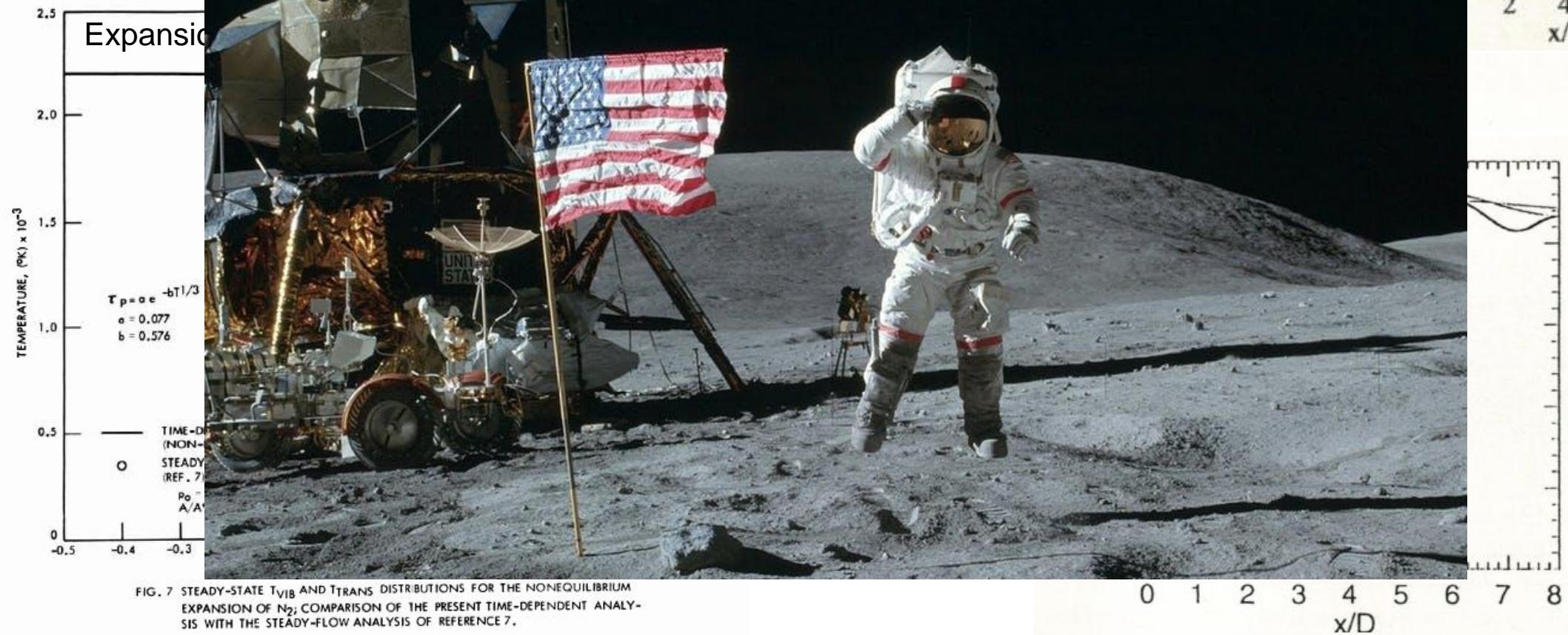
Cooling during expansion depends on
VV, VT, VR relaxation times.

Usually T_{rot} decouples from T_{vib}
In chemical non-equilibrium
expanding flows



Non-equilibrium in expanding ns plasmas

Rocket Science –
Chemical Non-Equilibrium
Hypersonic Flow



J. D. Anderson, Report Naval Ordonnance Labatory (1969)

M. Nichida, M. Matsumoto, Verlag d. Zeitung f. Naturforschung 0932-0784 (1997)

Solution for R(t) assuming a compressible liquid

$$r \frac{Dh}{Dt} + ru \frac{Du}{Dt} + (c + u) \left(h + \frac{u^2}{2} \right) + rc \frac{\partial h}{\partial r} + rcu \frac{\partial u}{\partial r} = 0$$

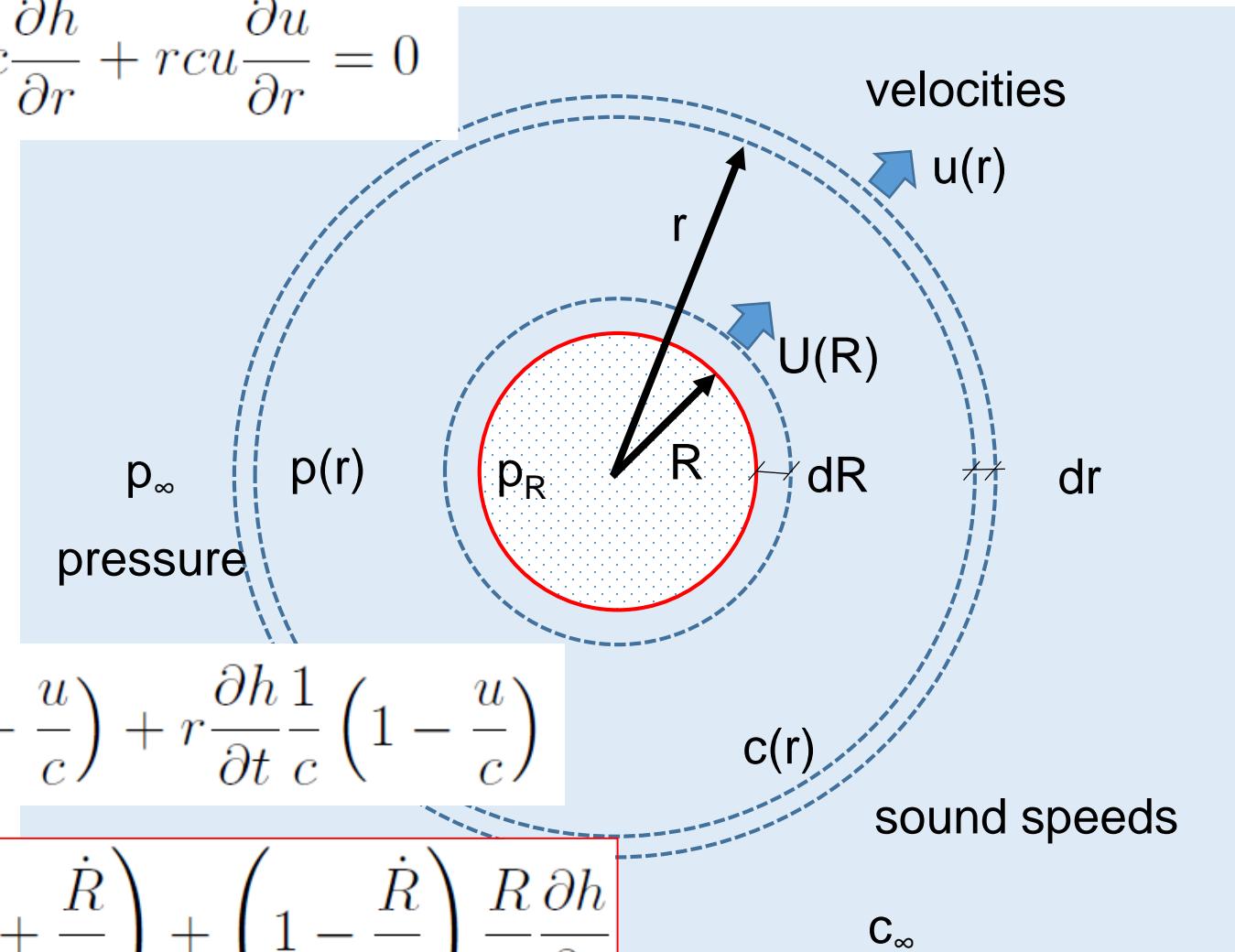
with momentum and mass conservation

$$\frac{Du}{Dt} = -\frac{\partial h}{\partial r}$$

$$\frac{1}{r^2} \frac{\partial}{\partial r} (r^2 u) = -\frac{1}{\rho} \frac{d\rho}{dt}$$

$$r \frac{Du}{Dt} \left(1 - \frac{u}{c} \right) + \frac{3}{2} u^2 \left(1 - \frac{u}{3c} \right) = h \left(1 + \frac{u}{c} \right) + r \frac{\partial h}{\partial t} \frac{1}{c} \left(1 - \frac{u}{c} \right)$$

$$R \ddot{R} \left(1 - \frac{\dot{R}}{c} \right) + \frac{3}{2} \dot{R}^2 \left(1 - \frac{\dot{R}}{3c} \right) = h \left(1 + \frac{\dot{R}}{c} \right) + \left(1 - \frac{\dot{R}}{c} \right) \frac{R}{c} \frac{\partial h}{\partial t}$$



Solution for R(t) assuming a compressible liquid

momentum balance equation (no Navier Stokes)

$$\frac{\partial}{\partial t} \left(-\vec{\nabla} \phi \right) + (\vec{u} \cdot \vec{\nabla}) \vec{u} = -\frac{\nabla p}{\rho}$$

Integrate over r

$$-\frac{\partial \phi}{\partial t} + \frac{1}{2} \vec{u}^2 = \int_{p_\infty}^{p(r)} -\frac{dp}{\rho}$$

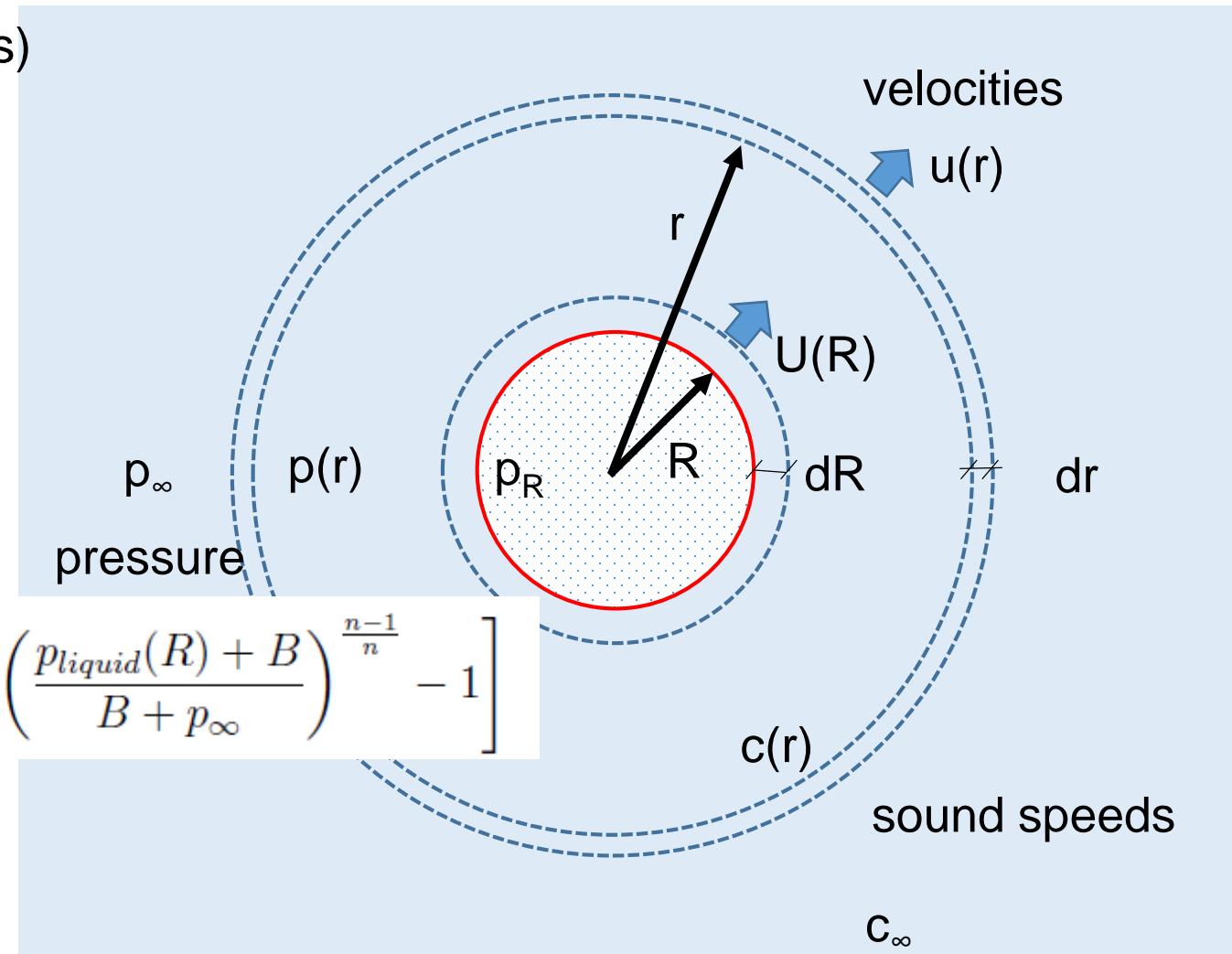
definition enthalpy

$$h = \int_{p_\infty}^{p(r)} -\frac{dp}{\rho}$$

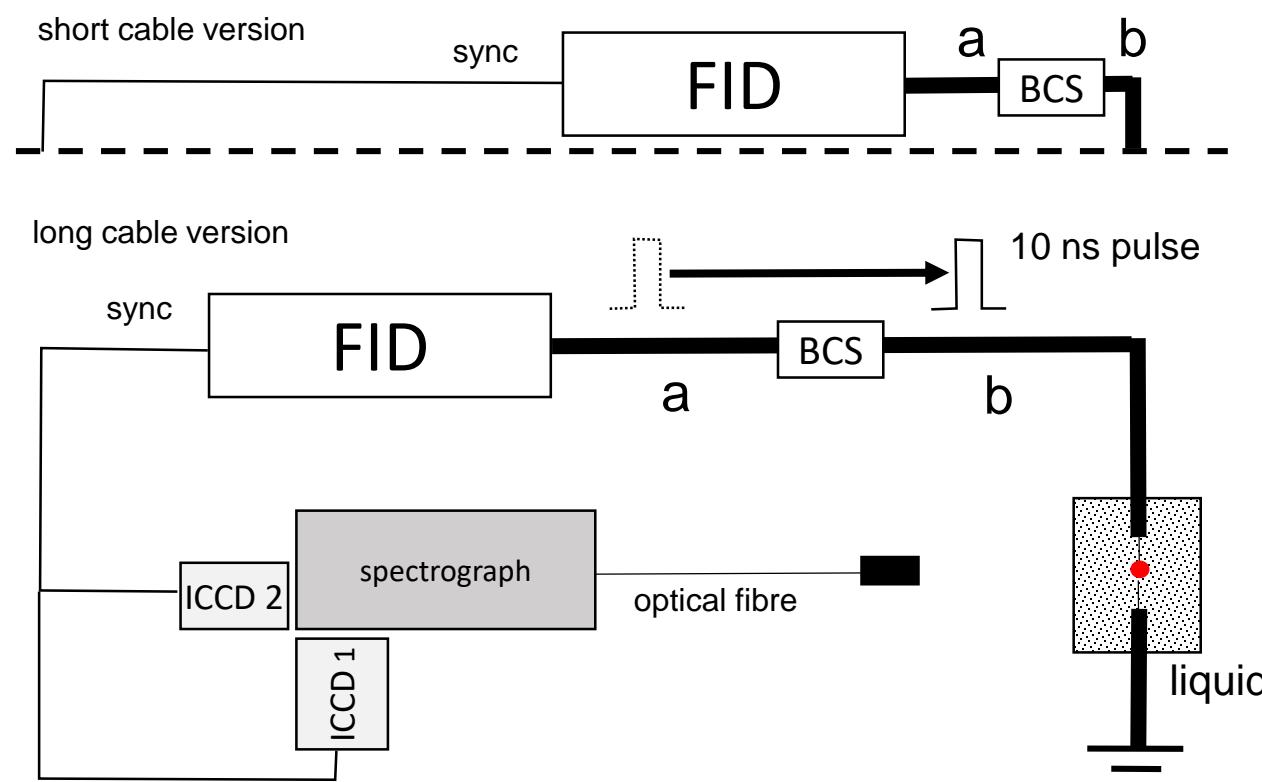
$$h = \frac{n}{n-1} \frac{1}{\rho_0} (B + p_\infty) \left[\left(\frac{p_{liquid}(R) + B}{B + p_\infty} \right)^{\frac{n-1}{n}} - 1 \right]$$

results in Bernoulli equation

$$\frac{\partial \phi}{\partial t} = h + \frac{1}{2} \vec{u}^2$$



Example 2: Plasma Chemistry in liquids – example H₂O₂ generation

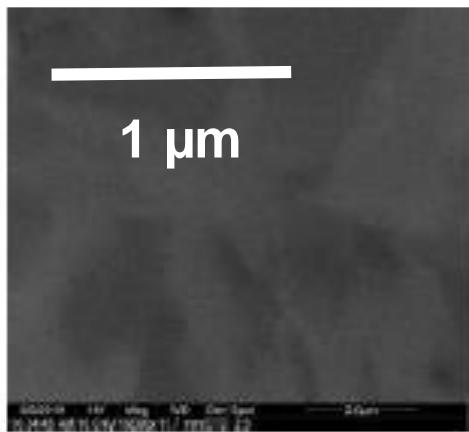


- 10 ns pulses, 14...30 kV, distilled water
- Shadowgraphy 2 ns ... 70 ns gate widths
- Optical emission spectroscopy with 2ns and 30 ns gate time

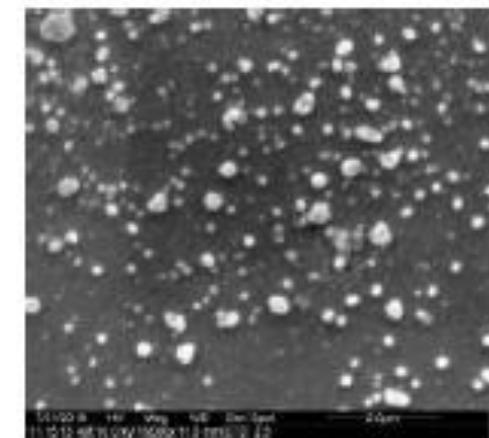
Nanosecond plasma based recovery of CuO nanocubes at copper electrodes

20 kV

distilled water

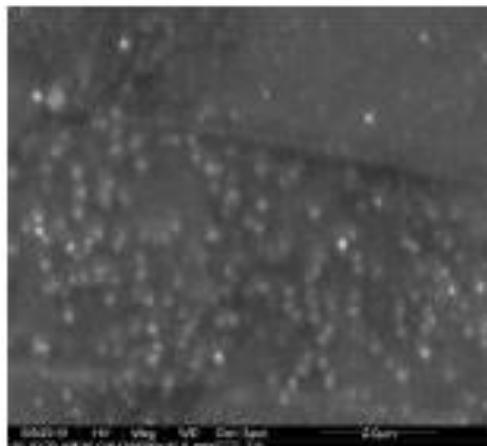


dist. Water + 0.0035M KCl

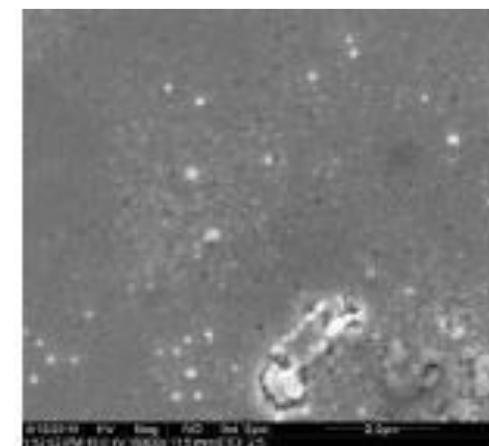


26 kV

distilled water

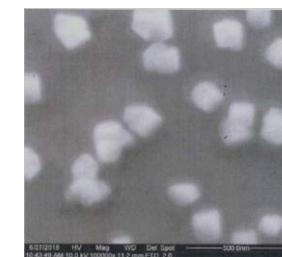


dist. Water + 0.0035M KCl



Creation of Cu-nanocubes possible, if

- no direct current to the sample, reduction of the oxide, plasma electrode distance to the sample important
- Cl necessary as nucleation sites
- Competition between oxidation due to H_2O_2 , OH and reduction by H, e^-



P. Grosse, B. Roldan et al.

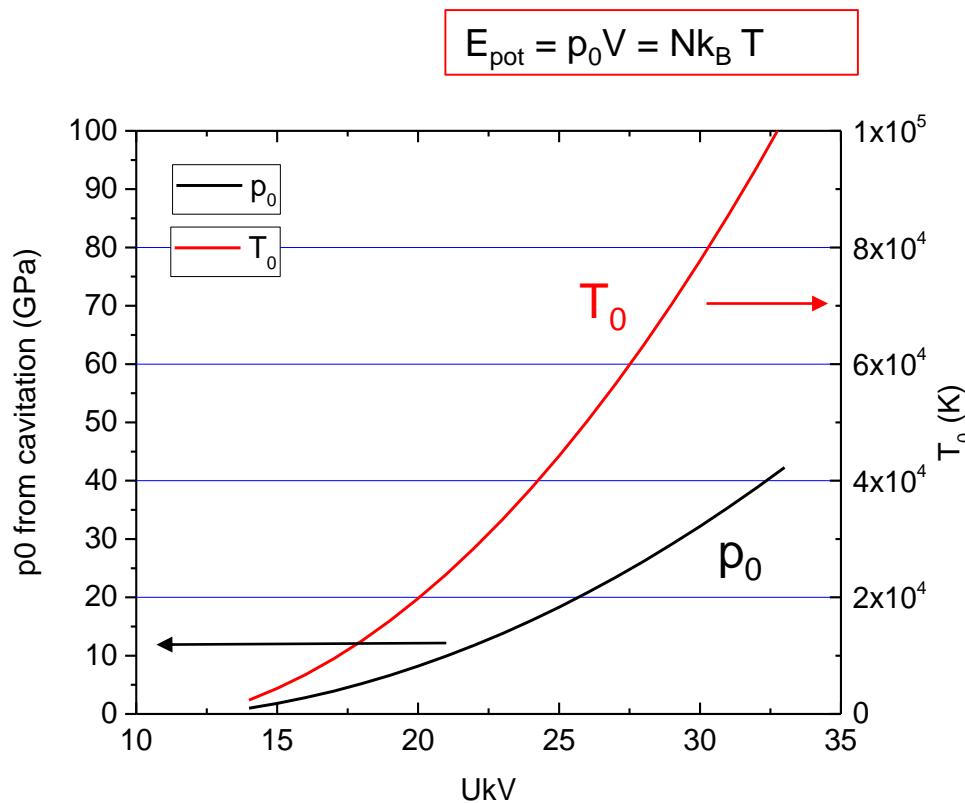
Reference plasma based low pressure creation of CuO nanocubes

Efficiency to create H₂O₂ with these plasmas

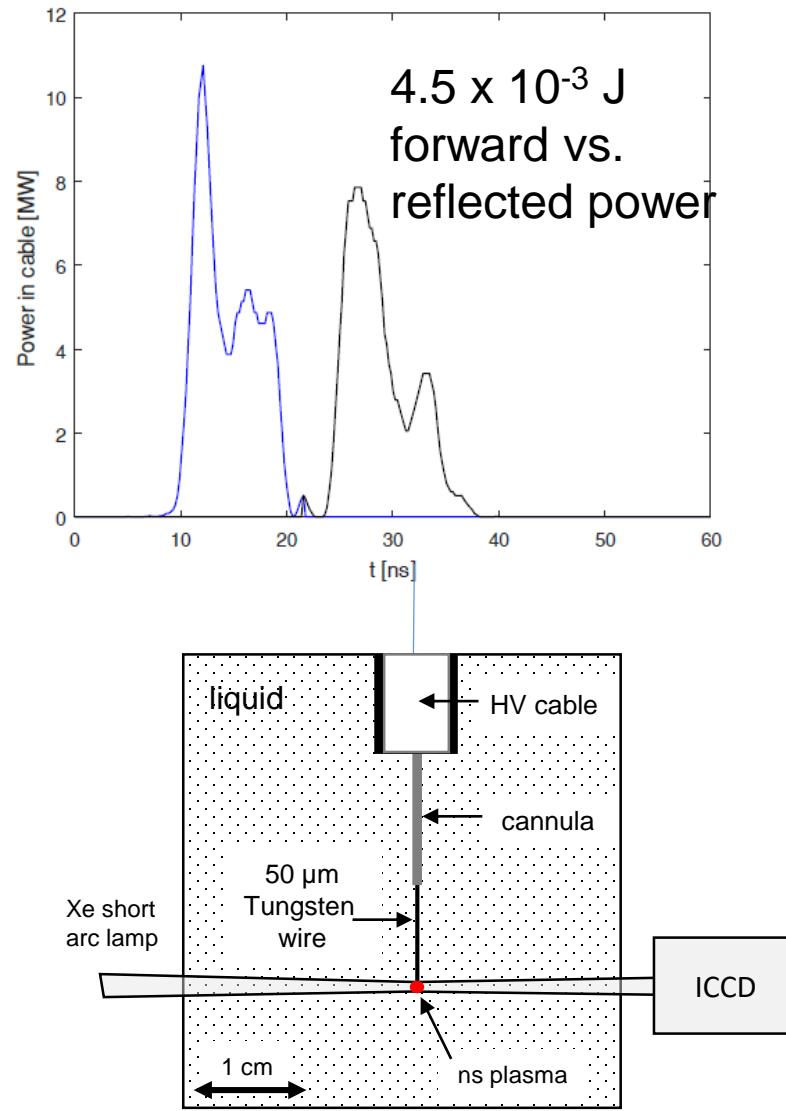
Modeling the chemistry

*S. Mededovic, B. Locke
JPD 40, 7734 (2007)
PCPP 32, 875 (2012)

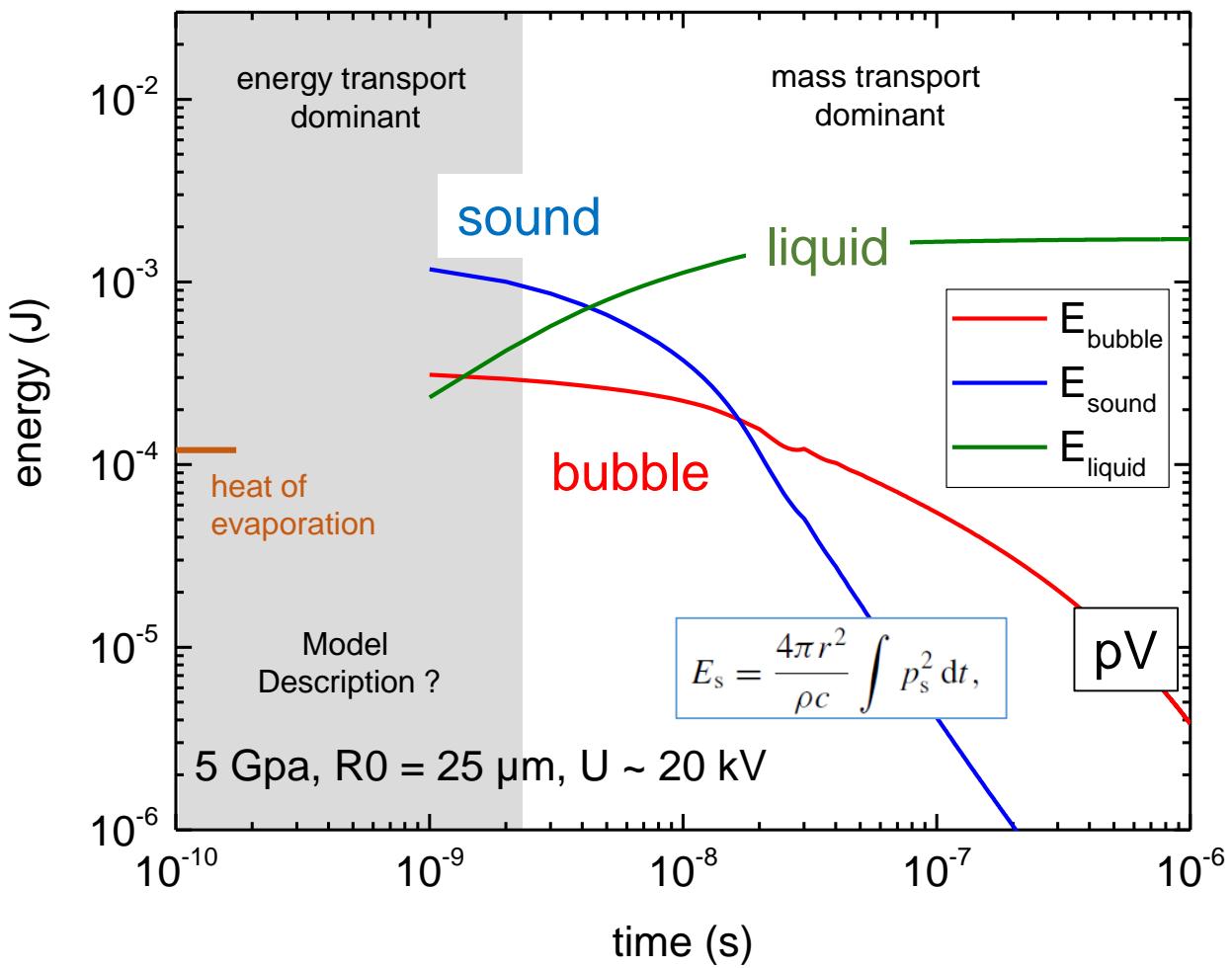
- Taking H₂O chemistry from spark plasma simulation*
- Taking the evolution of the temperature from cavitation theory



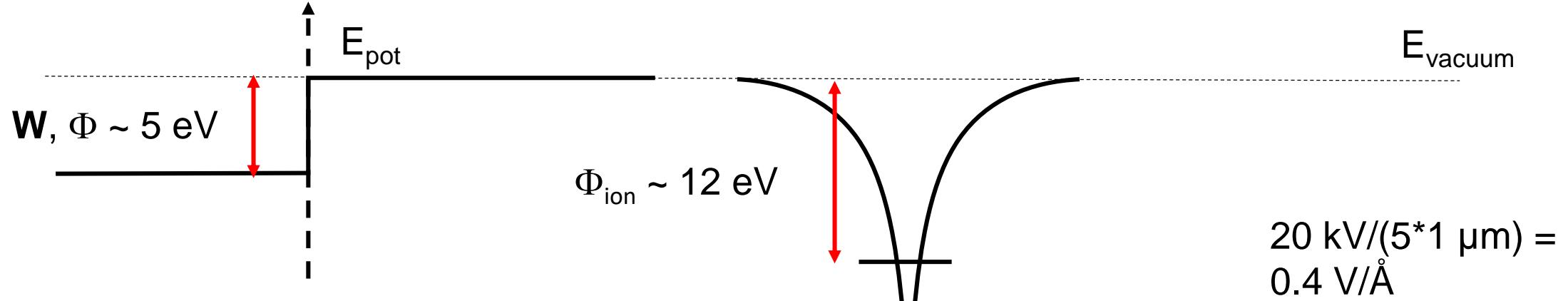
Dissipated energy during plasma ignition within the first 10 ns – comparison to experiments



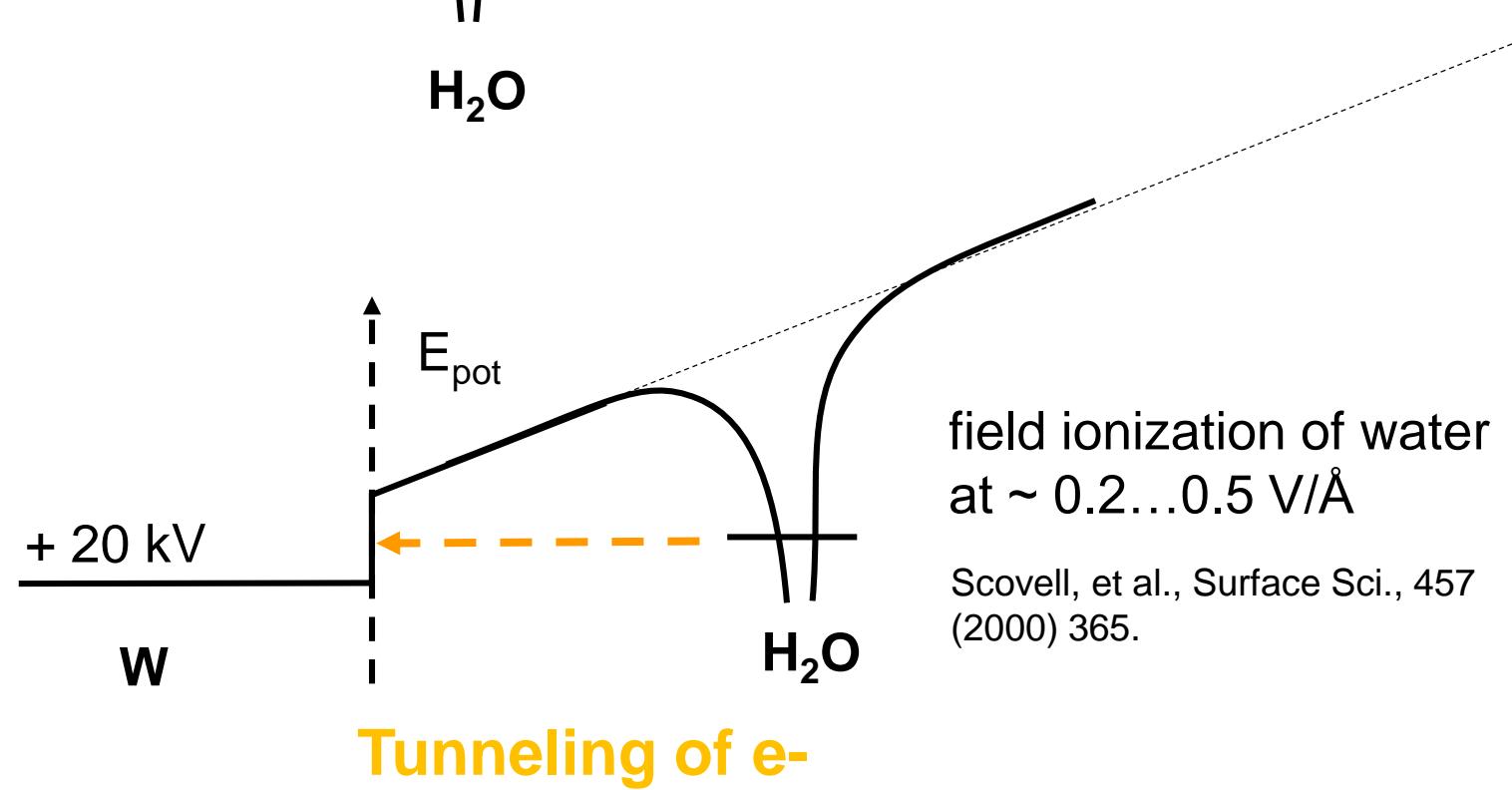
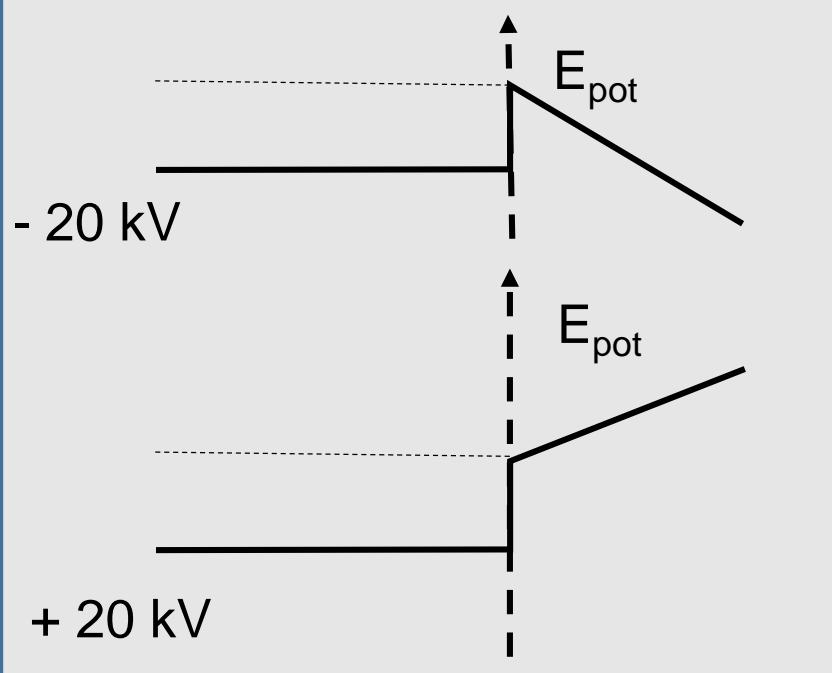
In ns plasmas at very high initial pressure, most of the energy is transferred to the sound wave



Field emission effects



Field effects at surfaces



field ionization of water
at $\sim 0.2 \dots 0.5 \text{ V}/\text{\AA}$
Scovell, et al., Surface Sci., 457
(2000) 365.

Streamer velocities in distilled water

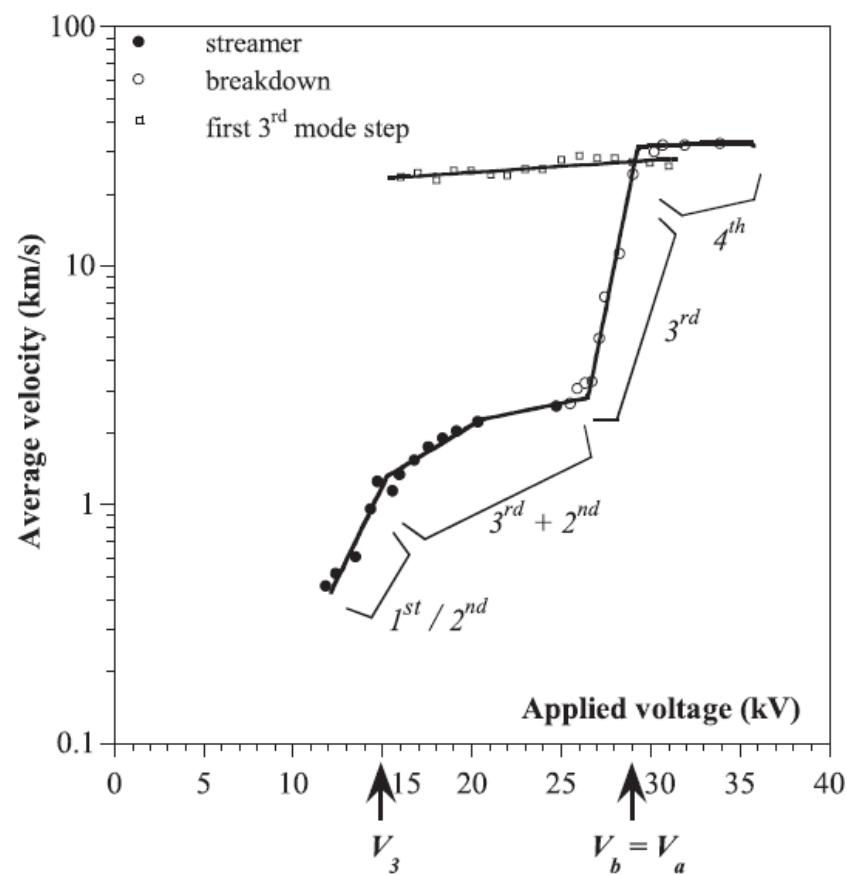
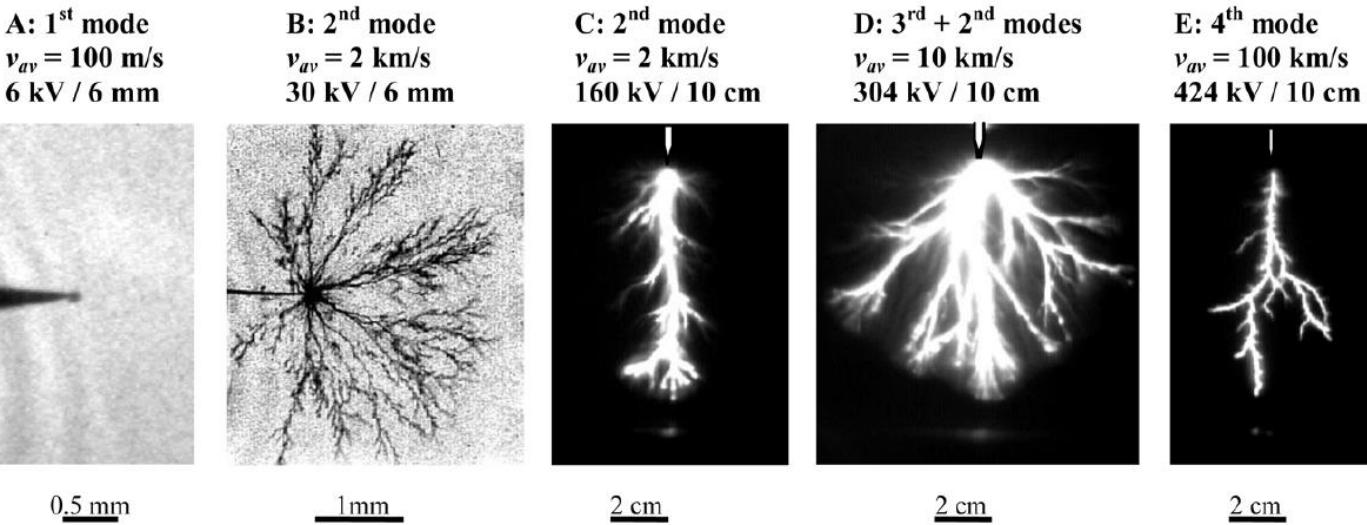


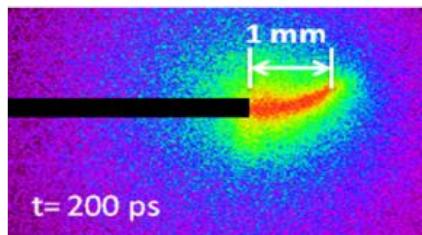
Figure 15. Average streamer propagation velocity v_{av} versus voltage in distilled water [52]. $d = 3\text{ cm}$, $r_p = 1\text{ }\mu\text{m}$.

Lesaint et al. JPD 49 144001 (2016)

Starikovsky et al. PSST 20, 24003 (2011)

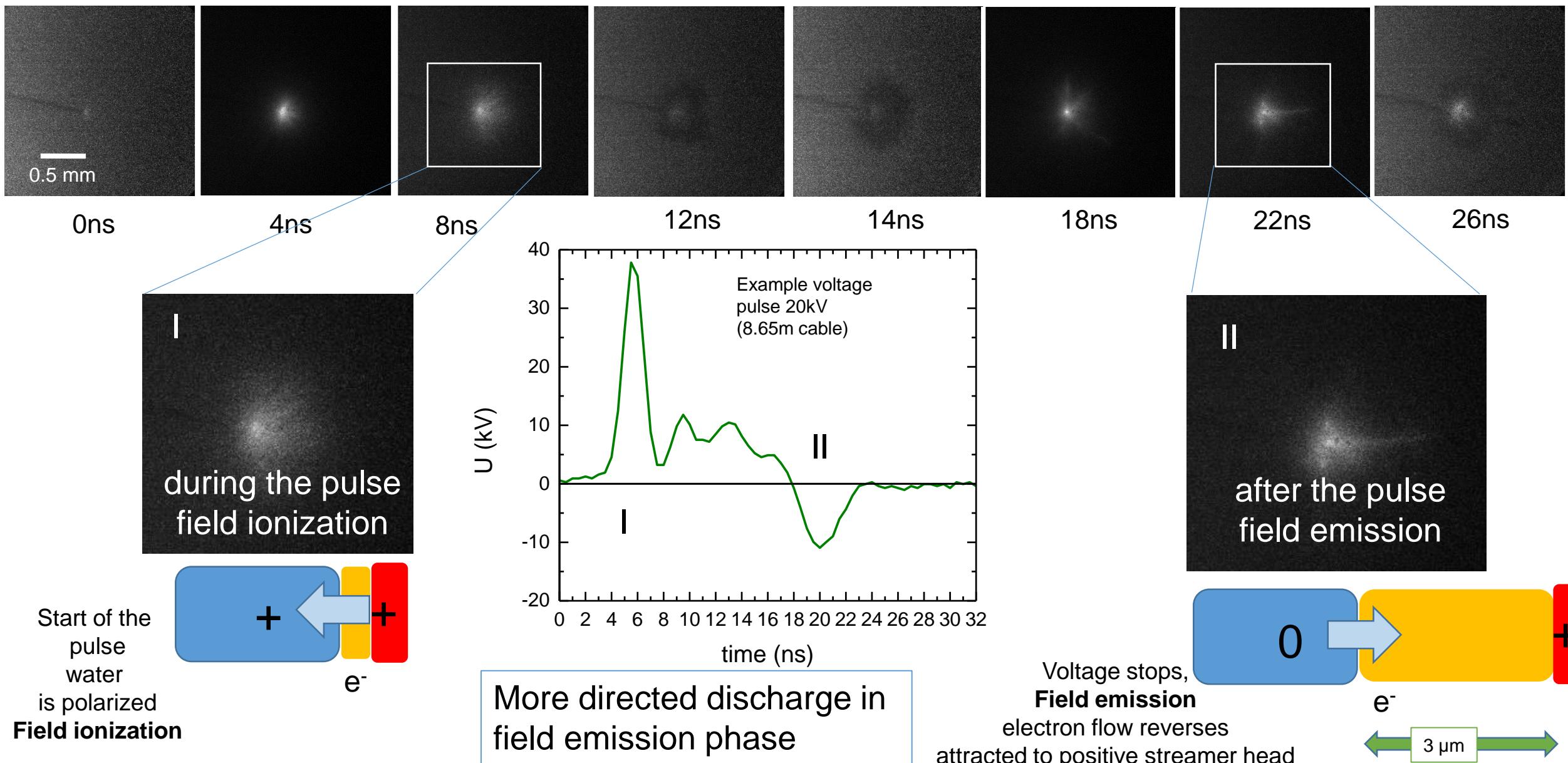


Propagation velocity of streamer in liquids (long pulses)
3 km/s ... 100 km/s in 10 ns $\Delta x = 3\text{ }\mu\text{m} \dots 0.1\text{ mm}$
Propagation velocity of streamer in liquids (short pulses)
1000...5000 km/s in 10 ns $\Delta x = 1\text{ mm} \dots 5\text{ mm}$

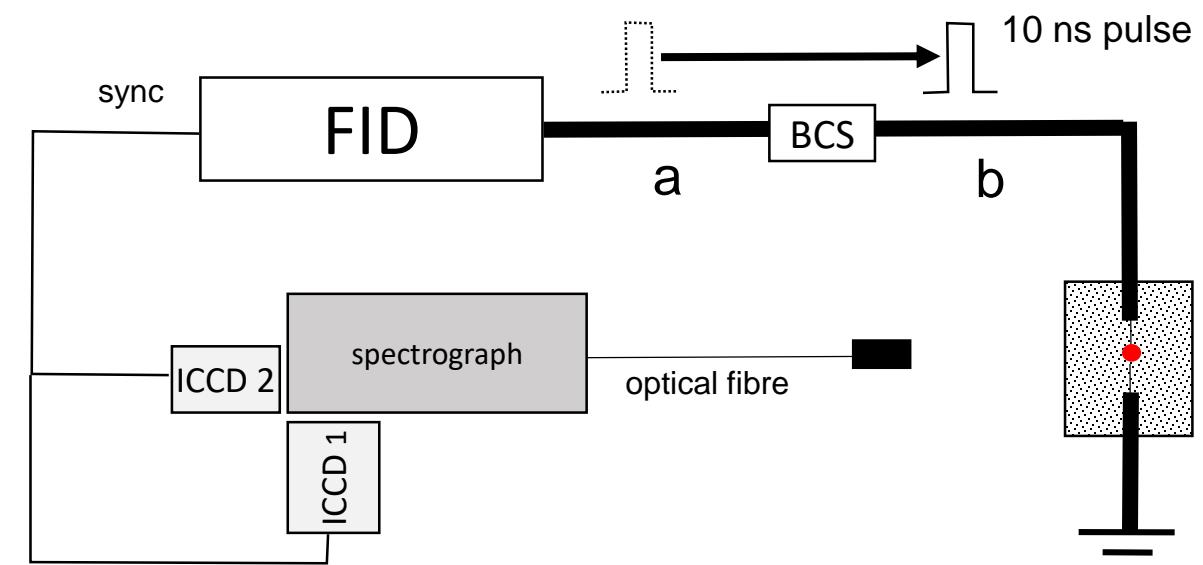
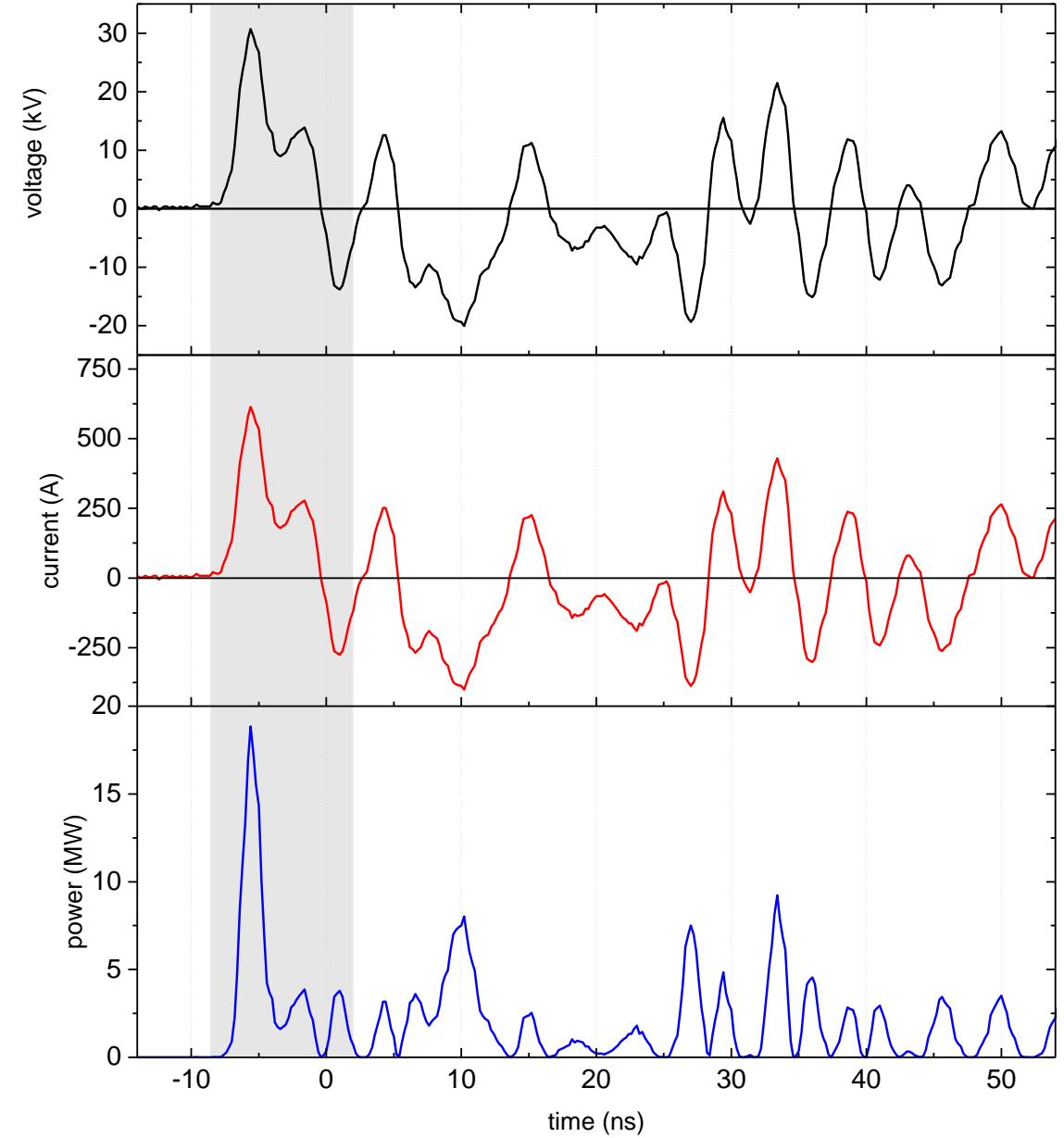


Comparison Streamer
in air $\sim 1000\text{ km/s}$

Plasma pulse from field ionization to field emission (gate time 2 ns, 22 kV)

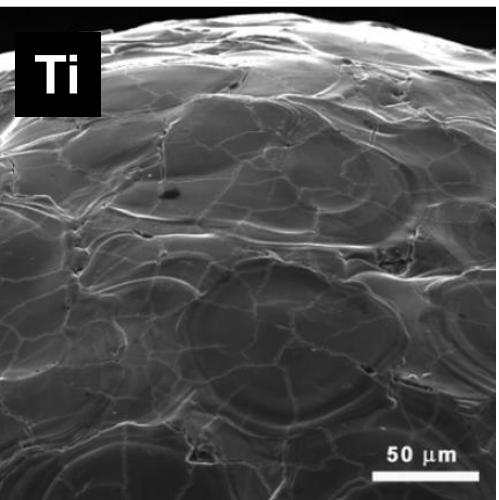


Current and Voltage signals for 10 ns 20 kV pulses using the 8.65 m cable

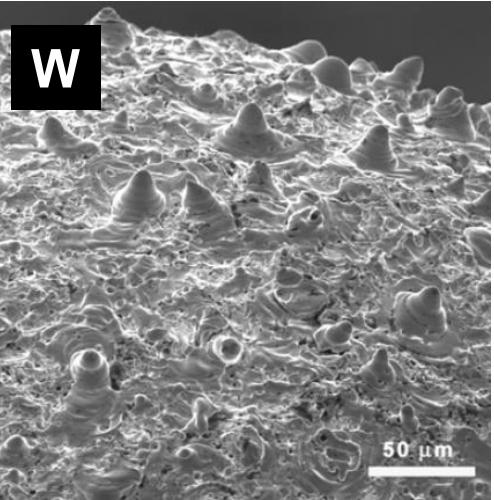


Tungsten vs. Titanium* tip (500 µm tip radius vs. 50 µm tip radius)

Spark 5..10 µs



3 J $1.5 \times 10^7 \text{ J/m}^2$



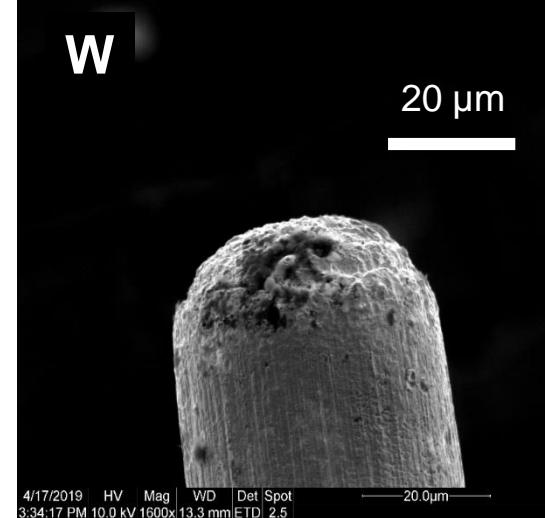
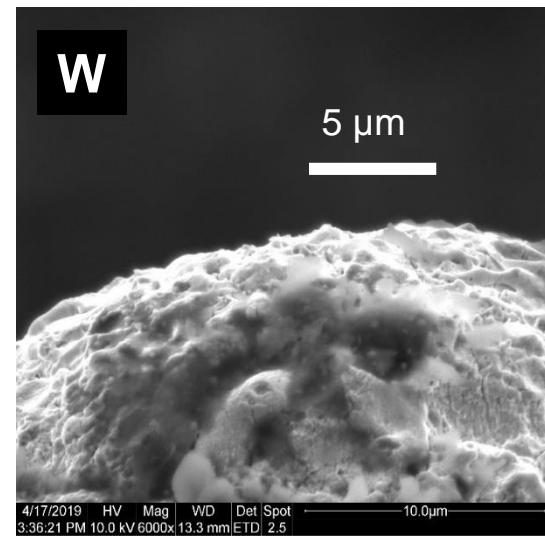
- similar energy per m² dissipated
- Sequence of melting and solidification

thermal conductivity @ 1000 °C
melting temp.

- Ti – 22 W/m K, 2000 K
- W – 115 W/m K, 3700 K
- Pt – 82 W/m K, 2500 K
- In case of W only small melting tips (for Pt, T, larger molten areas)
- similar to cathode spot in lamps

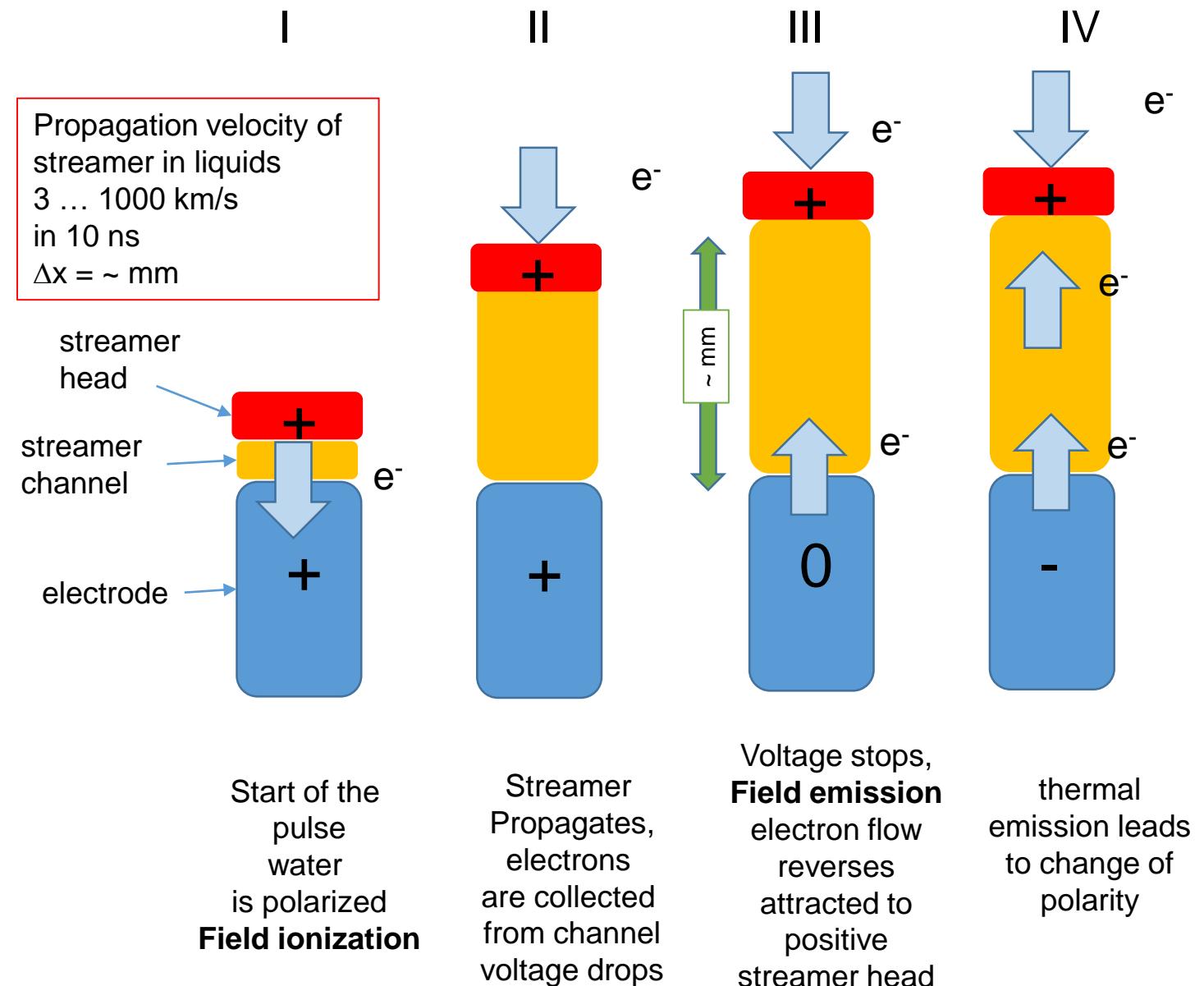
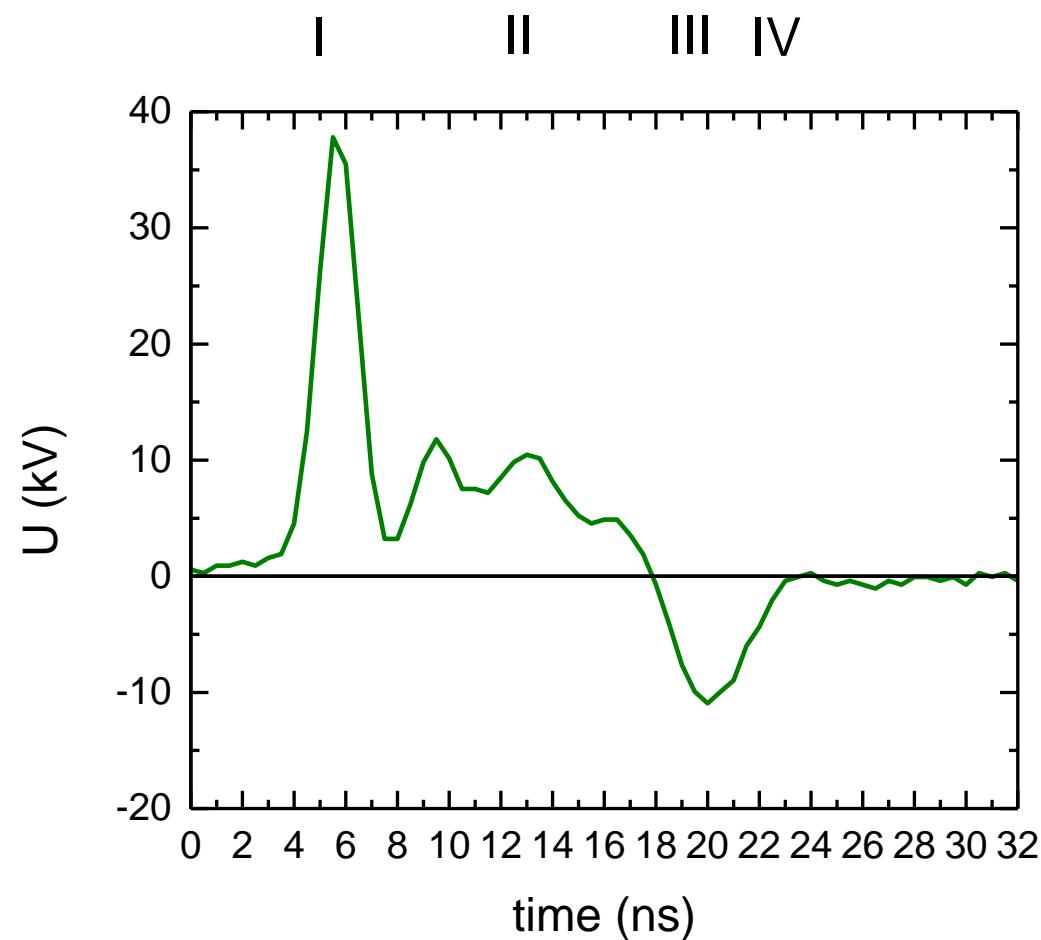
ns pulse 10 ns

0.03 J $1.5 \times 10^7 \text{ J/m}^2$



*P. Lukes et al. PSST 20, 3401 (2011)

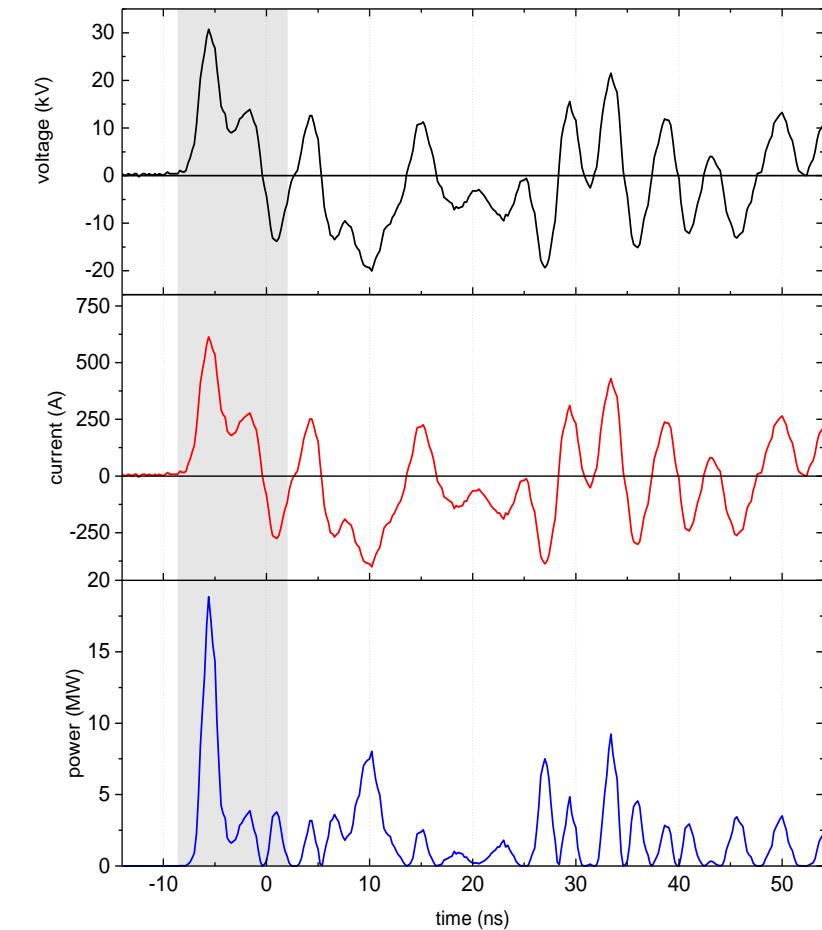
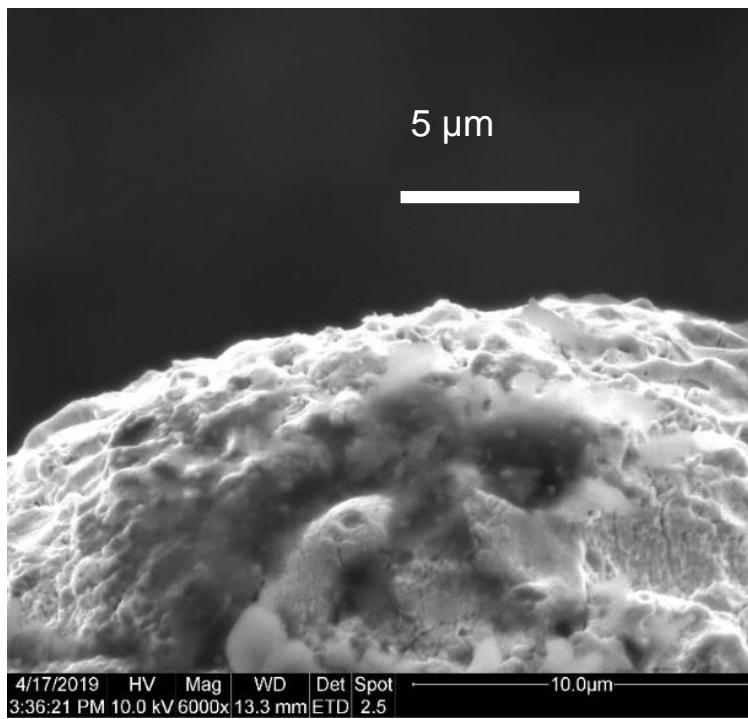
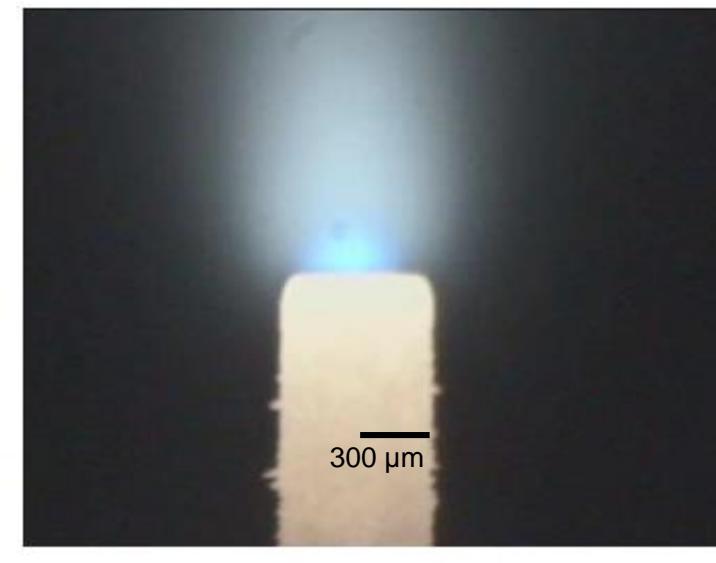
Sequence of the propagating discharge



Comparison to cathode spots at high pressure lamps

@500 A on a 2 μm radius spot
= magnetic field pressure of 1 Gpa

Pinch effects at W tips

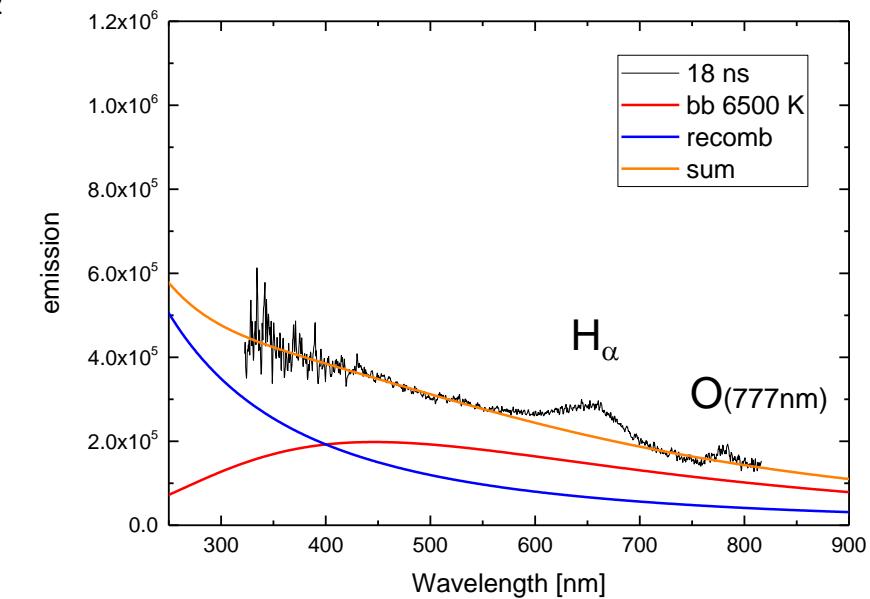
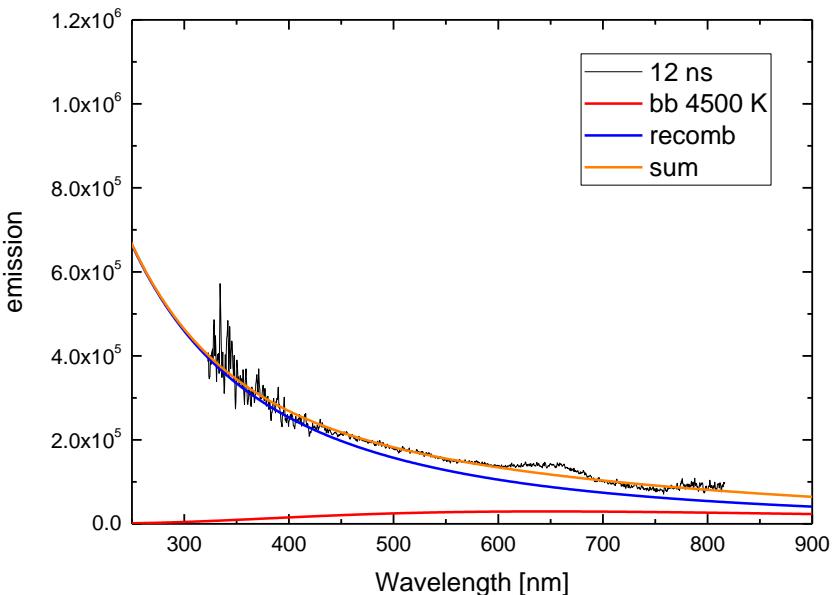
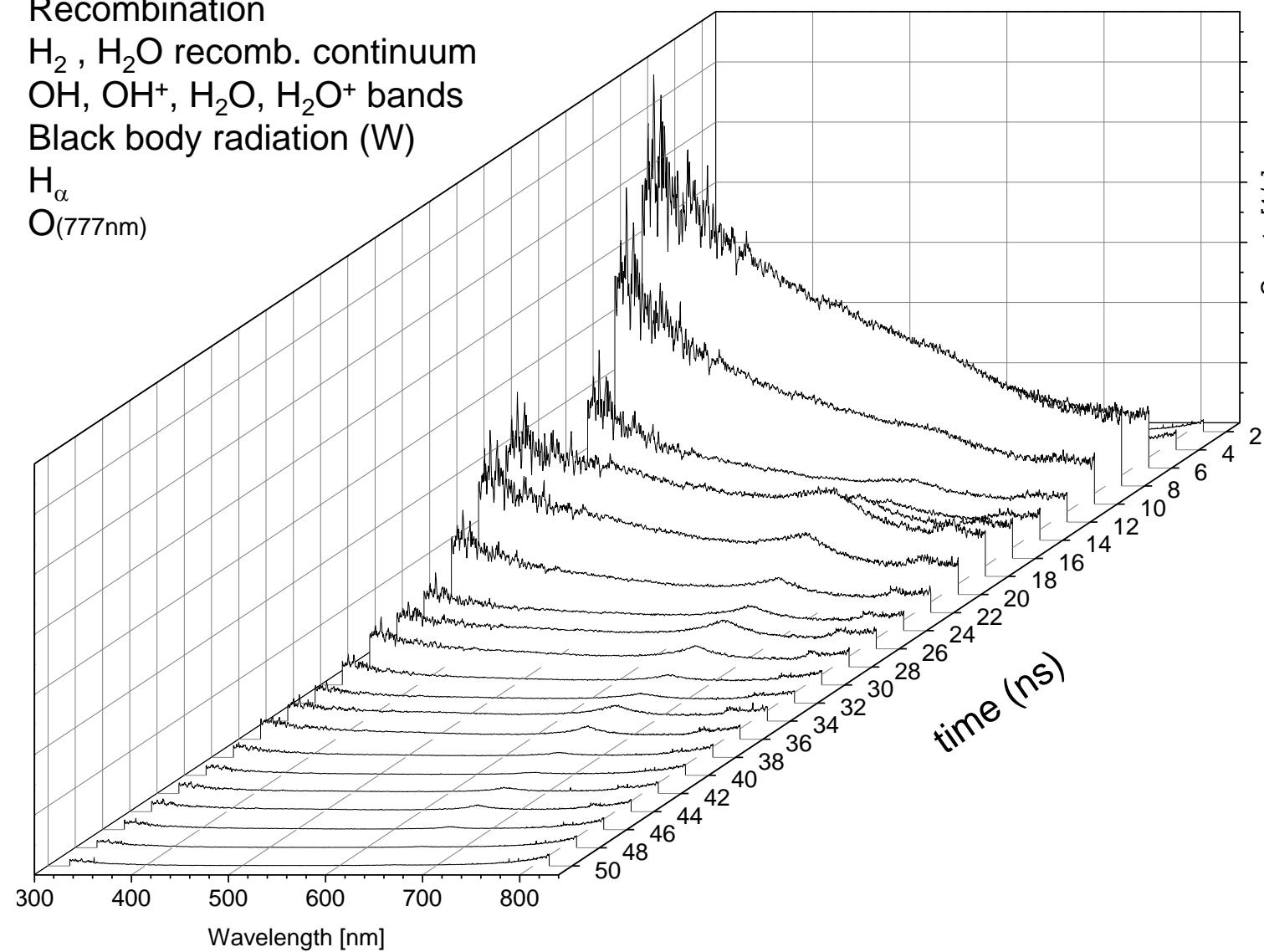


A. Bergner et al. PSST 23, 054005 (2014)

Temperatures from emission spectra of the plasma

Spectrum (first 50 ns, time resolution 2 ns)

- Recombination
- H₂, H₂O recomb. continuum
- OH, OH⁺, H₂O, H₂O⁺ bands
- Black body radiation (W)
- H_α
- O(777nm)

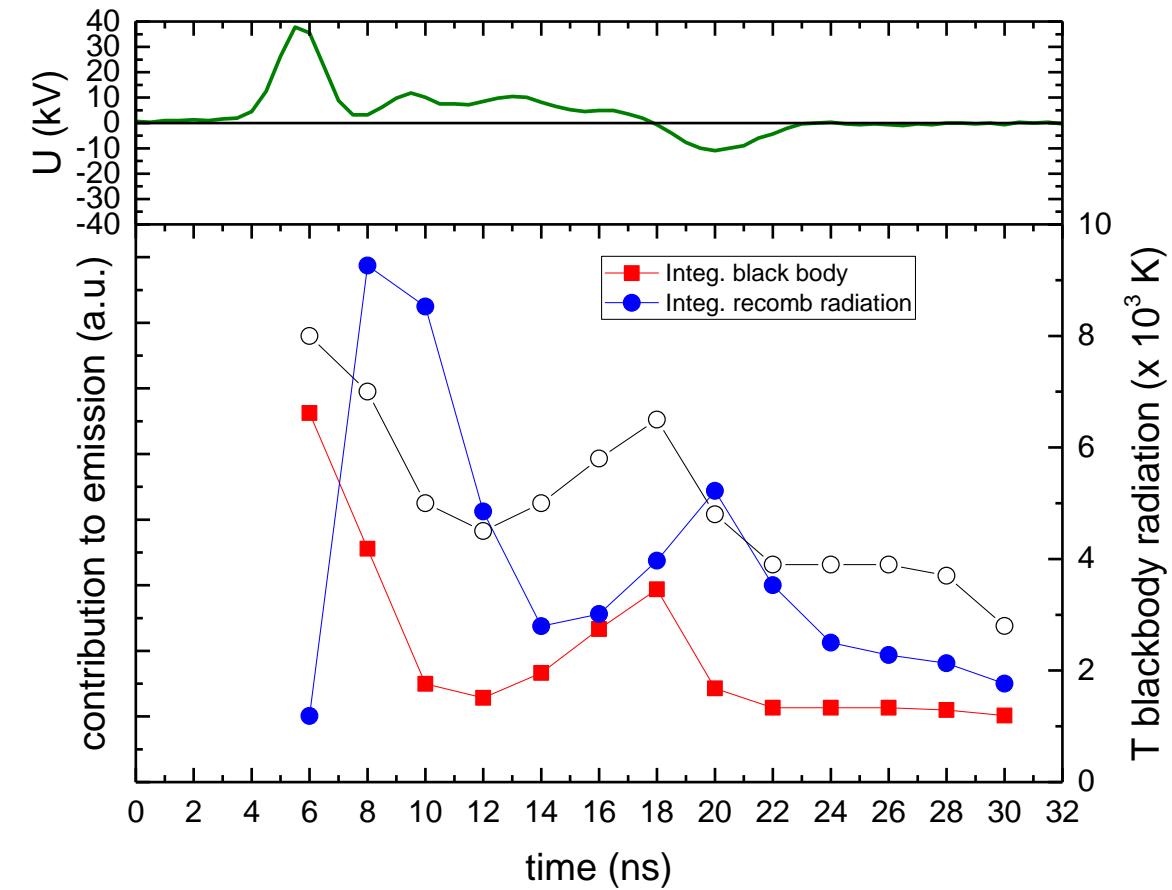
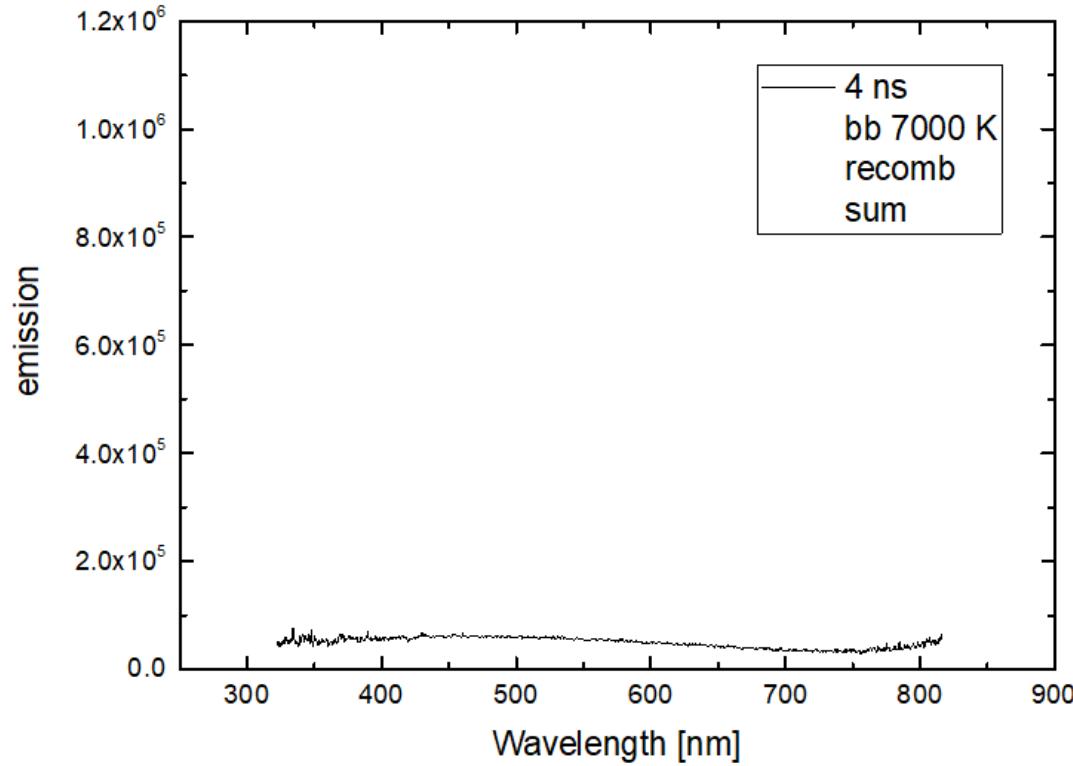


Temperatures from emission spectra of the plasma

For 20 kV Experiment

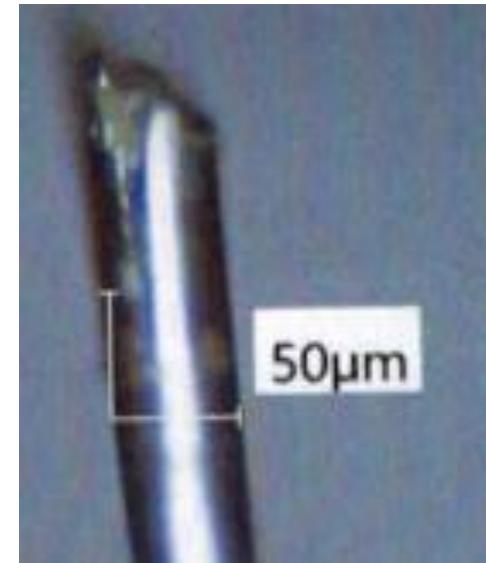
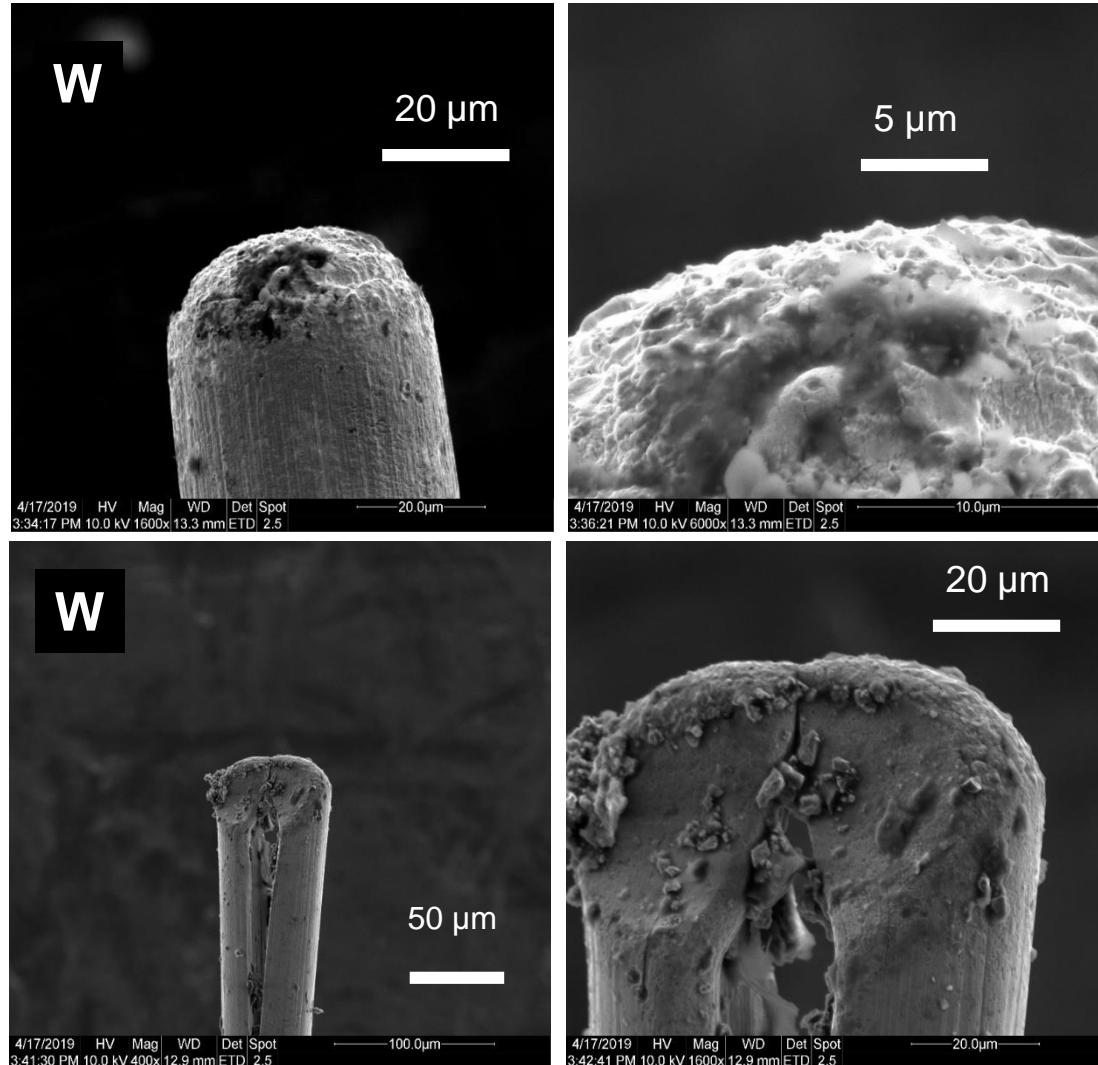
- Power absorption at the beginning of the pulse and due to reversal of the voltage
- Black body radiation, followed by $1/\lambda^x$ spectrum.

$$T_0 < 8000 \text{ K} @ 0 \text{ ns}$$



Tungsten tip

- Long term operation possible
- W as high melting temperature 3700 K
- W therm. conduct. 115 W / m K
- W oxidizes, melting temperature 1800 K
- Crystallites are formed

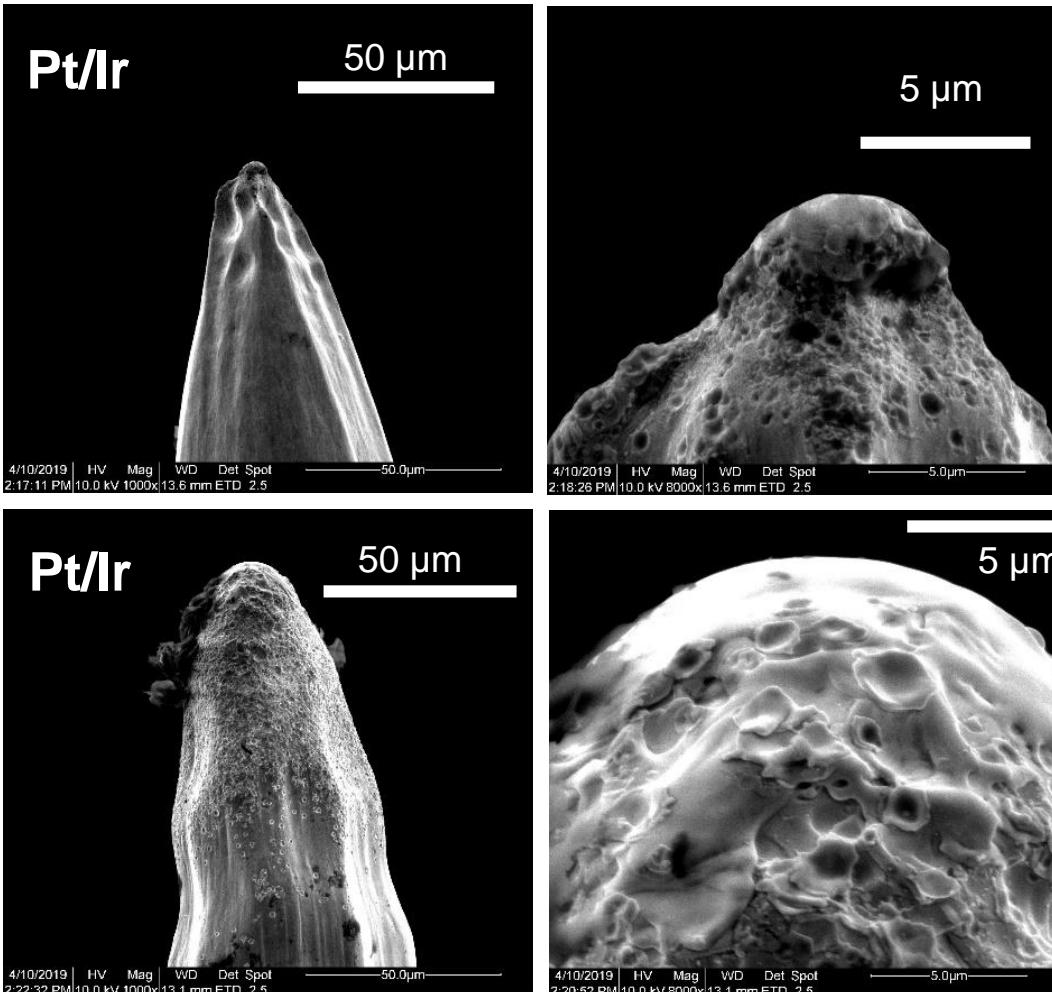


Short term use, < 7 h

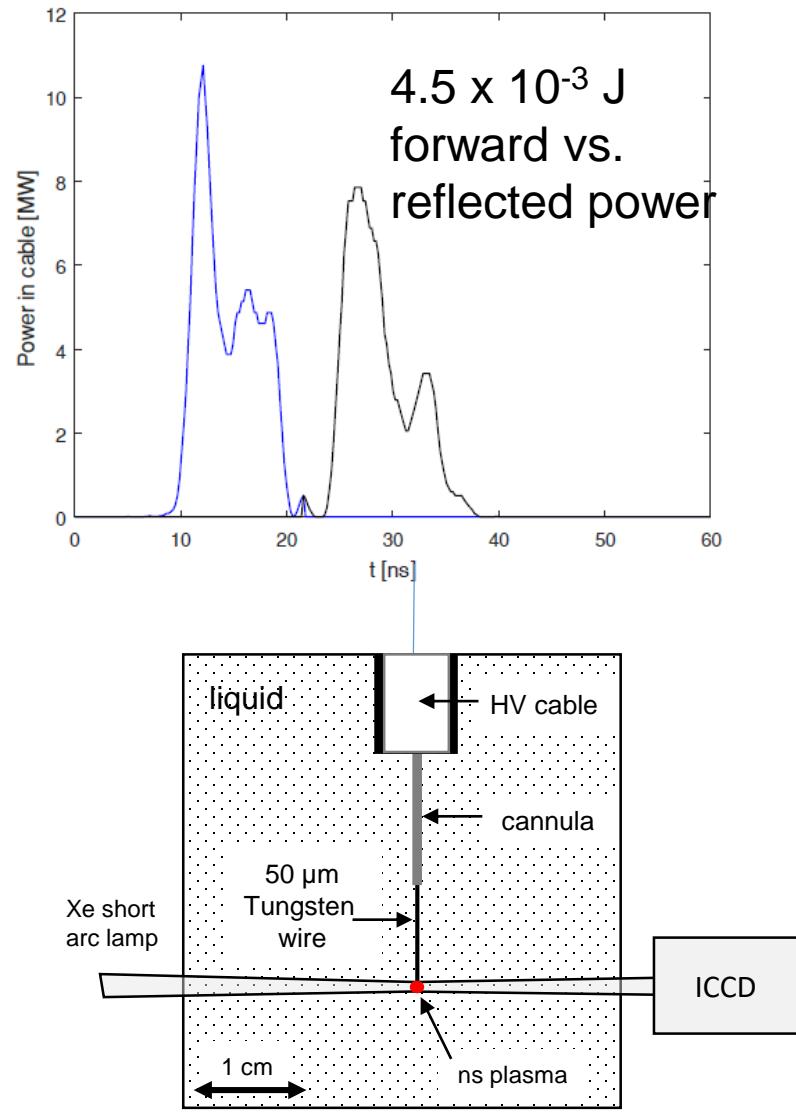
long term use, ~ 30 h

Platinum / Iridium tip (from AFM/SEM microscopes with 5..10 nm tip radius)

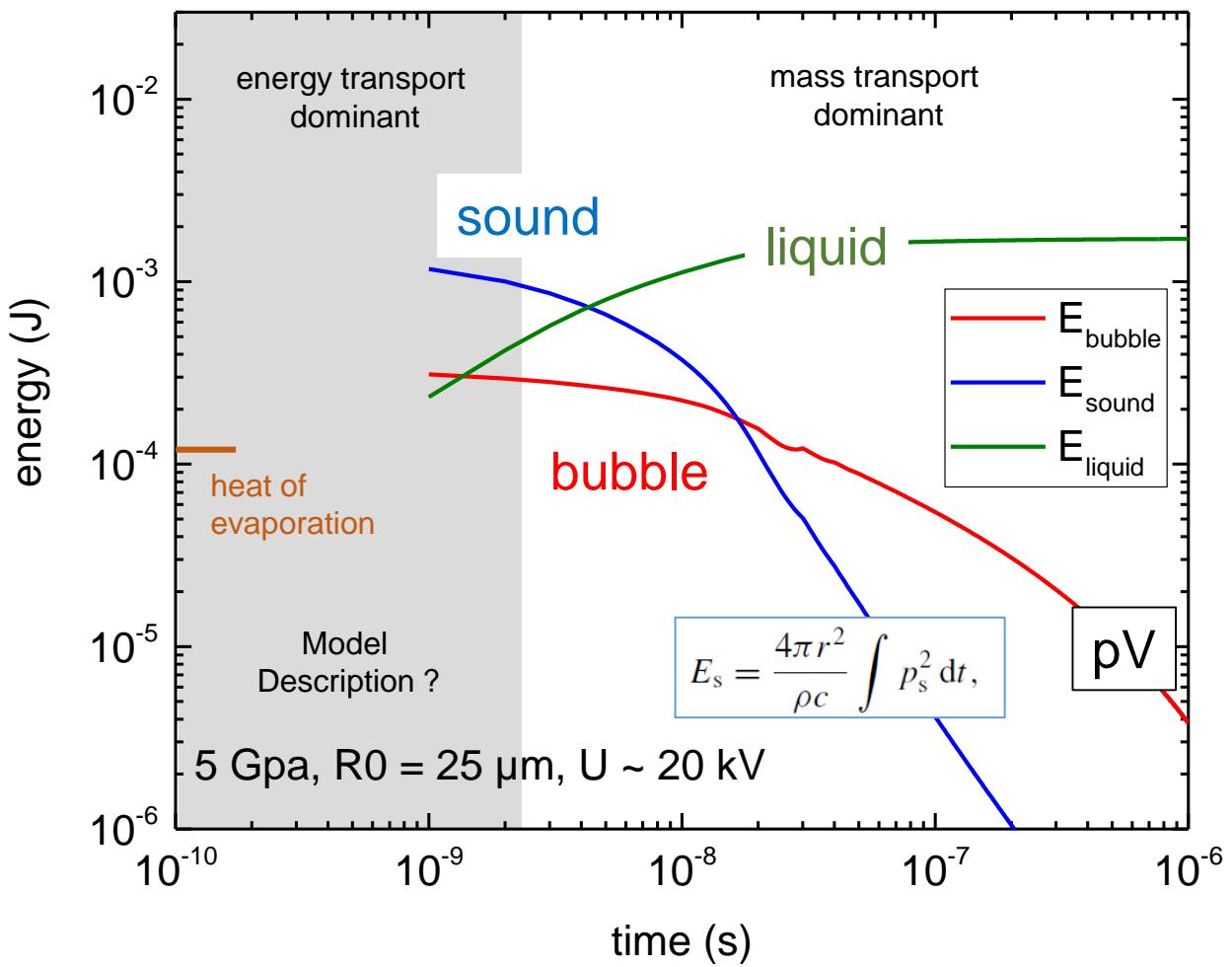
- Only short term operation possible
- After 1 h no plasma ignition possible anymore
- Pt melting temperature 2100 K
- Pt thermal conduct. 82 W/ m K
- Ir melting temperature 2700 K
- only Ir oxidizes

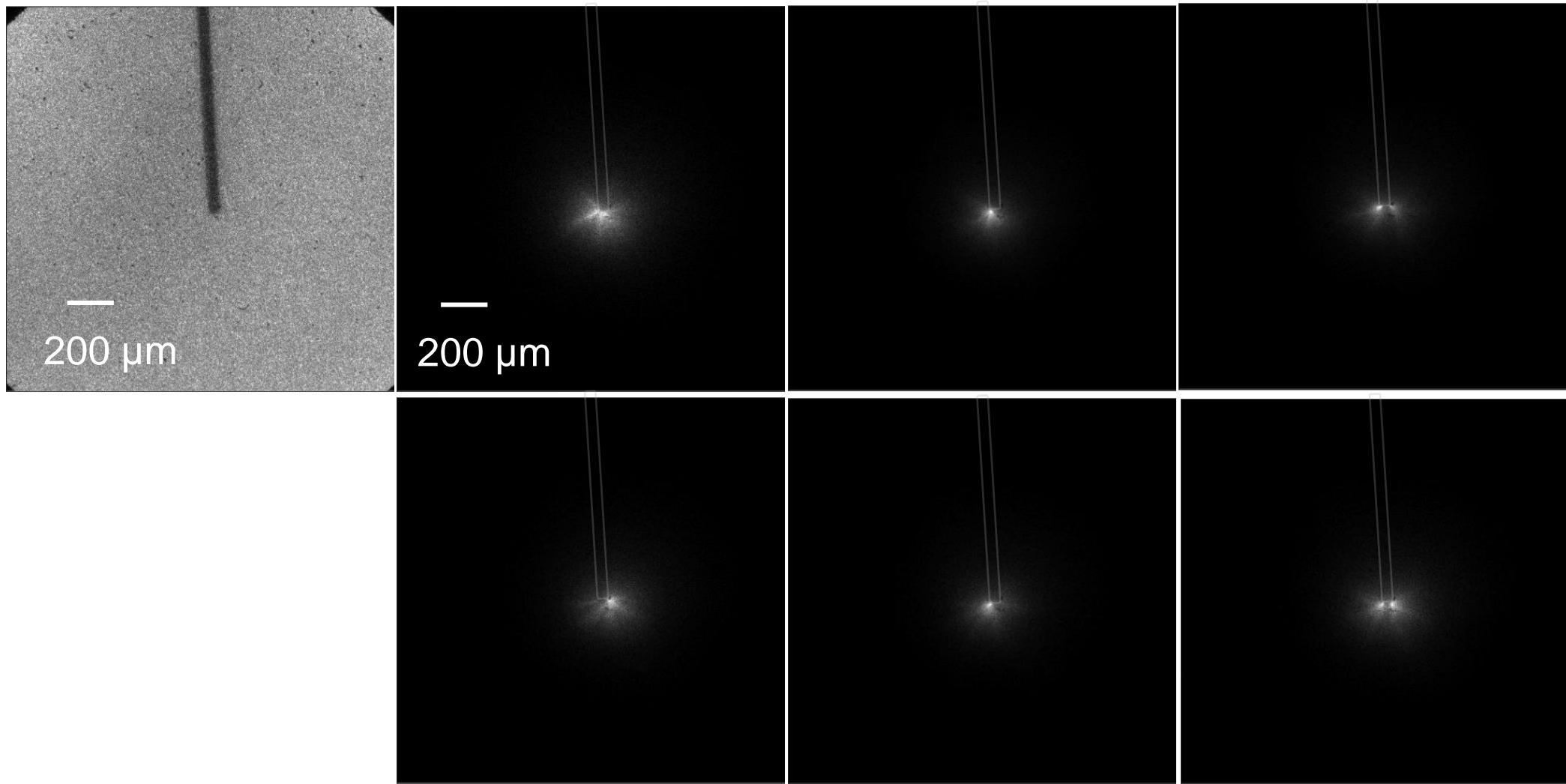


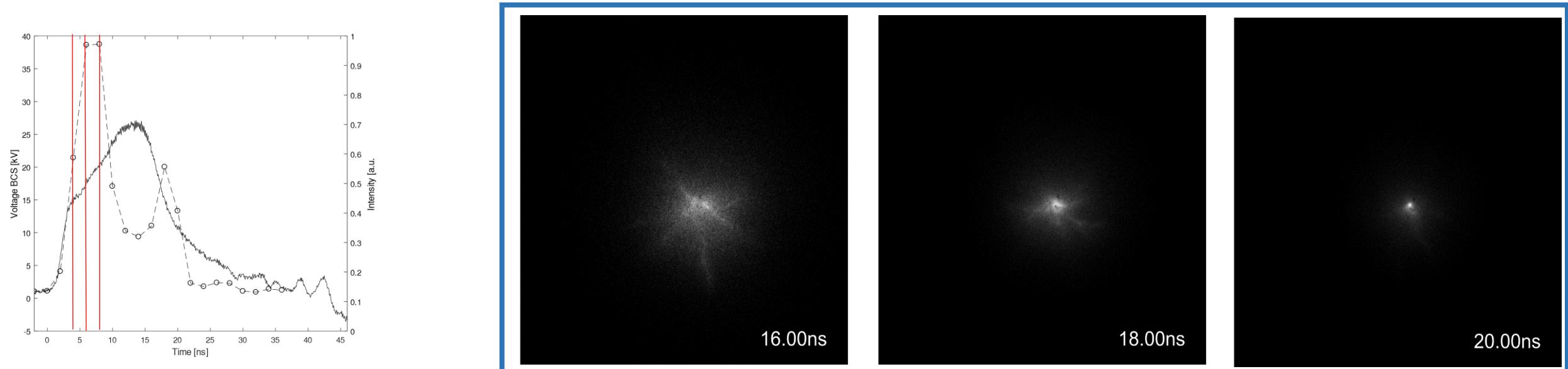
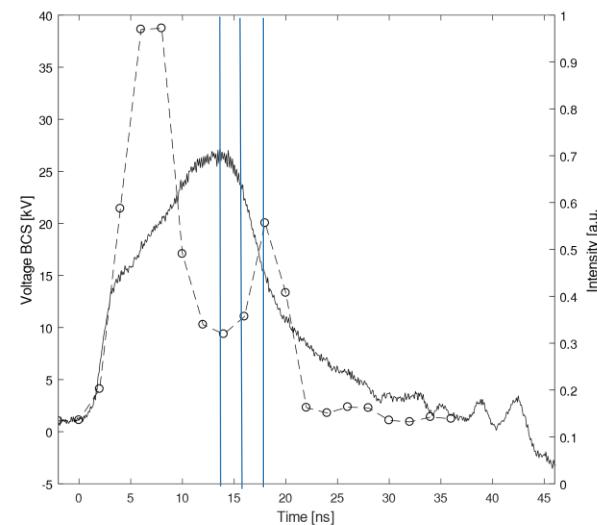
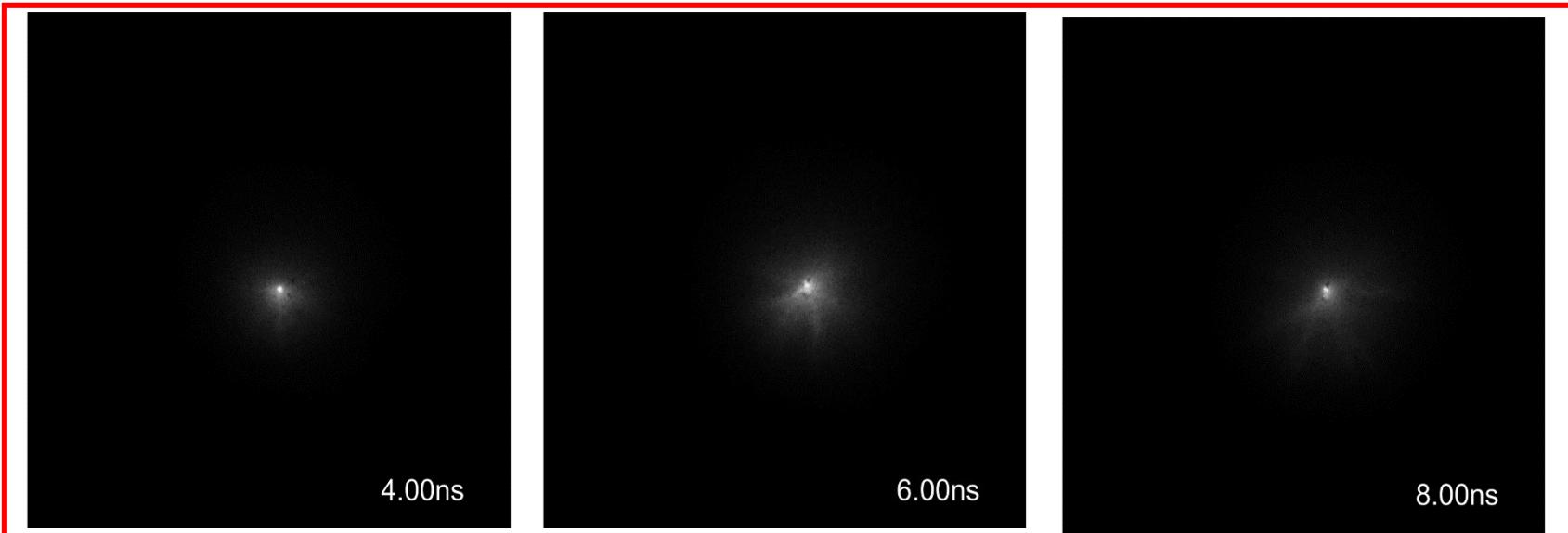
Dissipated energy during plasma ignition within the first 10 ns – comparison to experiments



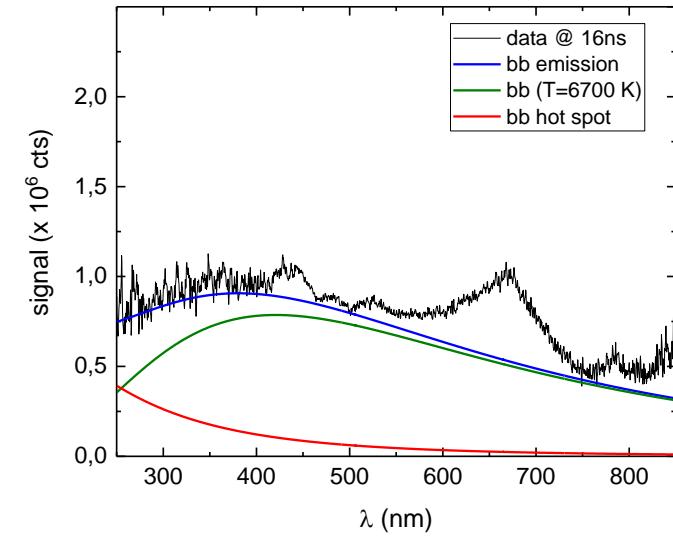
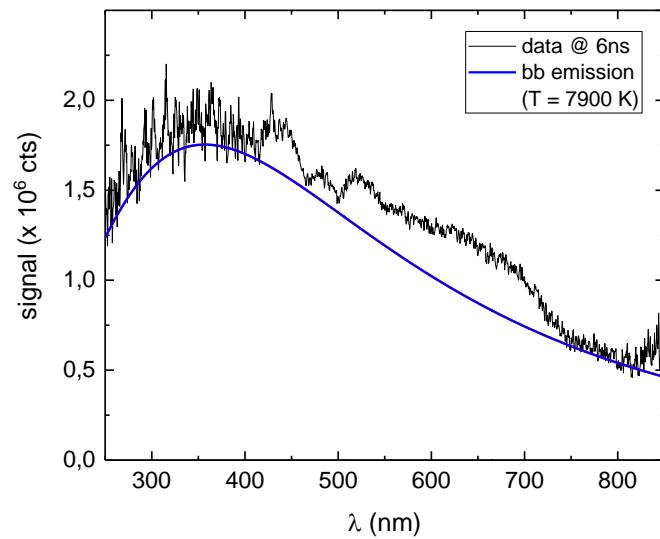
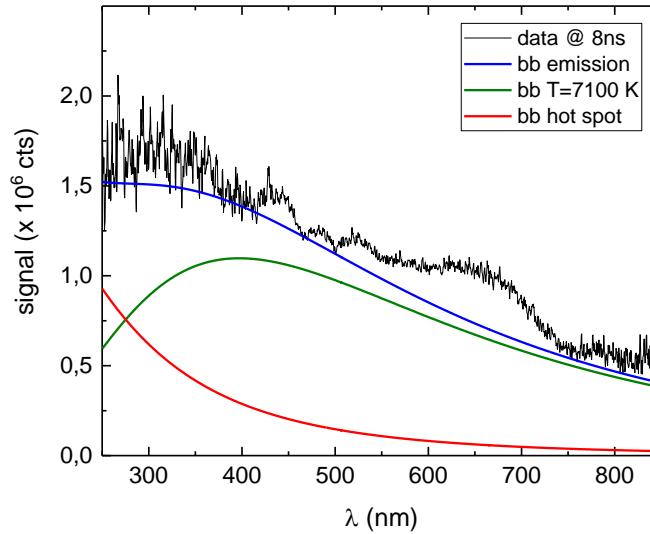
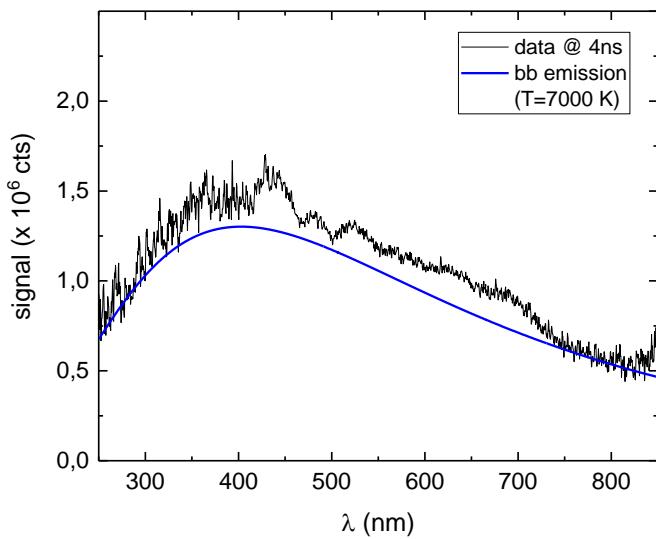
In ns plasmas at very high initial pressure, most of the energy is transferred to the sound wave







Nature of the $1/\lambda^x$ background in emission – cathode spots on the tungsten electrode



$\langle p \rangle_{\text{initial}}$ scales with the dissipated energy

$$p_0 V_0 / \Delta t = P_{\text{electrical}} = U^2 / R$$

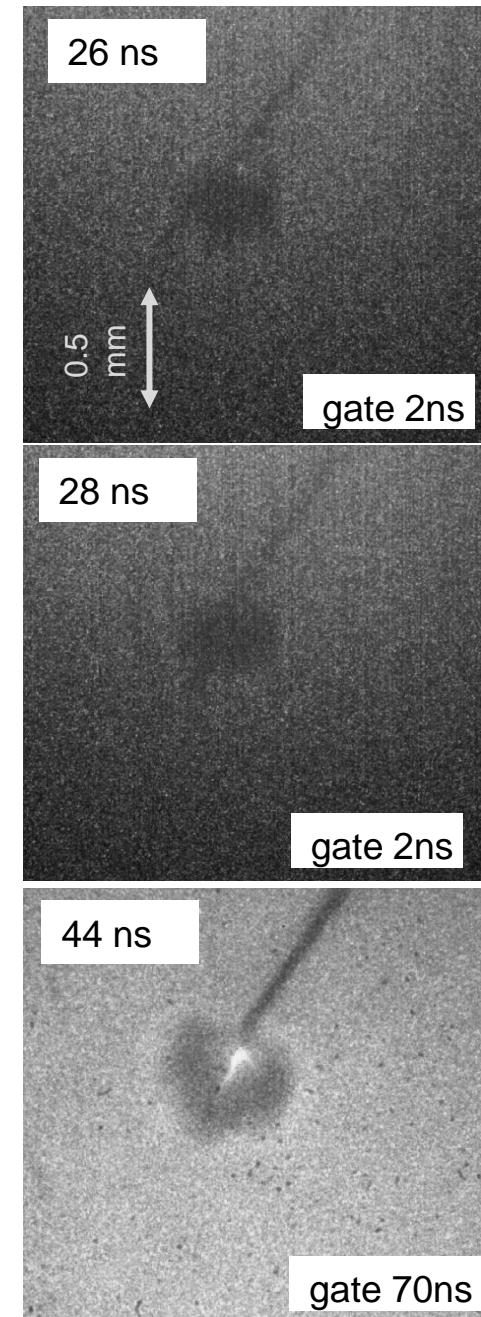
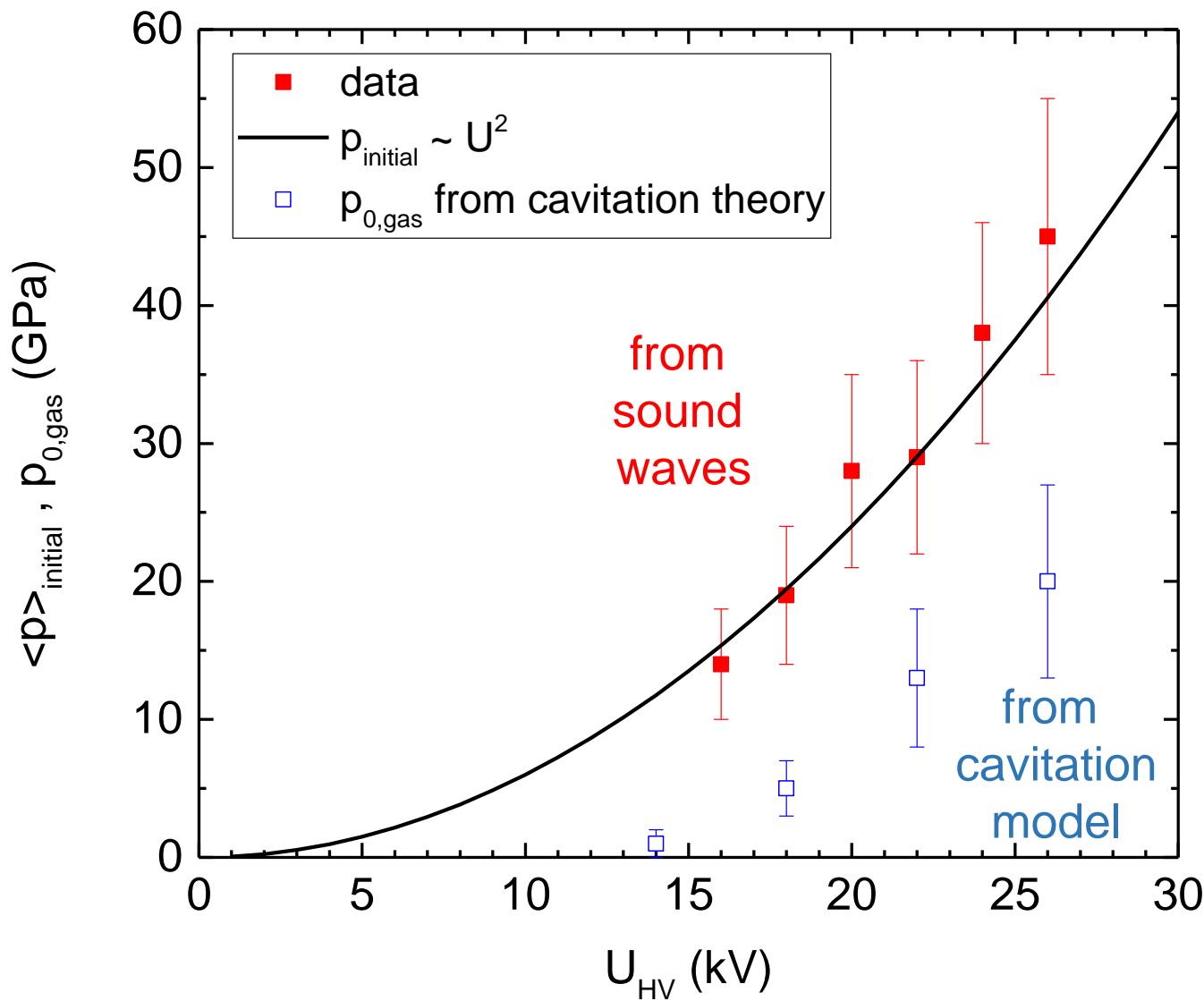
for $R = 2.5 \text{ kOhm}$

p_0, V_0

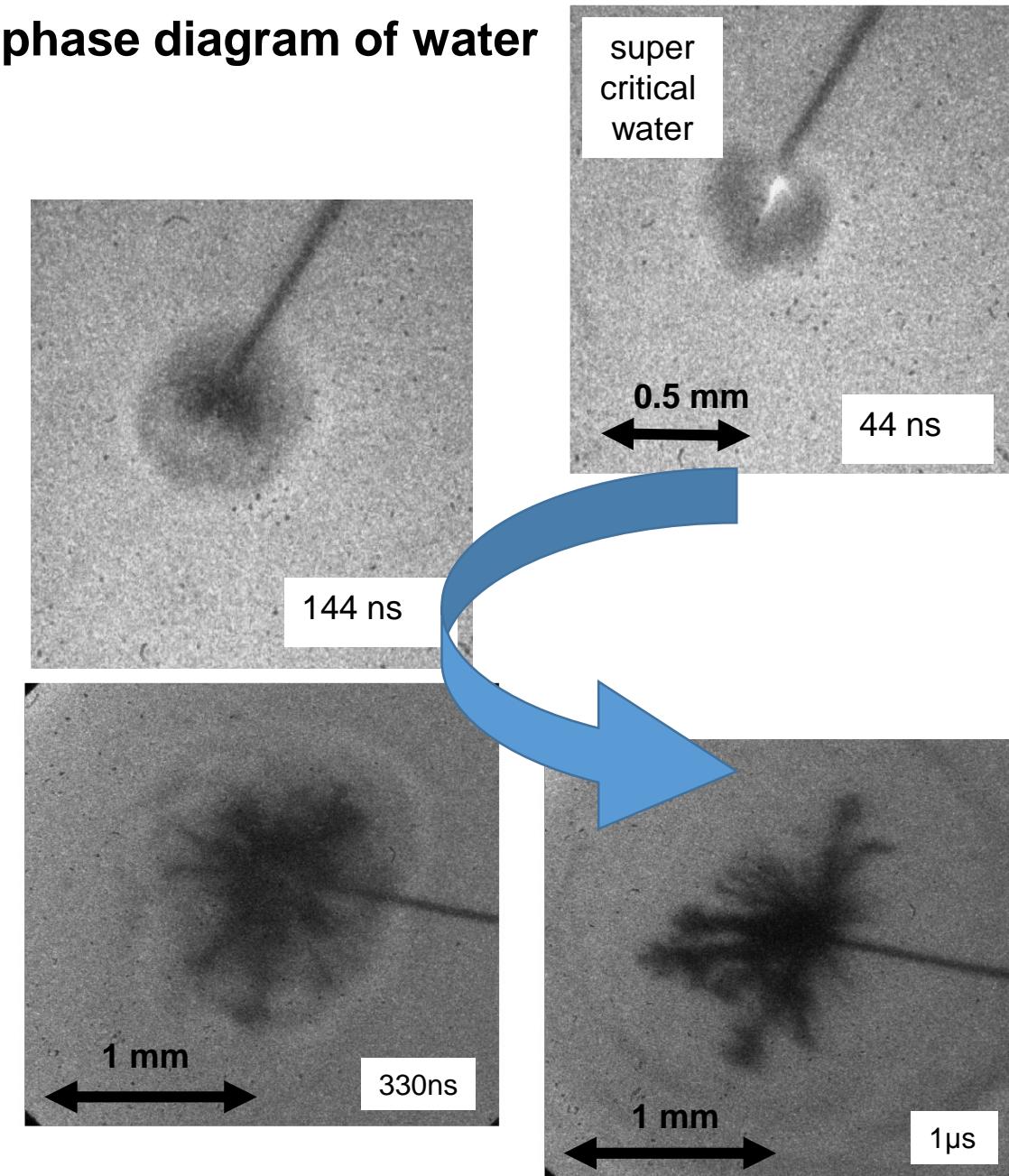
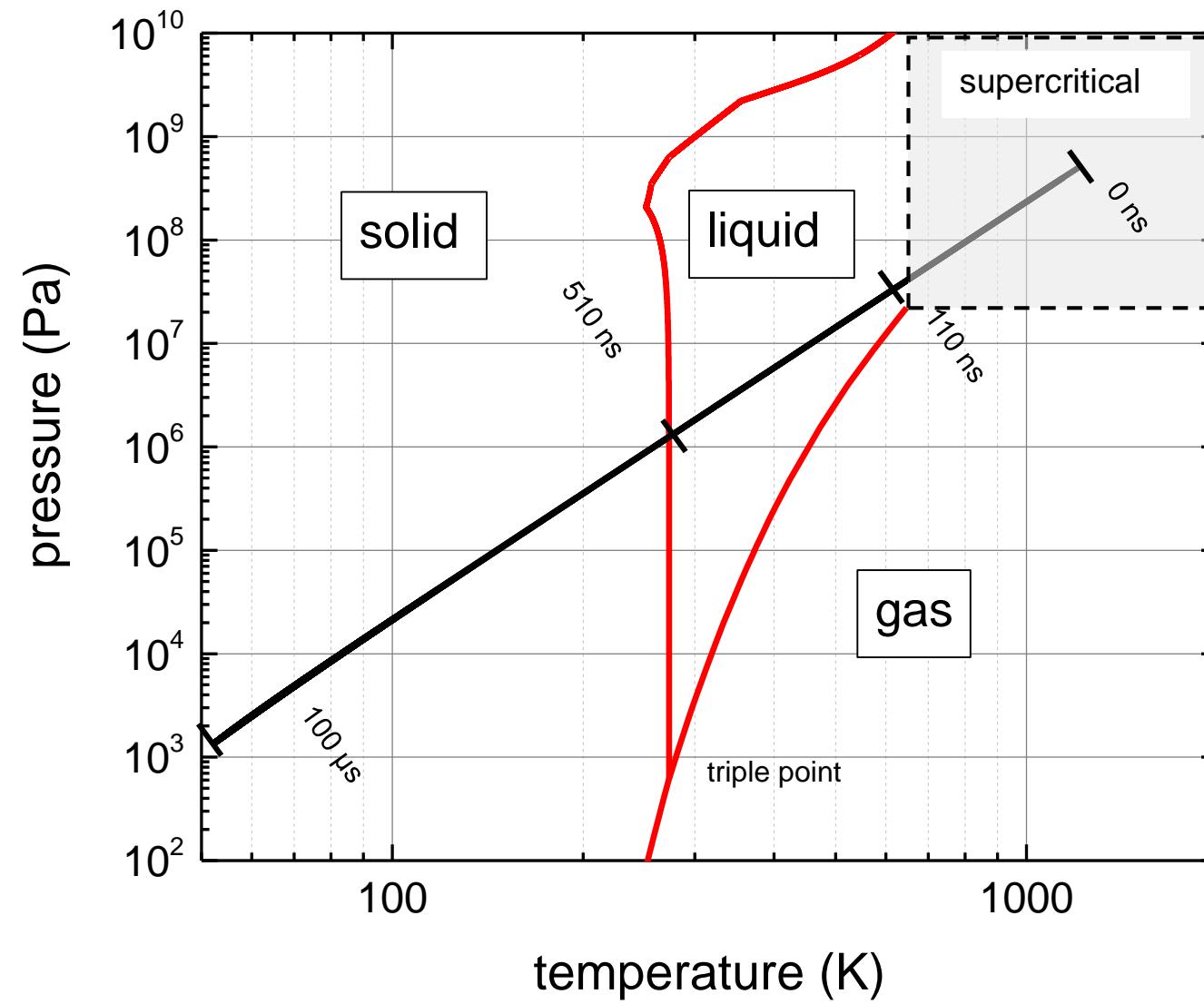
$r = 25 \mu\text{m}$
 $t = 10\text{ ns}$

$T \sim$
2000...
70000 K

Ignition
plasma



Transition of the water vapor in the bubble through the phase diagram of water



Solution for R(t) assuming a compressible liquid and condensation of vapor species

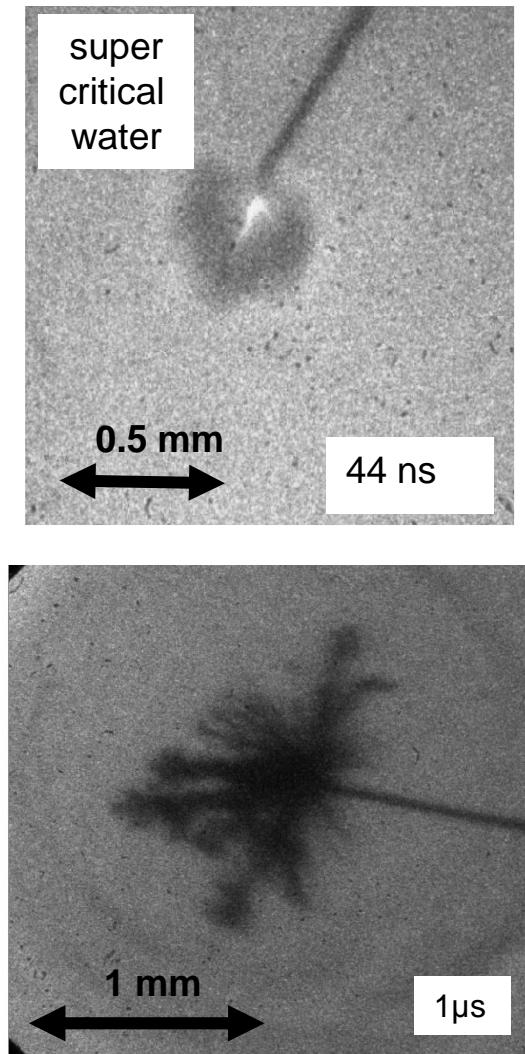
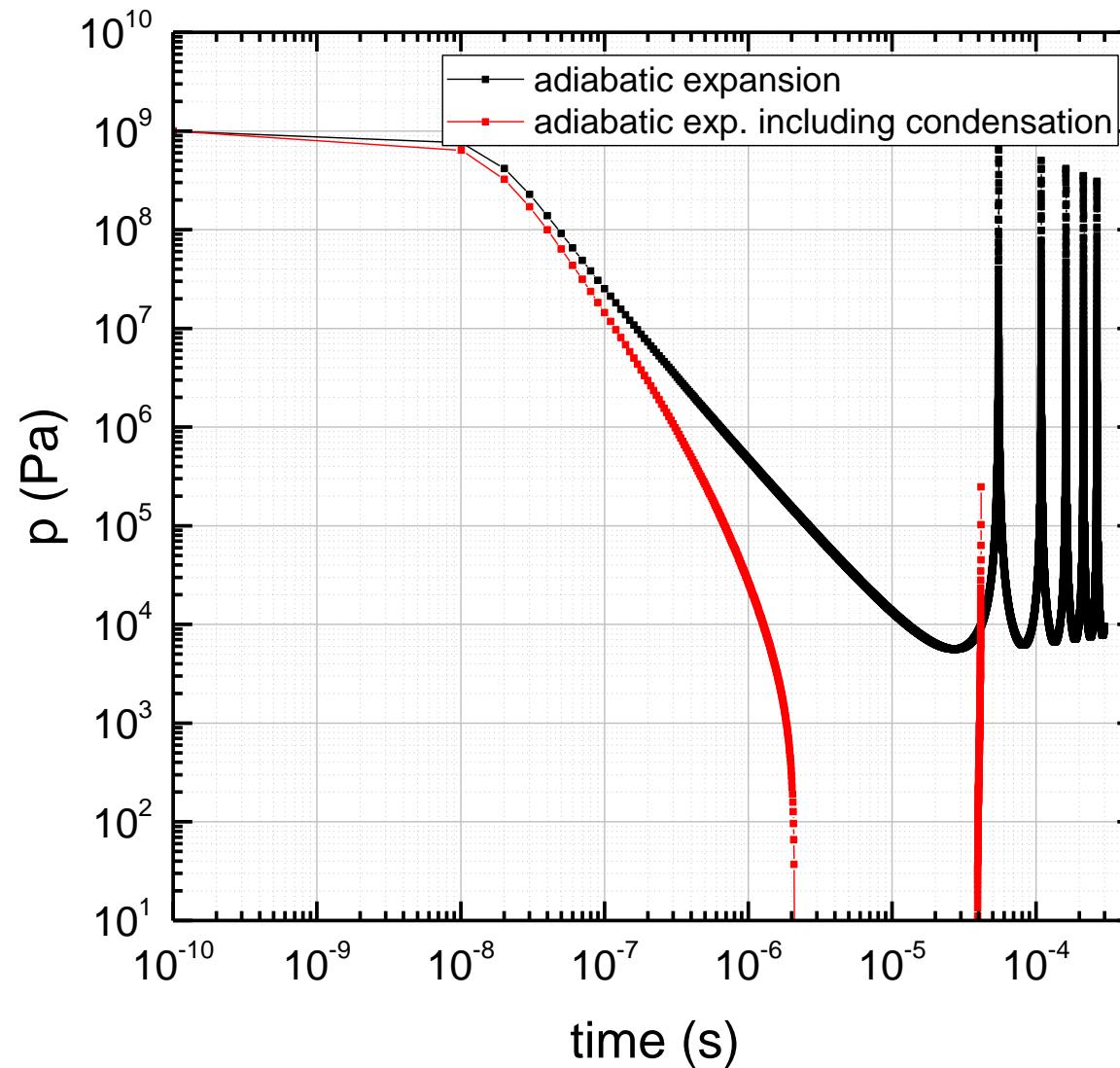
Balance between surface flux and volume loss depending on R,
 β =sticking coeff. = 1

$$\frac{1}{4} n v_{therm} 4\pi R^2 \beta = \frac{4\pi}{3} R^3 n \frac{1}{\tau}$$

$$\tau = \frac{4R}{3v_{therm}\beta}$$

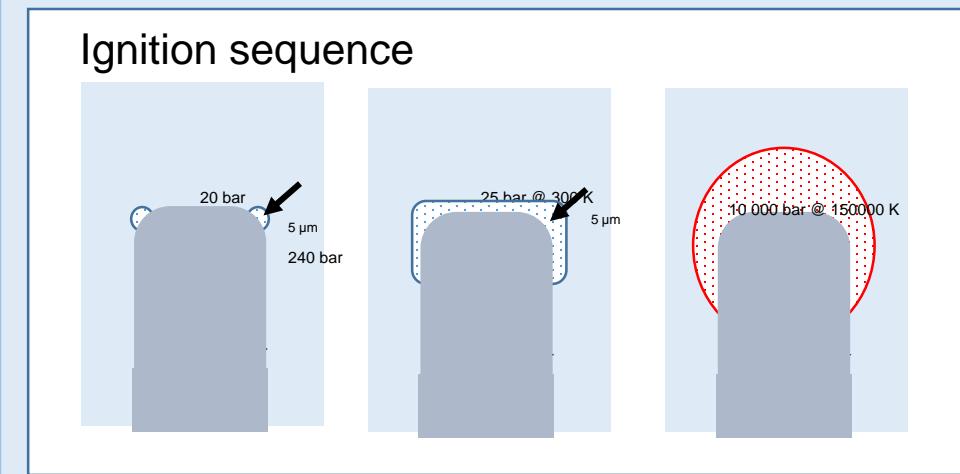
$$p_0 \rightarrow p_0 \exp\left(-\frac{t}{\tau}\right)$$

$$p(t) \rightarrow p_0 \exp\left(-\frac{t}{\tau}\right) \left(\frac{R_0}{R}\right)^{3\gamma}$$



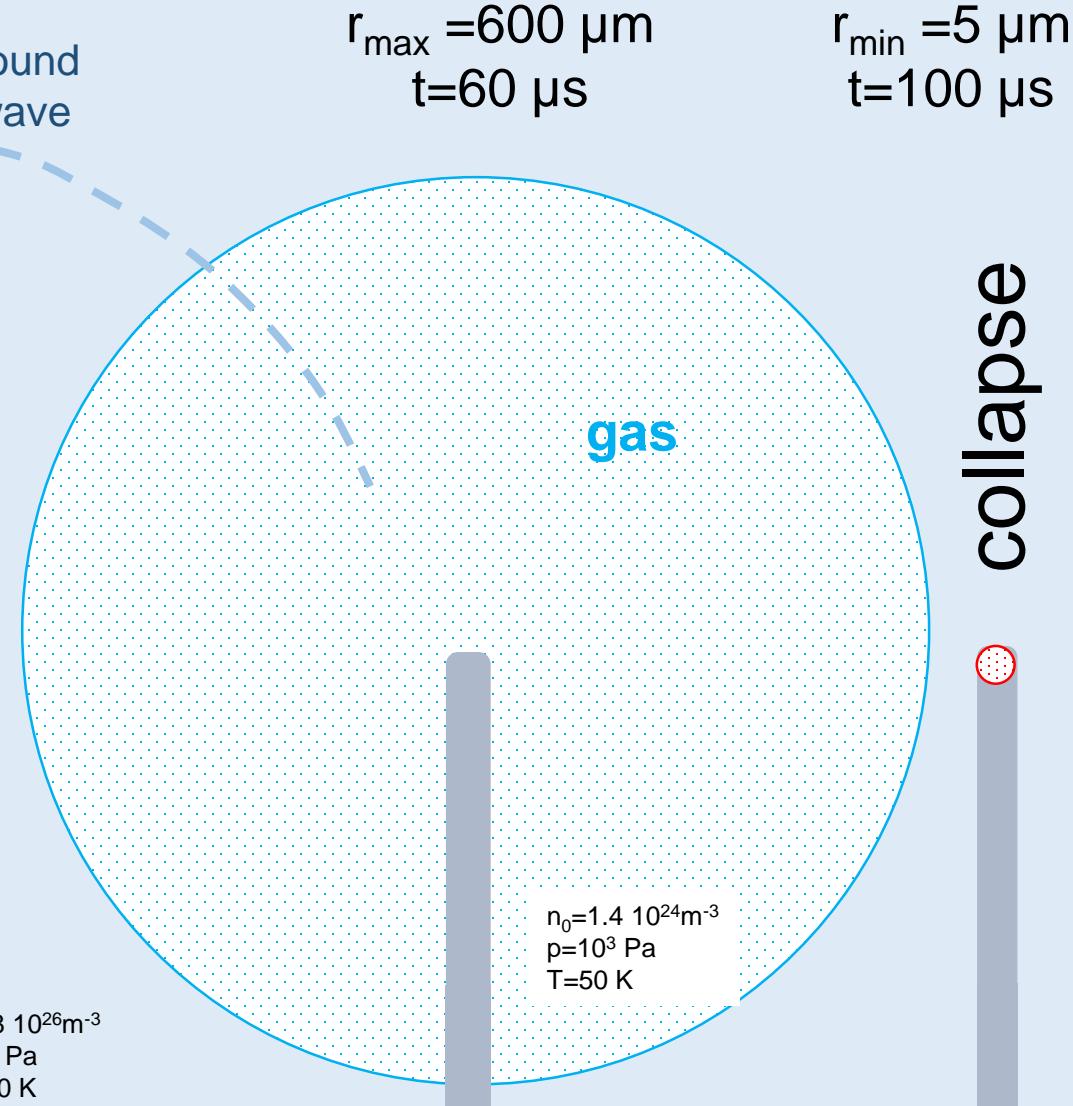
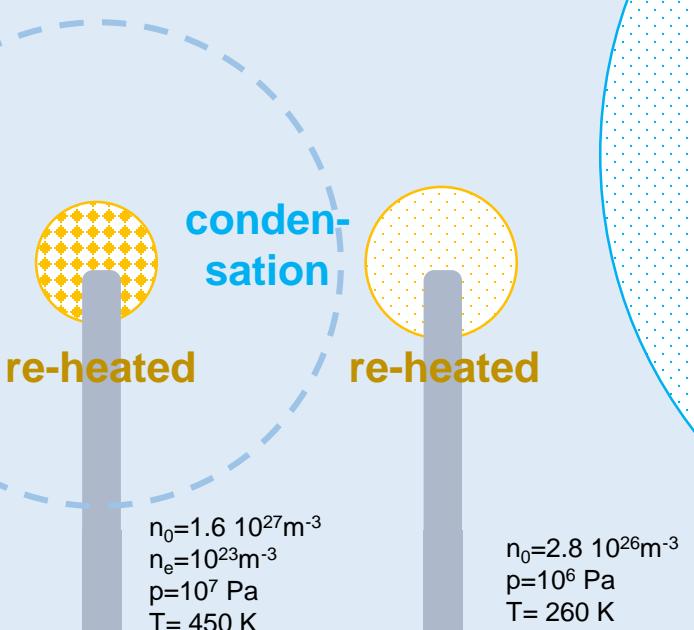
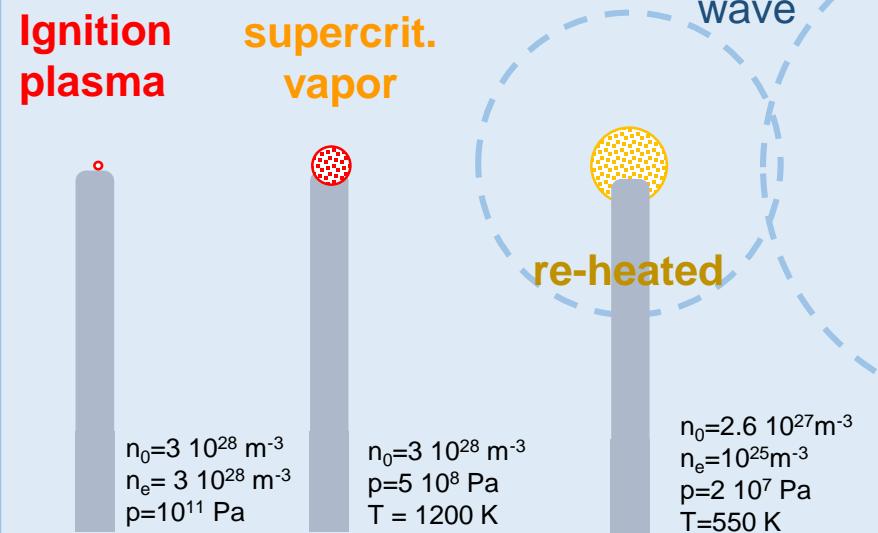
Plasma in a bubble expansion

everything to scale



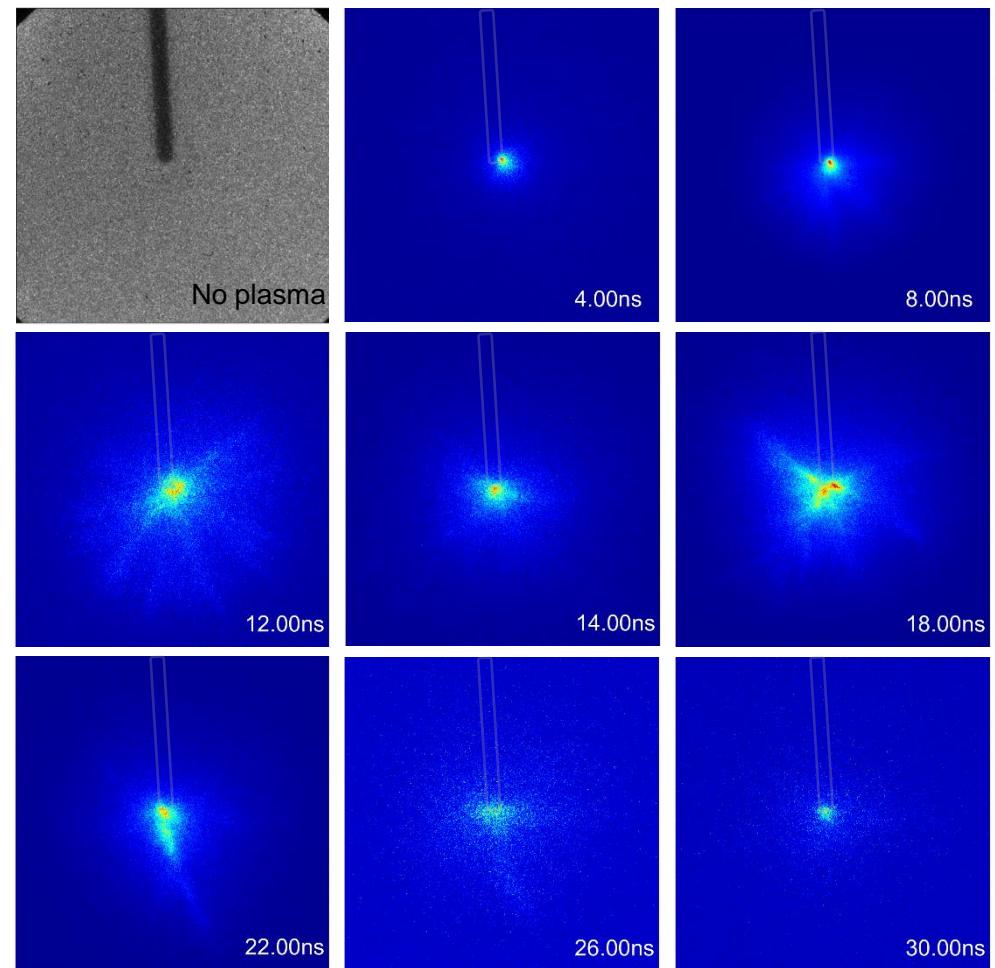
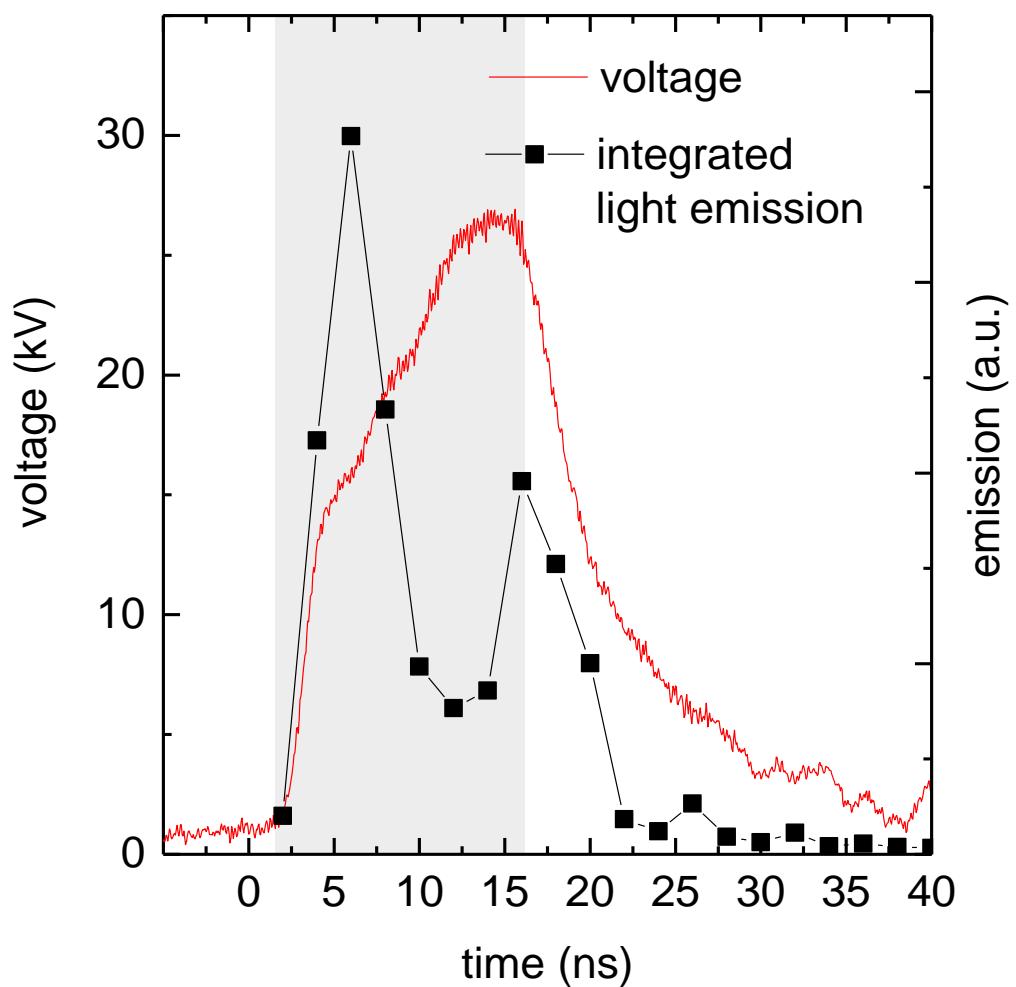
$r = 80 \mu\text{m}$ $t = 200\text{ ns}$

$r = 100 \mu\text{m}$ $t = 400\text{ ns}$



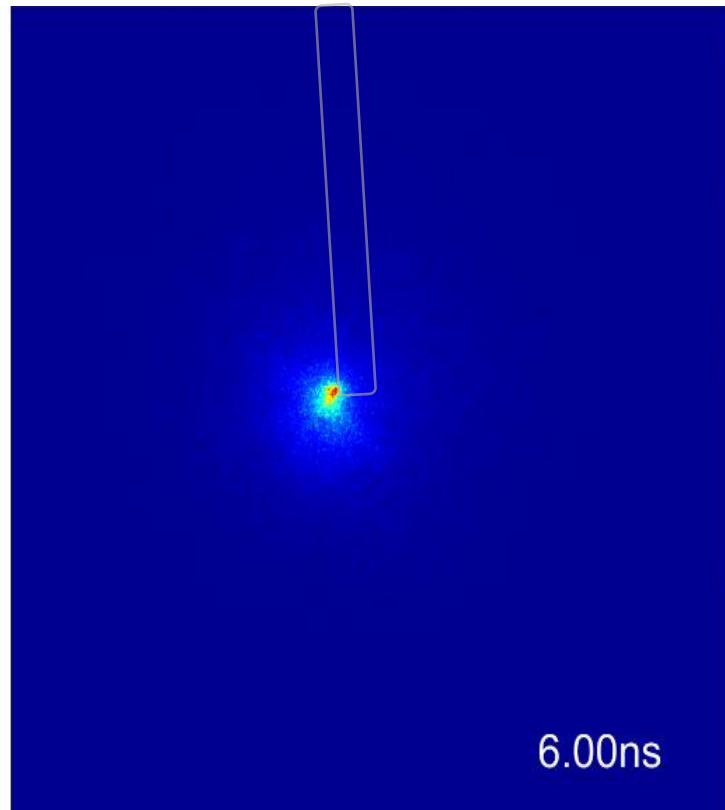
How does the plasma ignites ?

Experiment 22kV, 1Hz, distilled water, $t_{\text{gate}}=2\text{ns}$, $t_{\text{step}}=2\text{ns}$

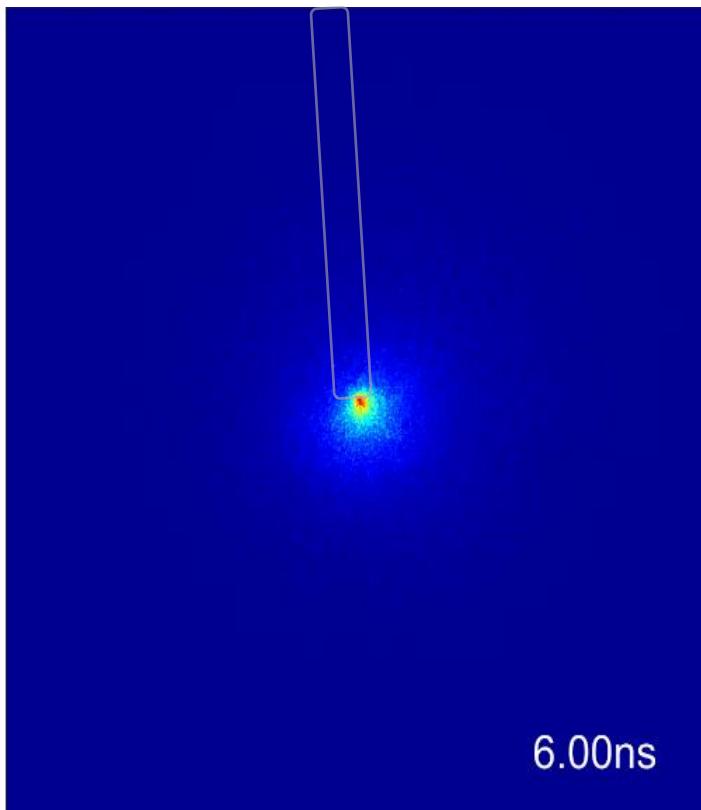


How does the plasma ignites ?

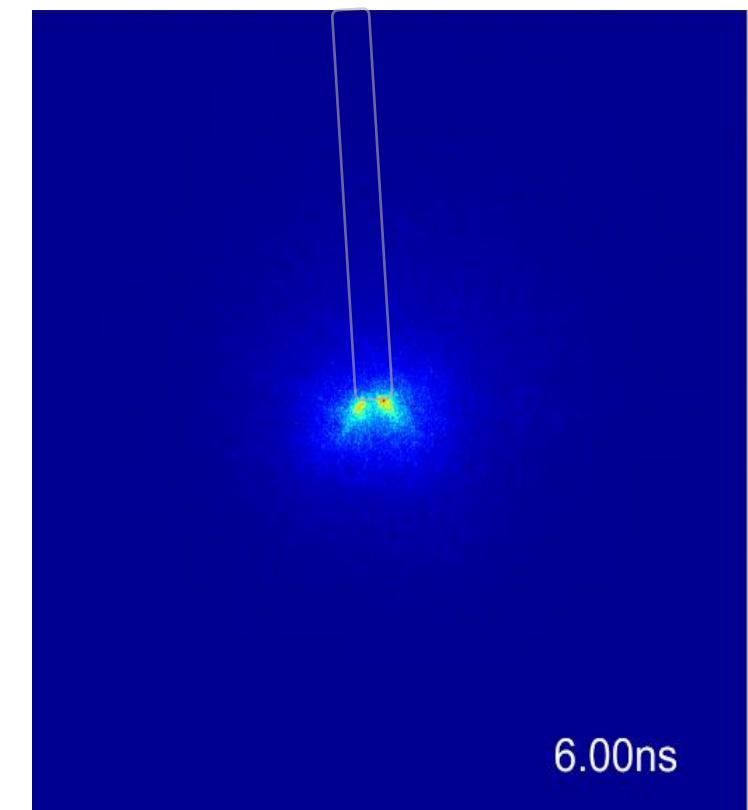
Experiment 22kV, 1Hz, distilled water, $t_{\text{gate}}=2\text{ns}$, $t_{\text{step}}=2\text{ns}$



6.00ns



6.00ns



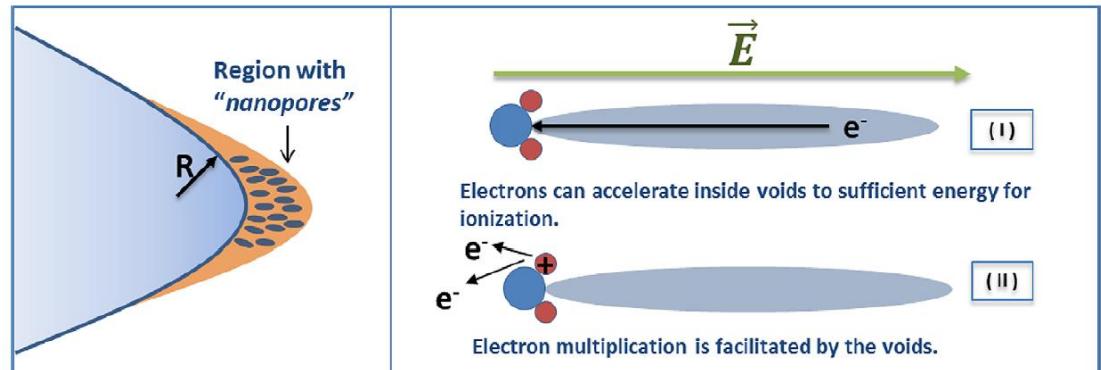
6.00ns

Comparison to Modeling of Streamer Propagation in 80 bar supercritical N₂

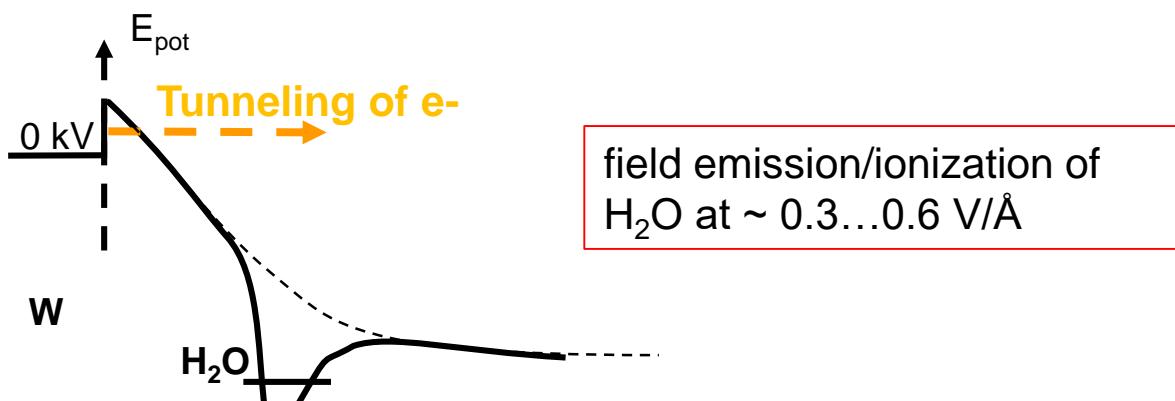
ns plasma in liquids:

Due to the inertia of the liquid, no bubble formation on ns time scale, so standard ignition in the gas trapped inside a bubble according to Paschen law is impossible

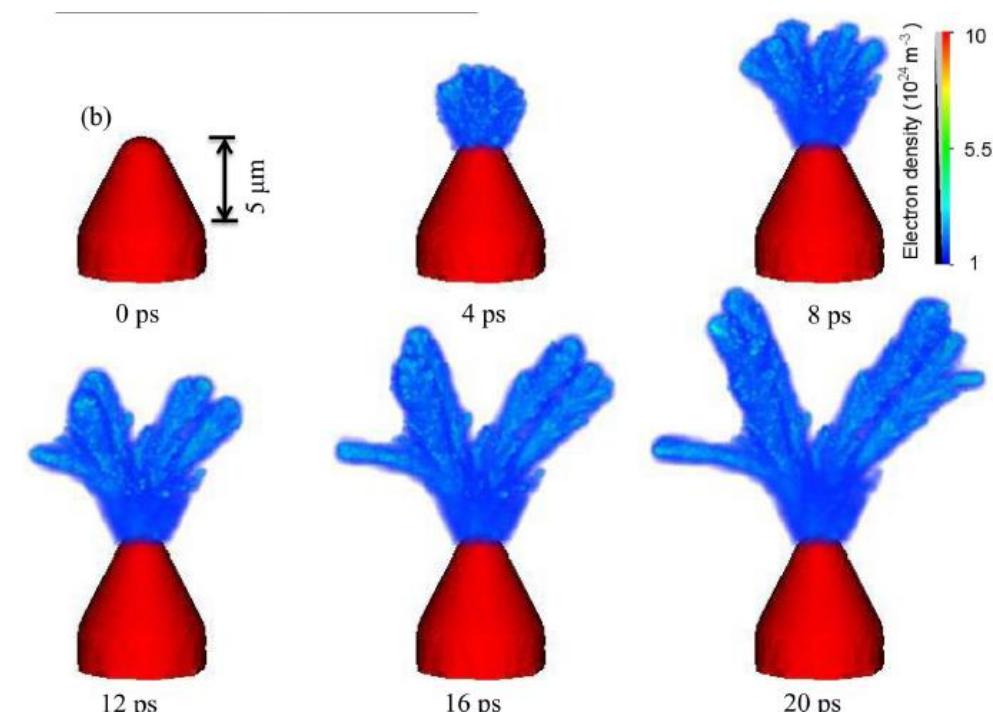
Nanopores (Dobryn et al. JPD 46, 355201 (2013))



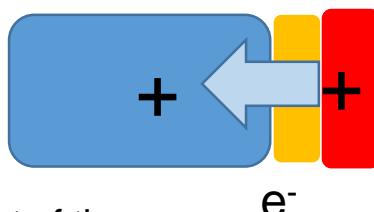
Field emission, field ionization



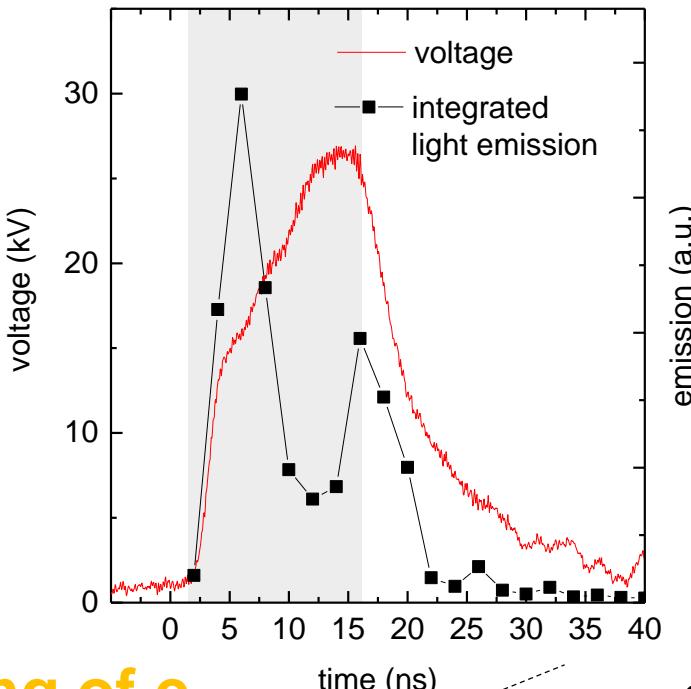
From A. Sun, J. Teunissen, U. Ebert
IEEE Trans. Plas. Sci. 42, 2416 (2014)



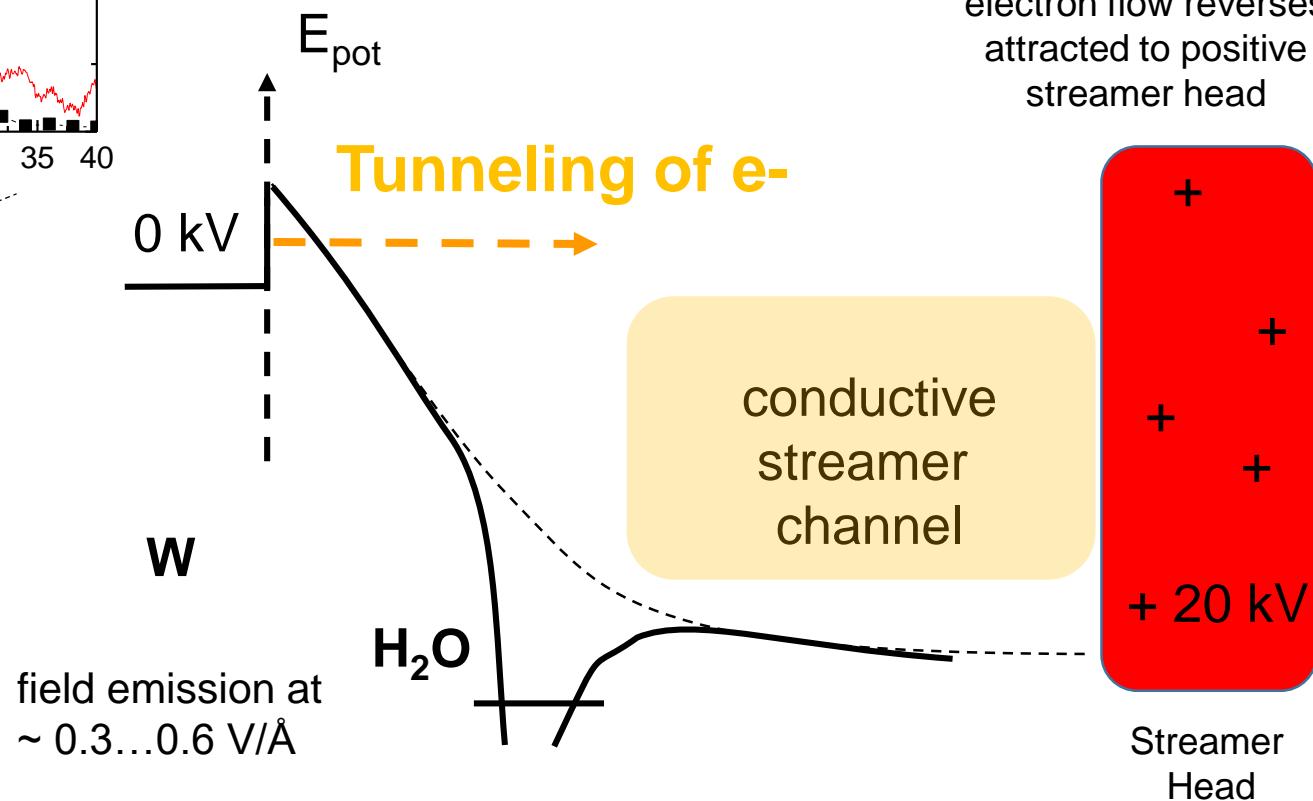
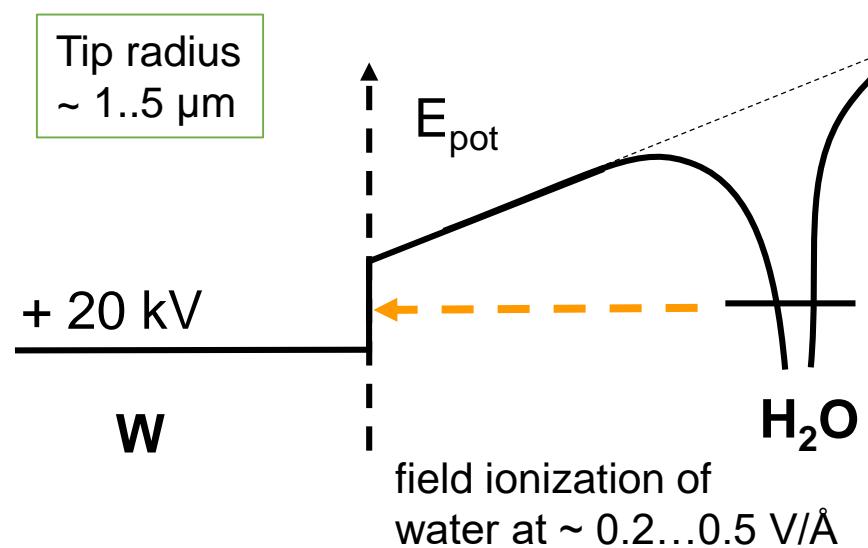
Field emission effects



Start of the pulse
water
is polarized
Field ionization

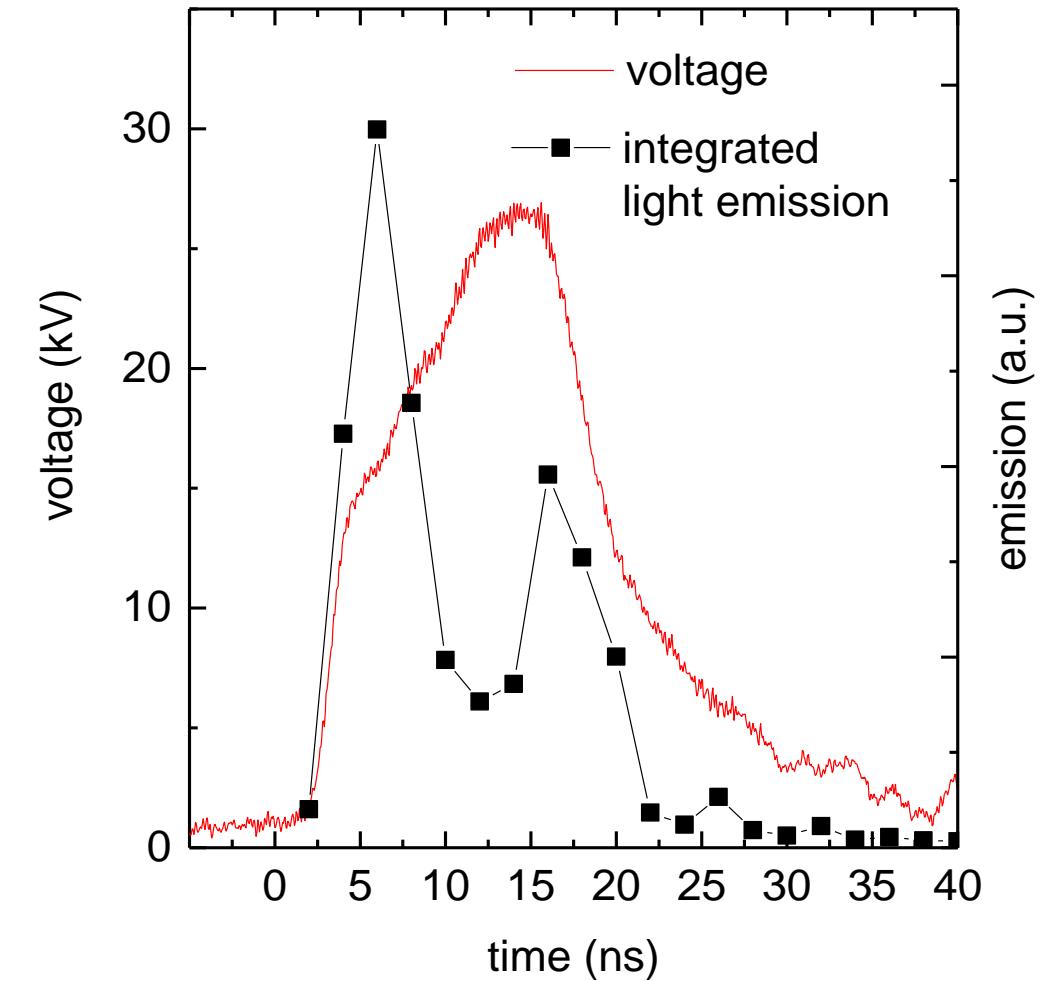
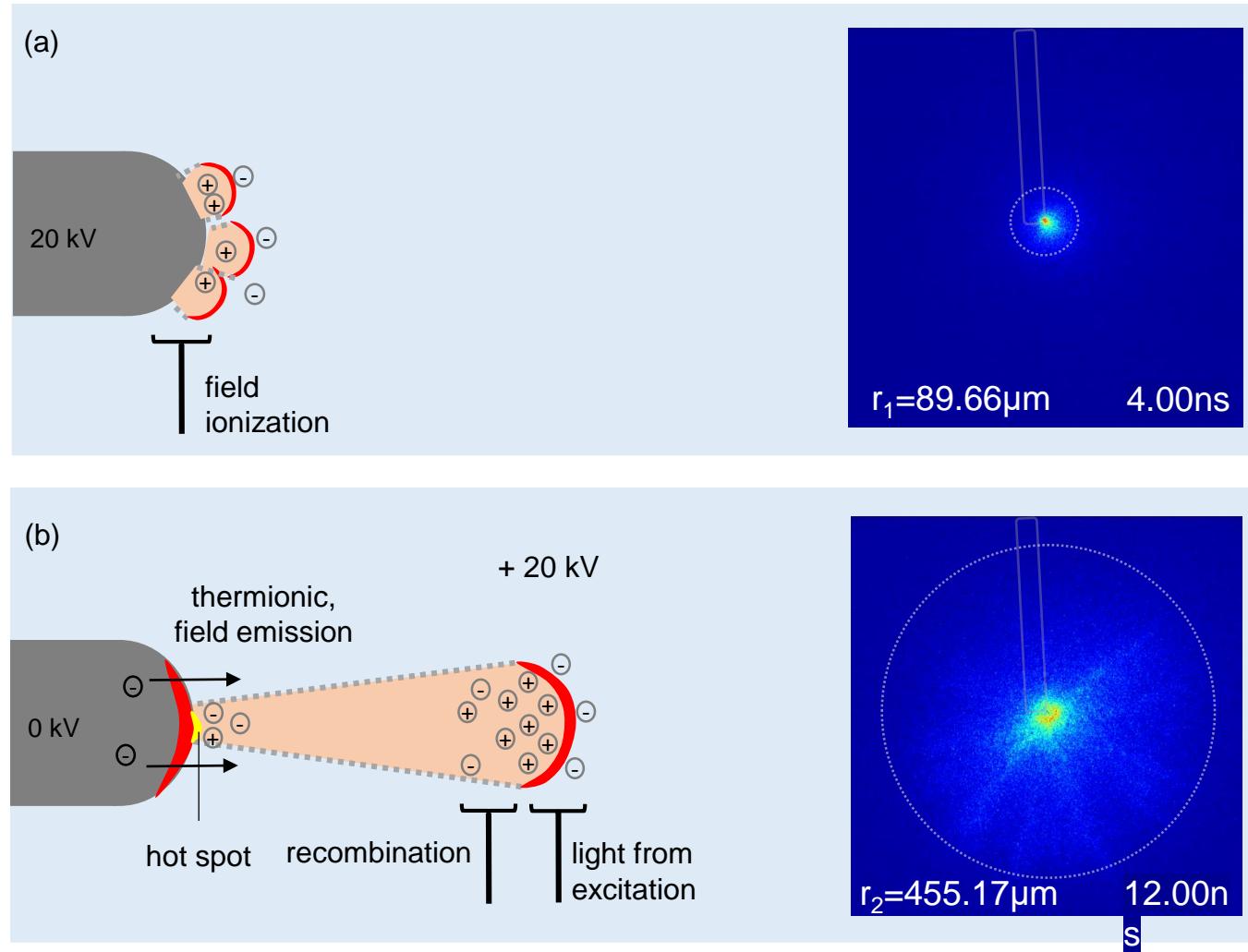


Tunneling of e-



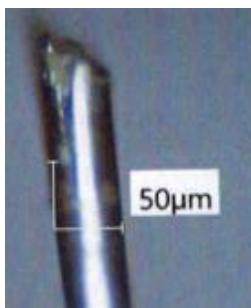
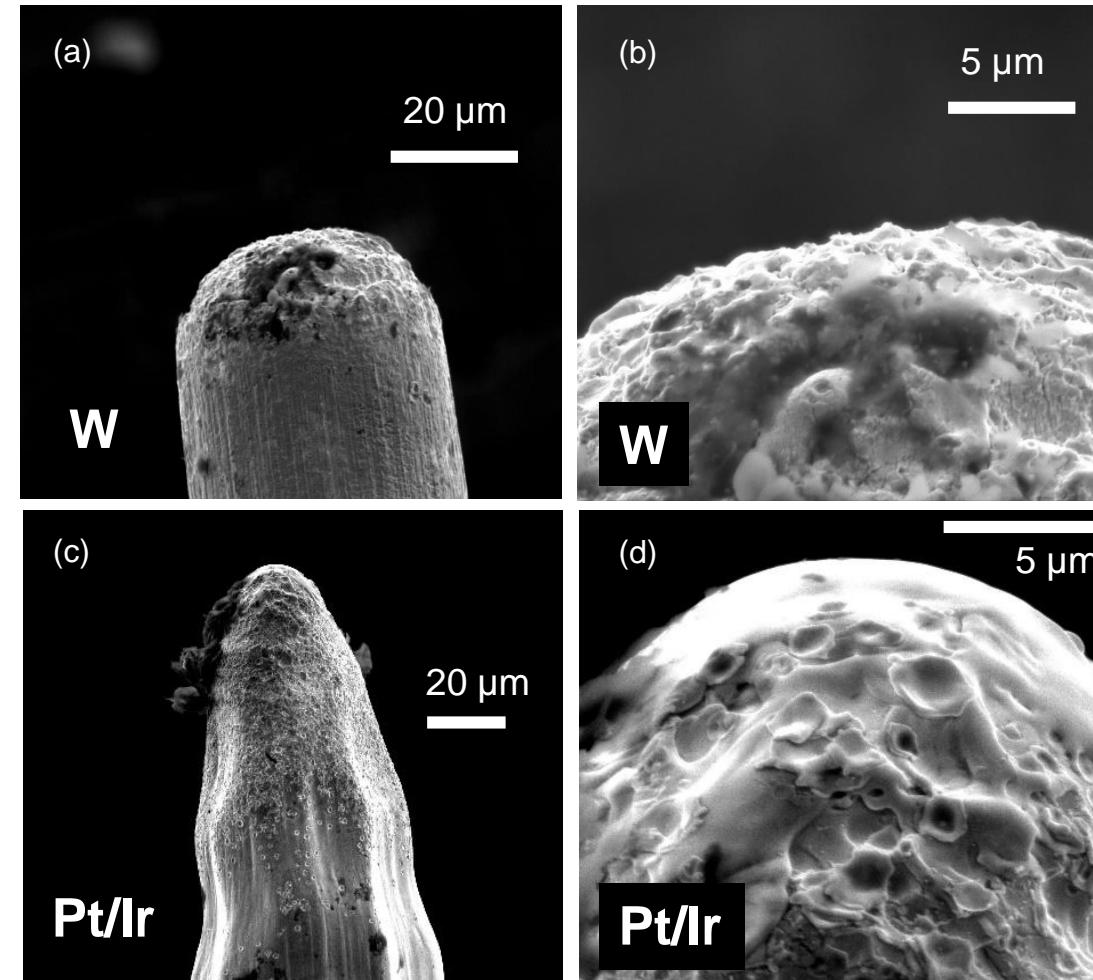
How does the plasma ignites ?

Experiment 22kV, 1Hz, distilled water, $t_{\text{gate}}=2\text{ns}$, $t_{\text{step}}=2\text{ns}$



Tungsten tip vs. Platinum/Iridium tip

- W as high melting temperature 3700 K
- W therm. conduct. 115 W / m K
- W oxidizes, Crystallites are formed
- Pt melting temperature 2100 K
- Pt thermal conduct. 82 W/ m K
- Ir melting temperature 2700 K



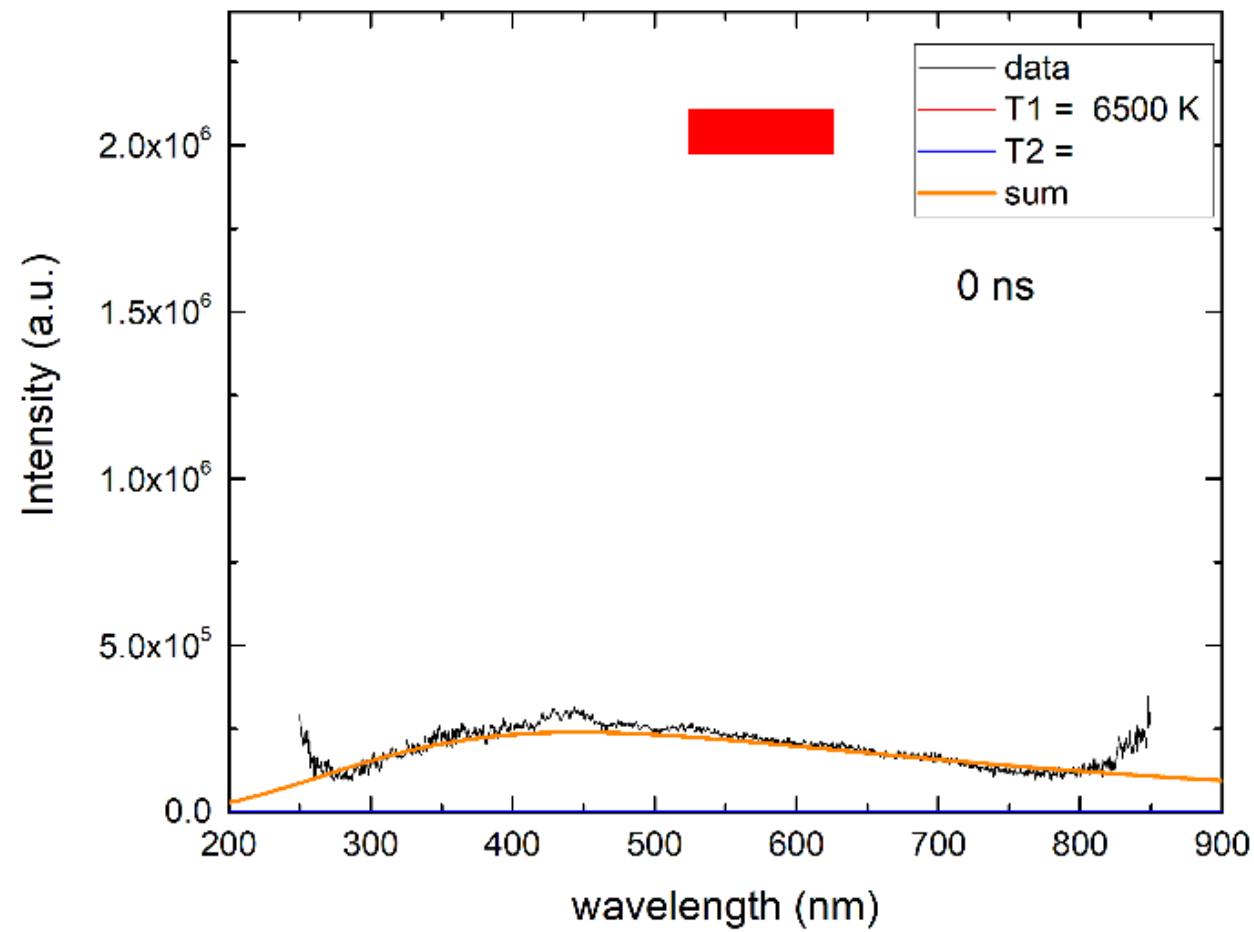
Permanent operation possible

Ignition stops after 1 h of operation

Temperatures from emission spectra of the plasma into the UV

Spectrum (first 50 ns, time resolution 2 ns)

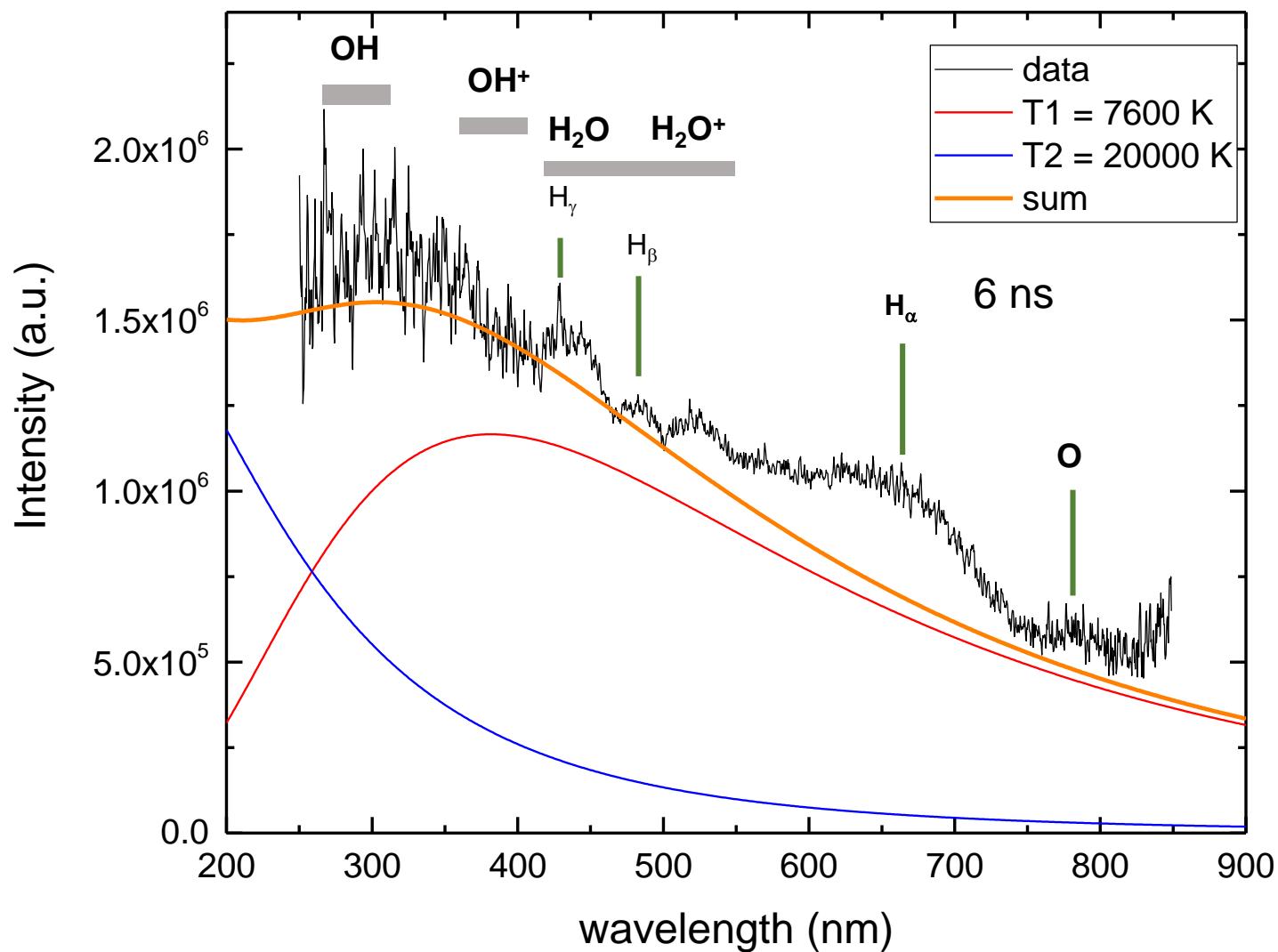
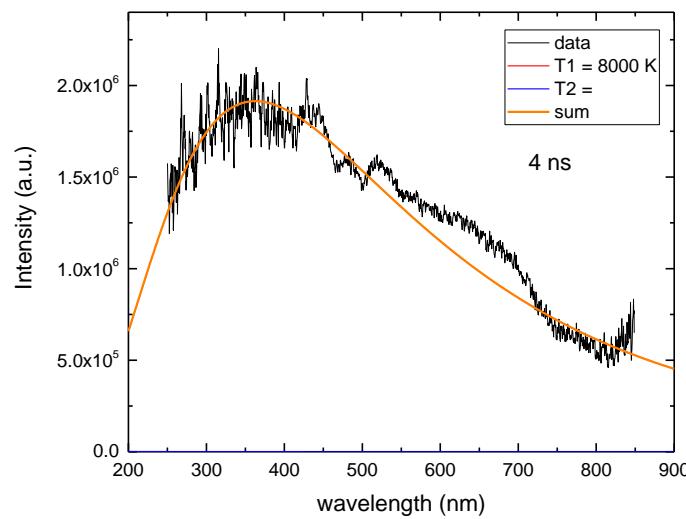
- Recombination
- H₂, H₂O recomb. continuum
- OH, OH⁺, H₂O, H₂O⁺ bands
- Black body radiation (W)
- H_α
- O(777nm)



Temperatures from emission spectra of the plasma UV part

Spectrum (first 50 ns, time resolution 2 ns)

- Recombination
- H_2 , H_2O recomb. continuum
- OH, OH^+ , H_2O , H_2O^+ bands
- Black body radiation (W)
- H_α
- O(777nm)

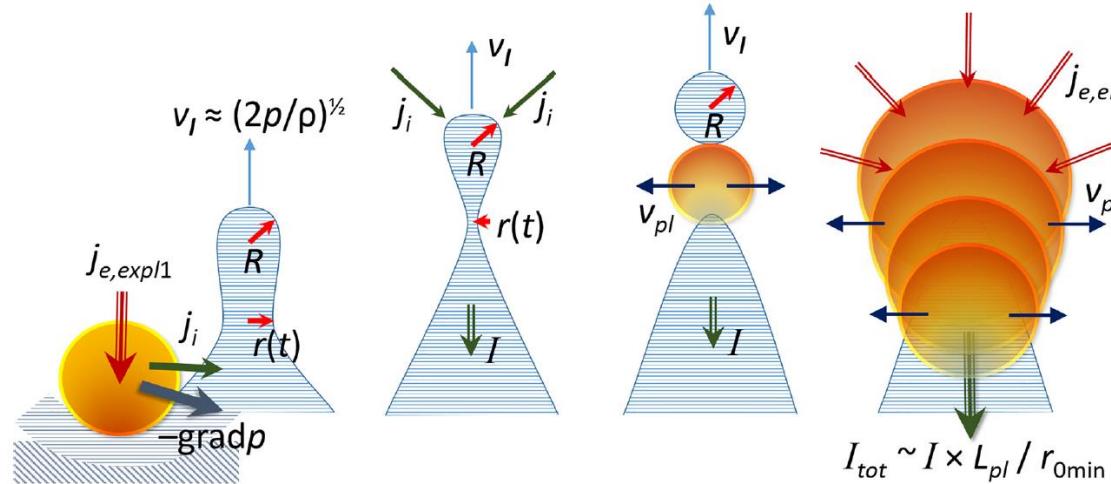


Nature of the $1/\lambda^x$ background in emission

Possible sources for $1/\lambda^x$ background

- H₂O, H₂ recombination background (300...500 nm, no part down to 250 nm)
- $1/\lambda^2$ recombination radiation free-free, bound-free
- $1/\lambda^5$ long wavelength part of the emission of a hot cathode spot

Example cathode spot in vacuum arcs T (15000...2000 K)

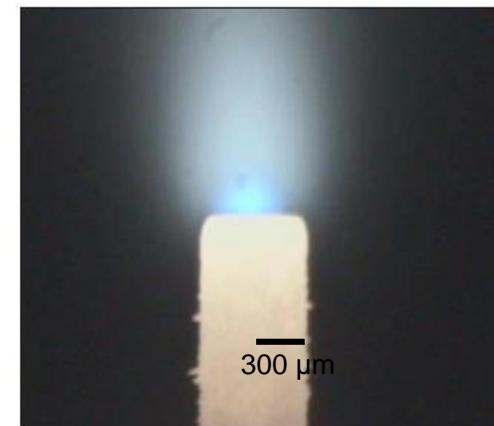


initial
droplet $R=1\mu\text{m}$

Tsvetoukh,
Phys. Plasmas 25, 53504 (2018)

Coulomb-explosion of the neck
leads to metal plasma
(10 μm diameter, $n_e \sim 10^{26} \text{ m}^{-3}$,)

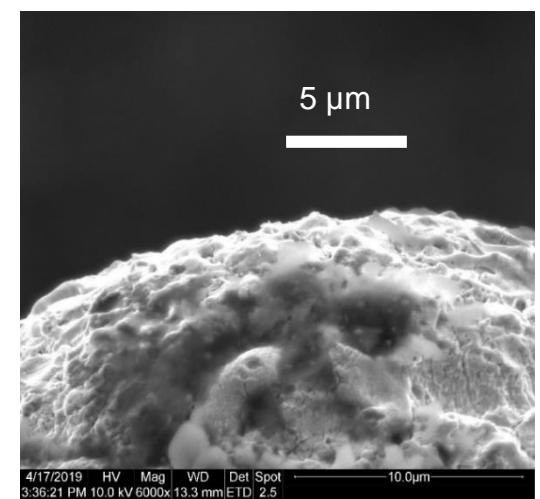
Example cathode spots
in arc lamps



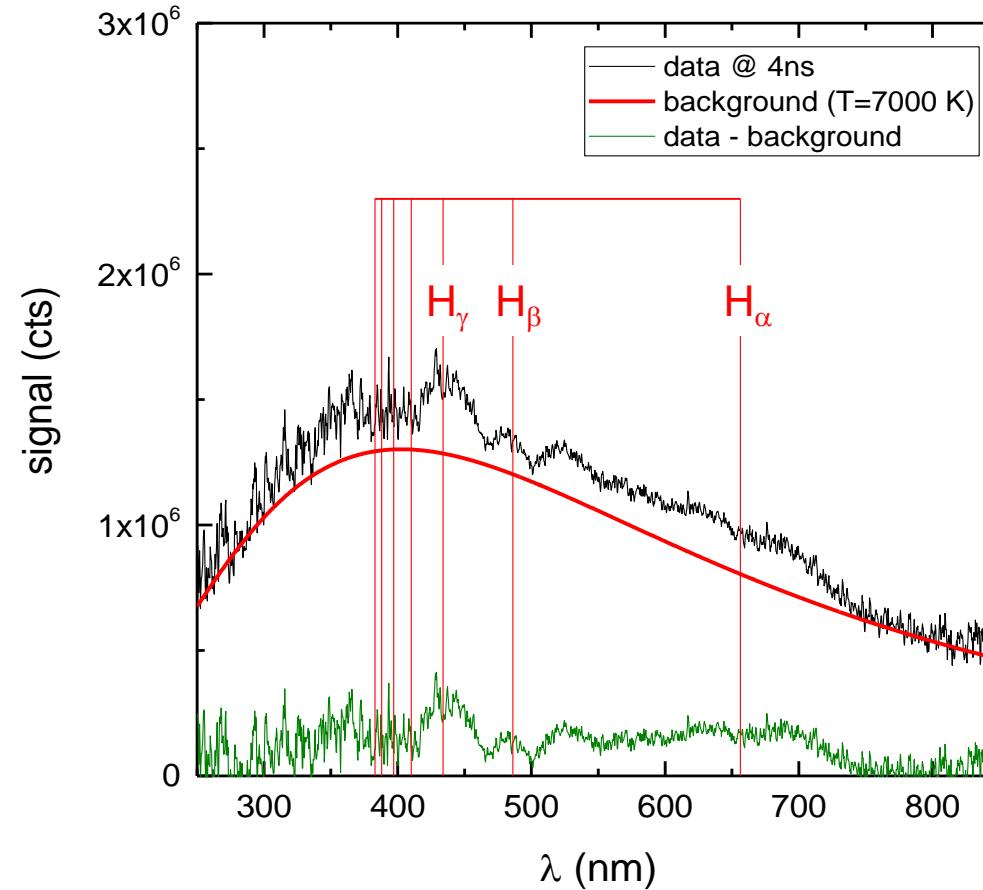
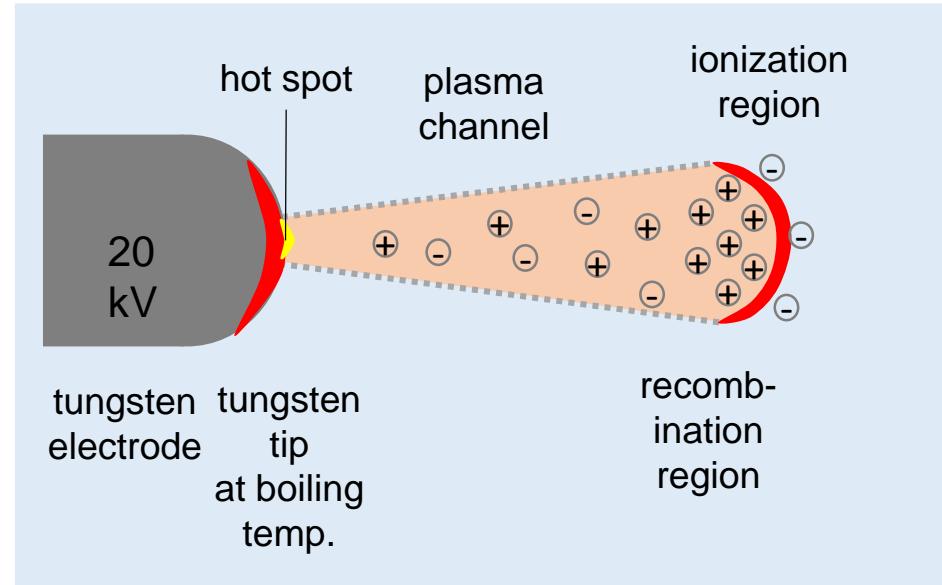
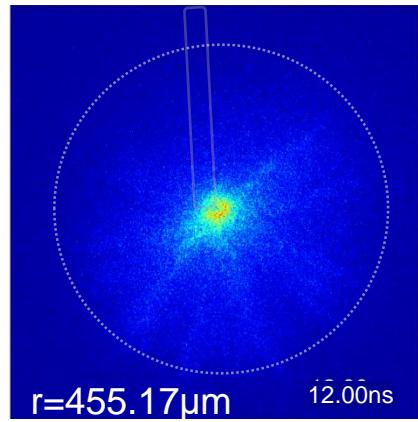
(d) $i_{arc} = 15.0 \text{ A}$

A. Bergner et al. PSST 23, 054005 (2014)

Cathode spots at W tips

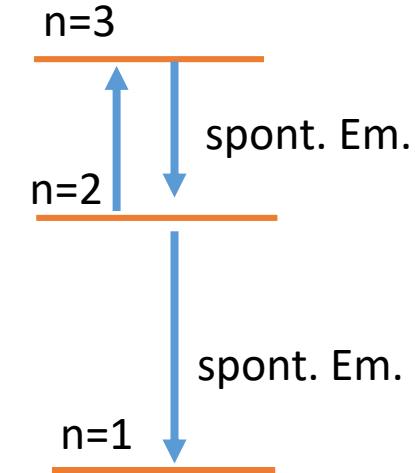
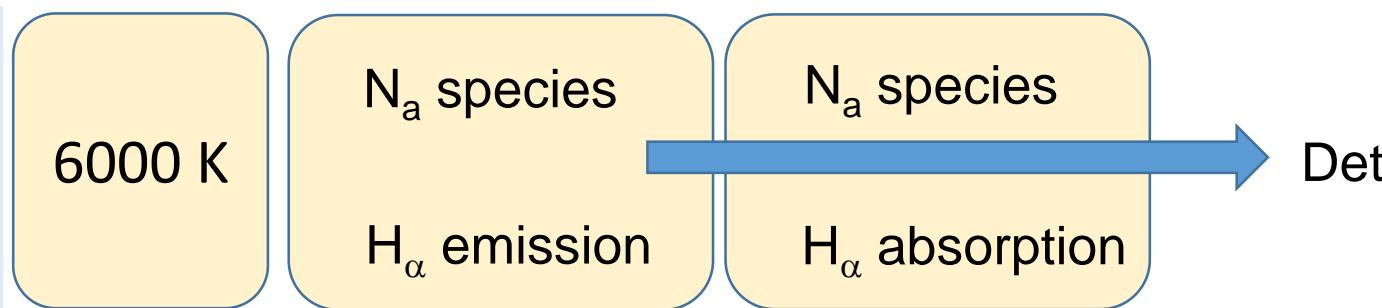
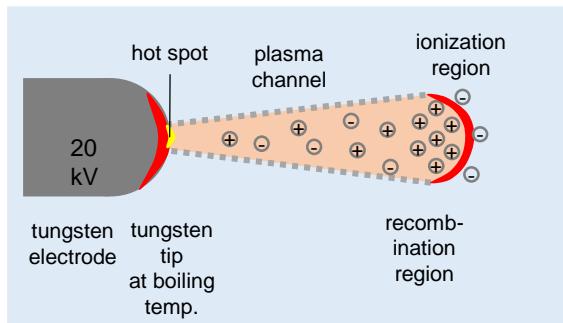
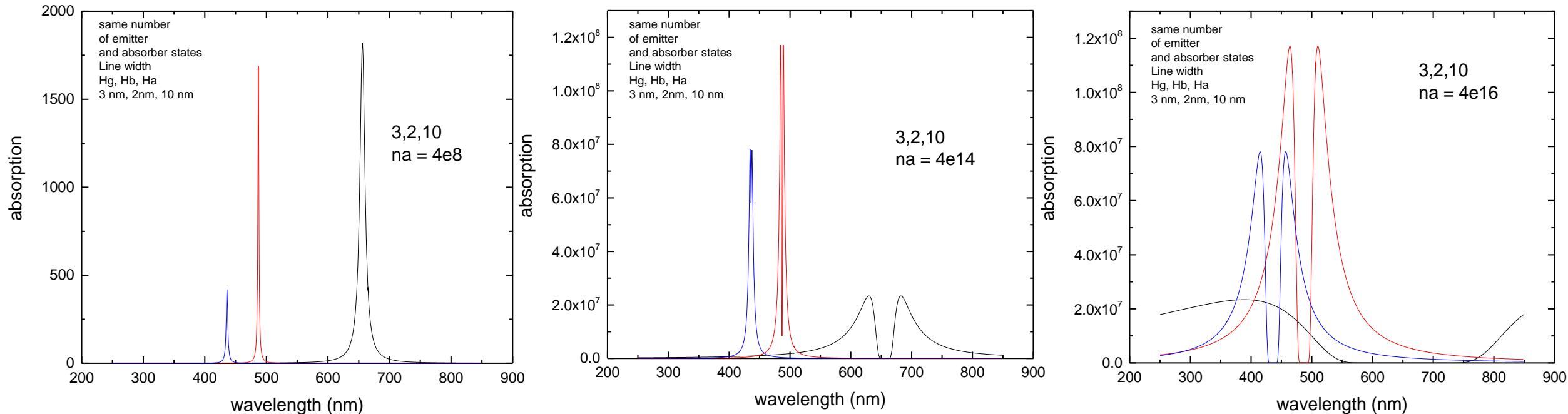


Line emission – Hydrogen Balmer Series

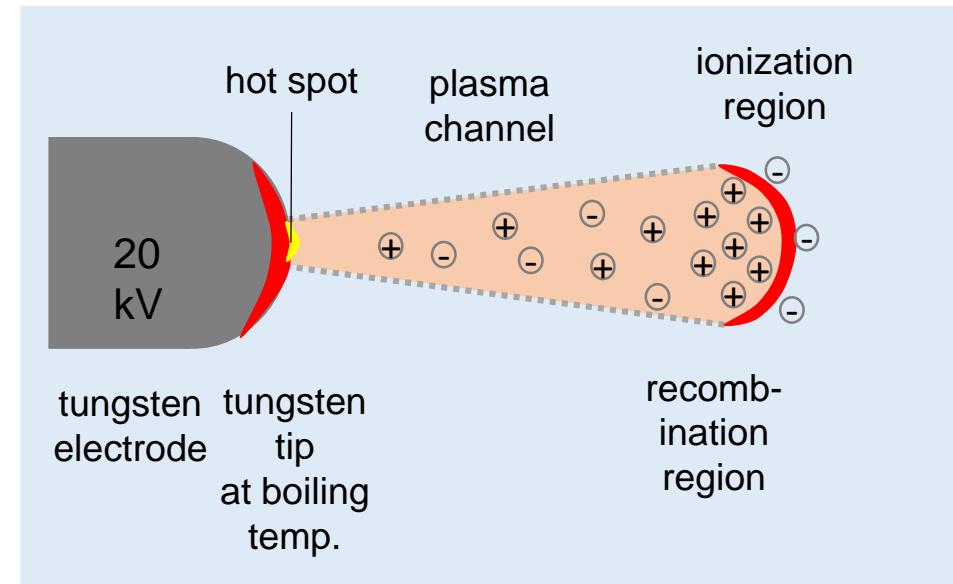
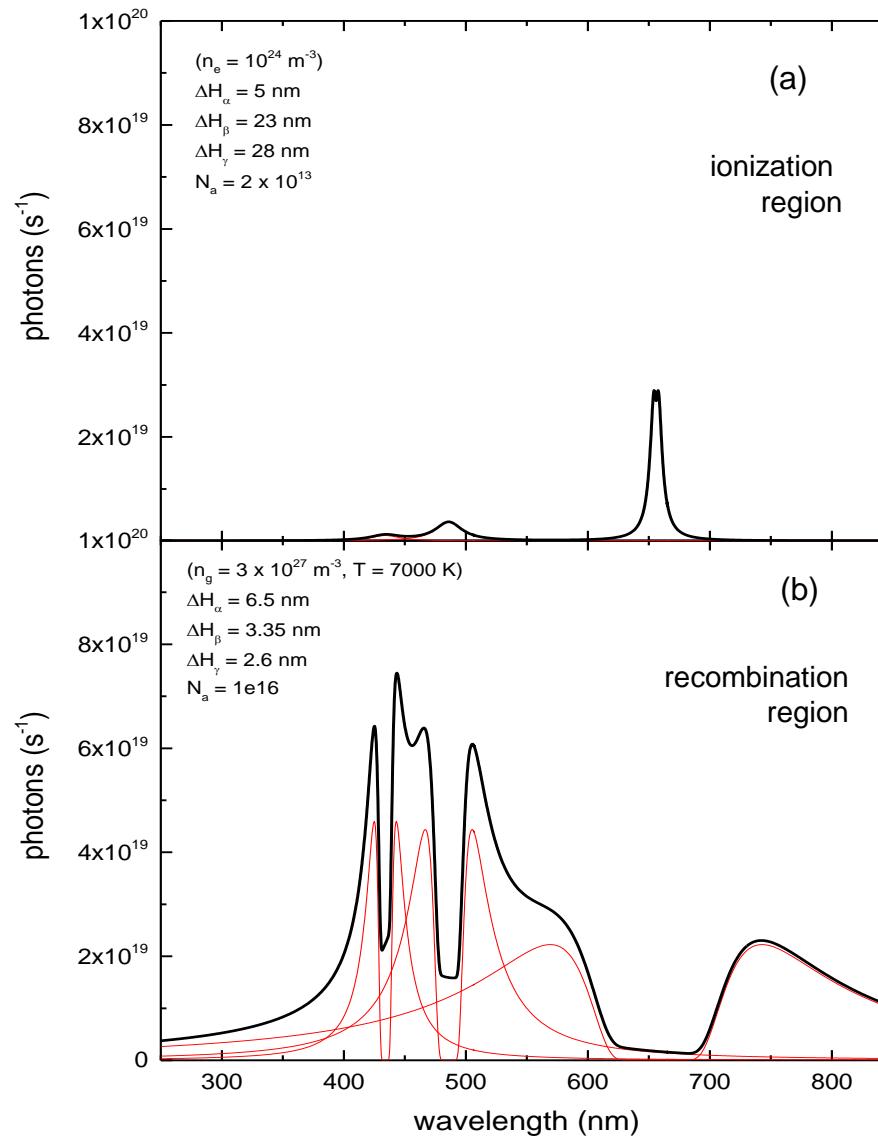


- H Balmer lines clearly visible
- Why is H γ more pronounced than H α ?

Line emission – Selfabsorption of Hydrogen Balmer Series



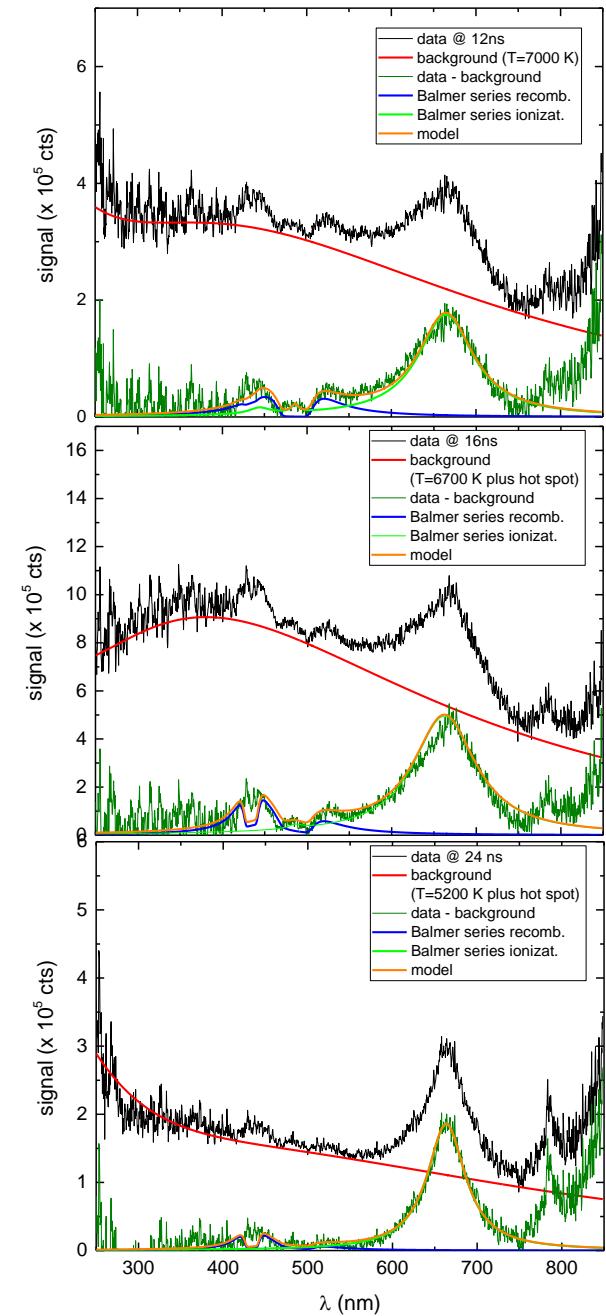
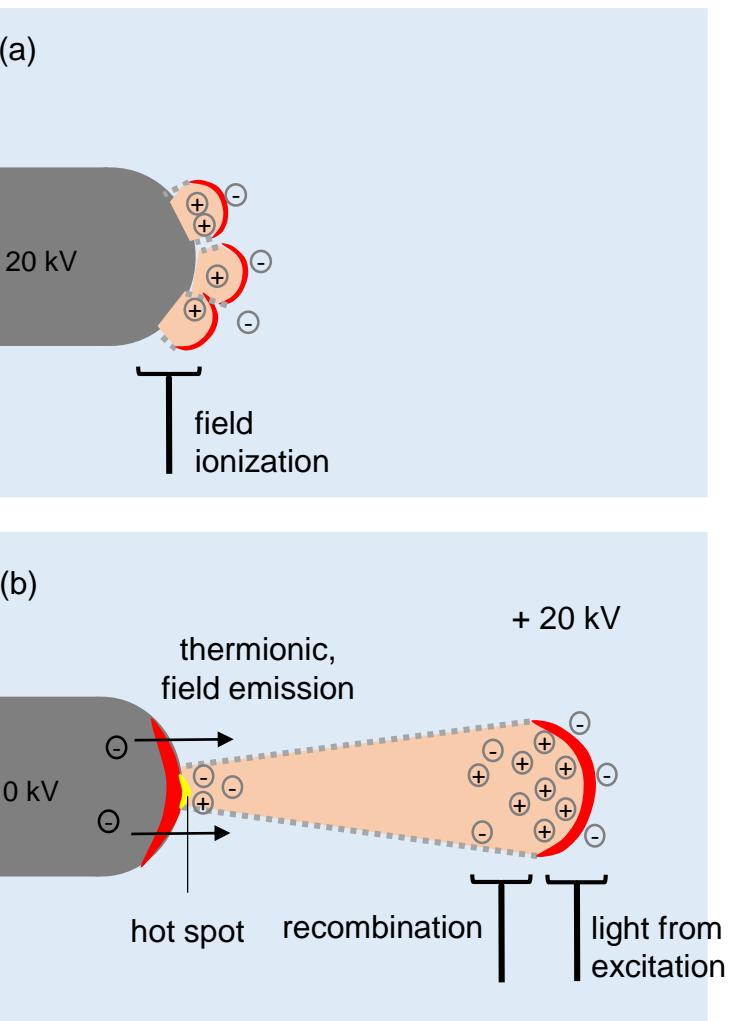
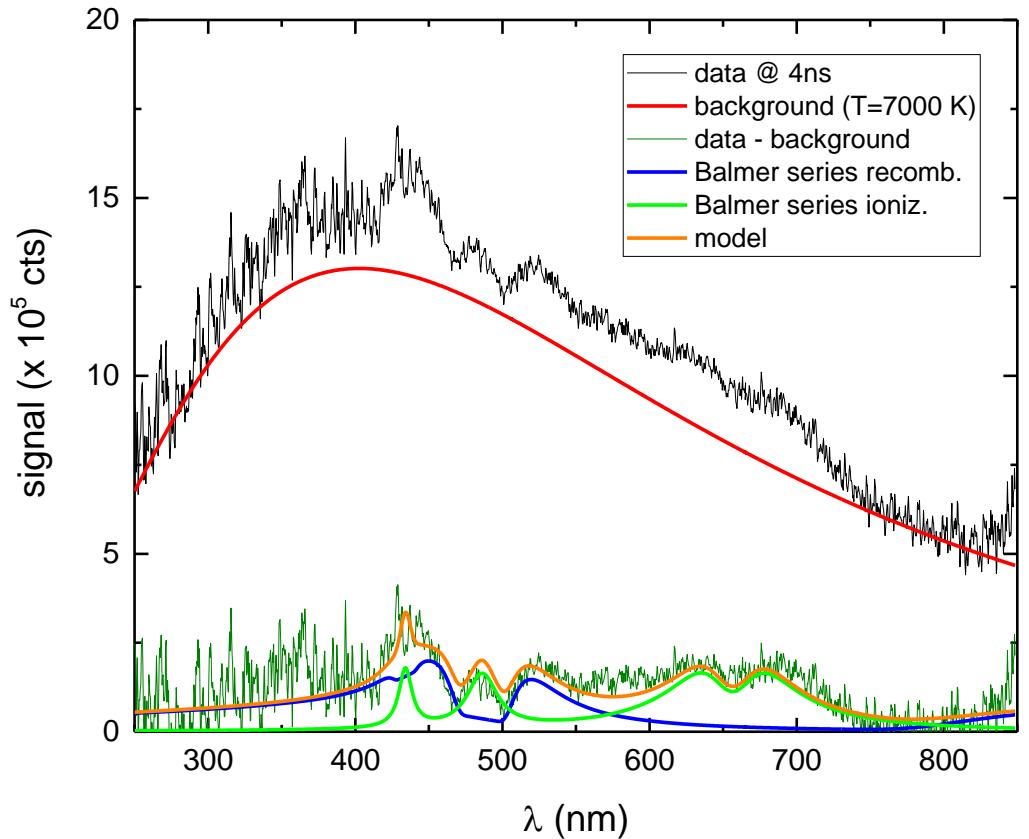
Nature of the $1/\lambda^x$ background in emission – cathode spots on the tungsten electrode



Three contributions

- Continuum blackbody $\sim 6000 - 7000 \text{ K}$
- H Balmer strong self absorption - recombination region (H γ dominates)
- H Balmer small self absorption - ionization region (H α dominates)

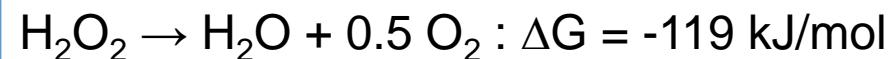
Line emission – Selfabsorption of Hydrogen Balmer Series



Efficiency to create H₂O₂ with these plasmas

Modeling the chemistry

- Take T₀ from cavitation
- use chemical equilibrium



0th order reaction (no local H₂O depletion)

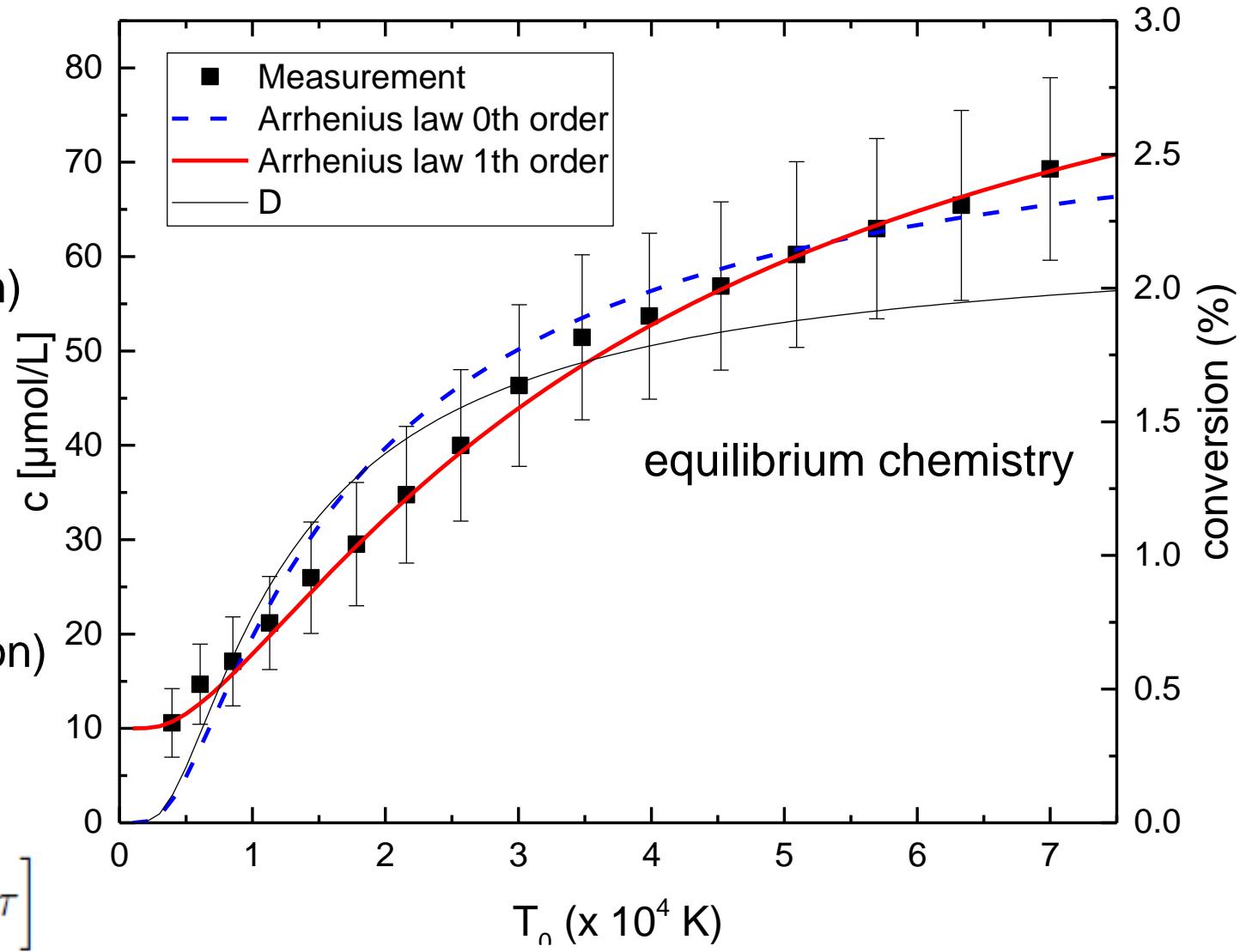
$$\frac{dn_{\text{H}_2\text{O}_2}}{dt} = kn_{\text{H}_2\text{O}}$$

$$n_{\text{H}_2\text{O}_2} = \exp\left(-\frac{E_a}{RT}\right) n_{\text{H}_2\text{O}} \tau + n_{\text{H}_2\text{O}_2}|_0$$

1st order reaction (with local H₂O depletion)

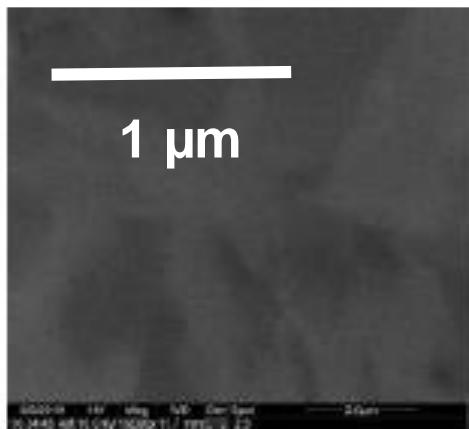
$$\frac{dn_{\text{H}_2\text{O}_2}}{dt} = -\frac{dn_{\text{H}_2\text{O}}}{dt} = kn_{\text{H}_2\text{O}}$$

$$n_{\text{H}_2\text{O}_2} = n_{\text{H}_2\text{O}_2}|_0 \exp\left[\exp\left(-\frac{E_a}{RT}\right) n_{\text{H}_2\text{O}} \tau\right]$$

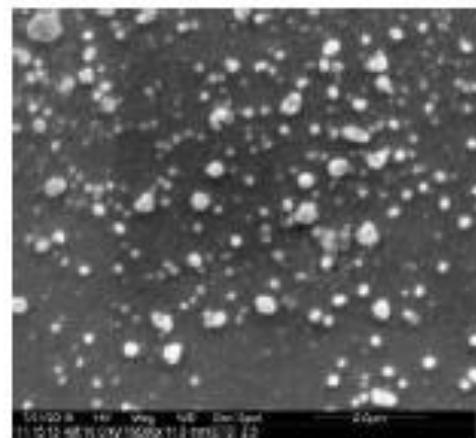


Nanosecond plasma based recovery of CuO nanocubes at copper electrodes

distilled water



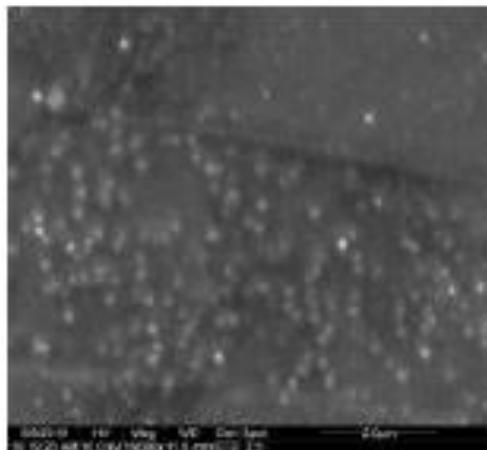
dist. Water + 0.0035M KCl



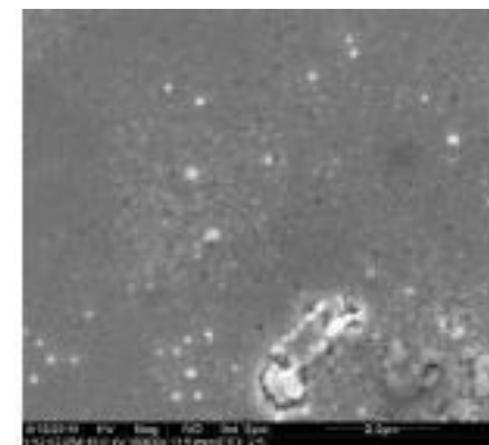
Creation of Cu-nanocubes possible, if

- no direct current to the sample, reduction of the oxide, plasma electrode distance to the sample important
- Competition between oxidation due to H_2O_2 , OH and reduction by H, e^-

distilled water



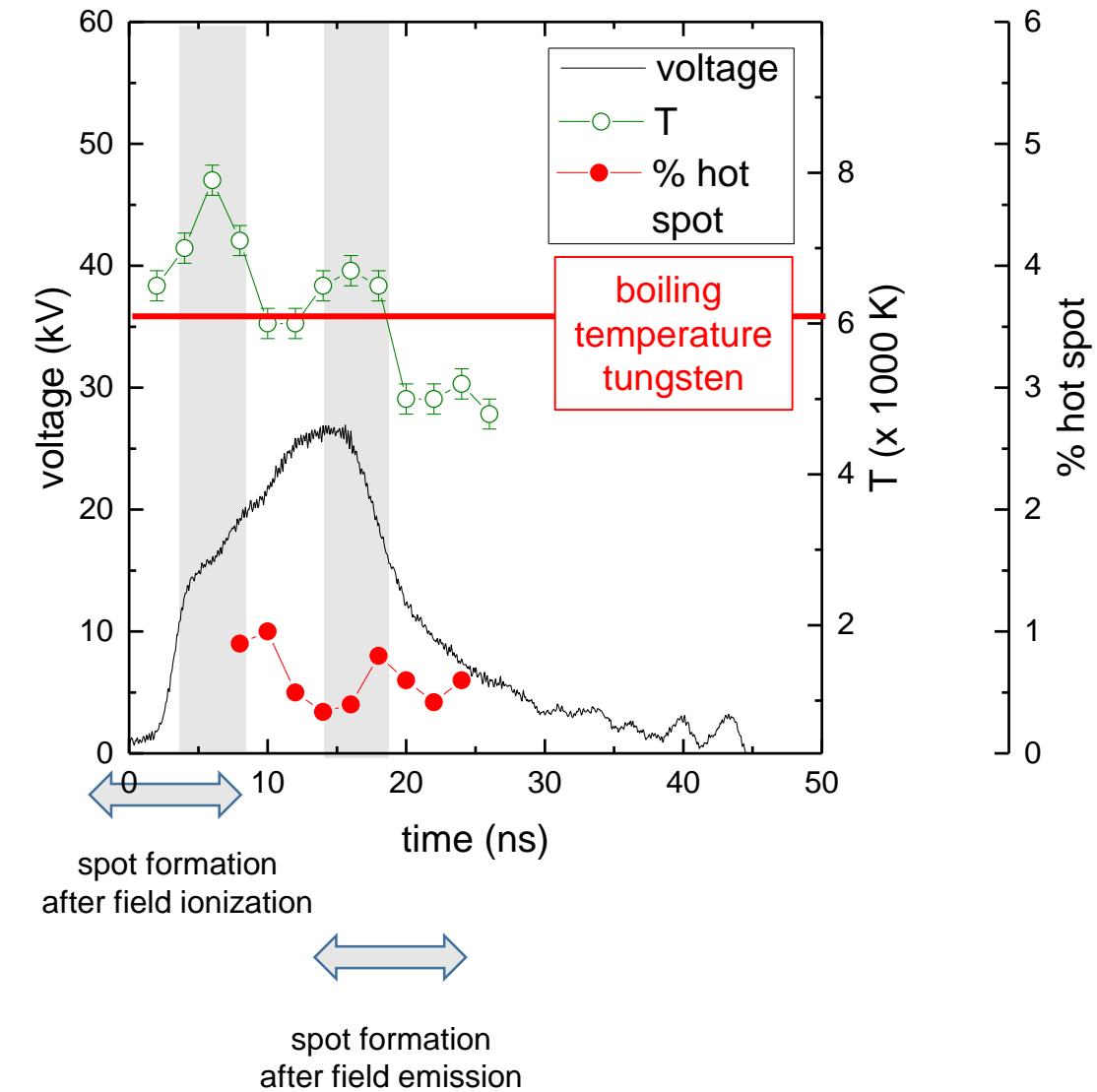
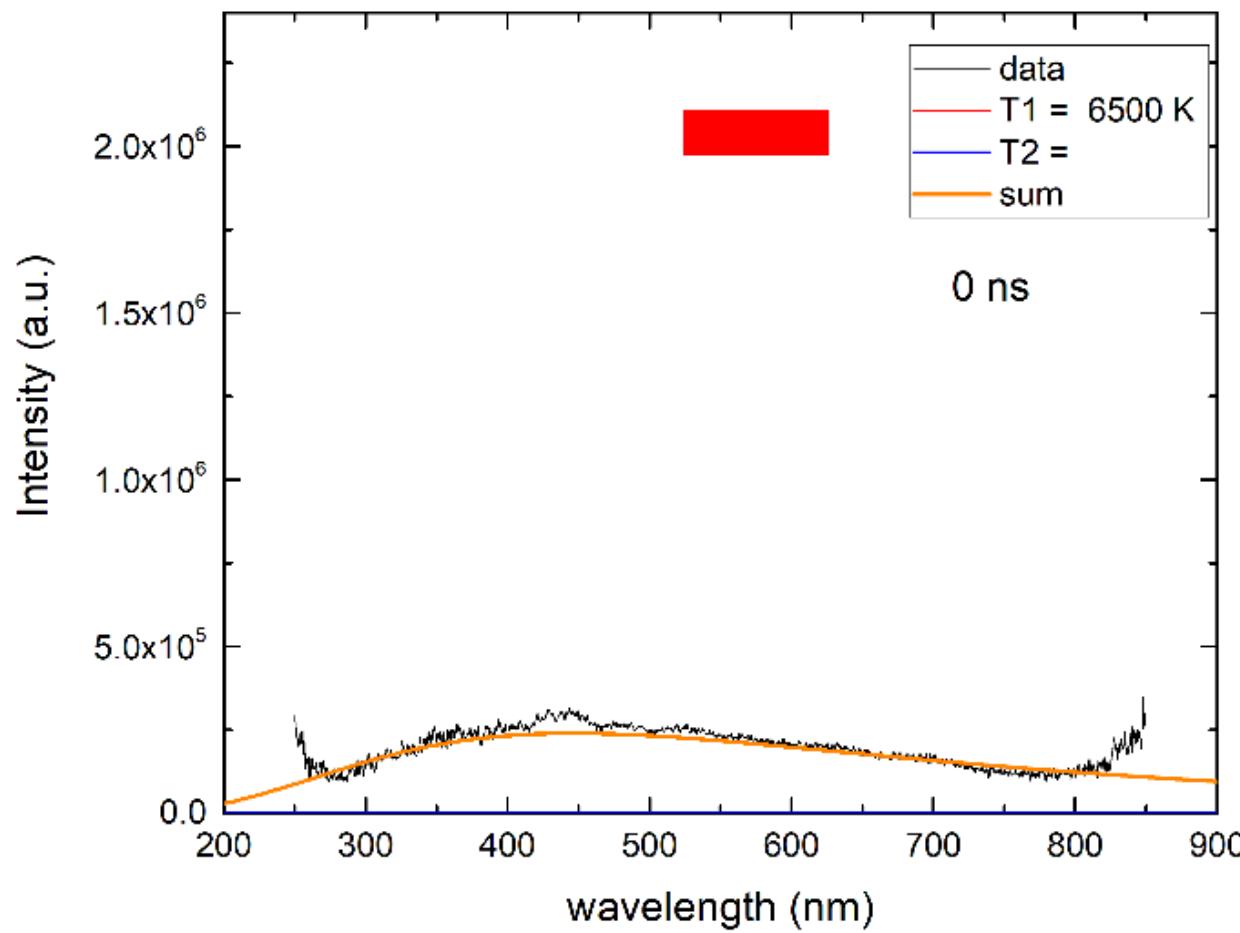
dist. Water + 0.0035M KCl



P. Grosse, B. Roldan et al.

Reference plasma based low pressure creation of CuO nanocubes

Nature of the $1/\lambda^x$ background in emission – cathode spots on the tungsten electrode



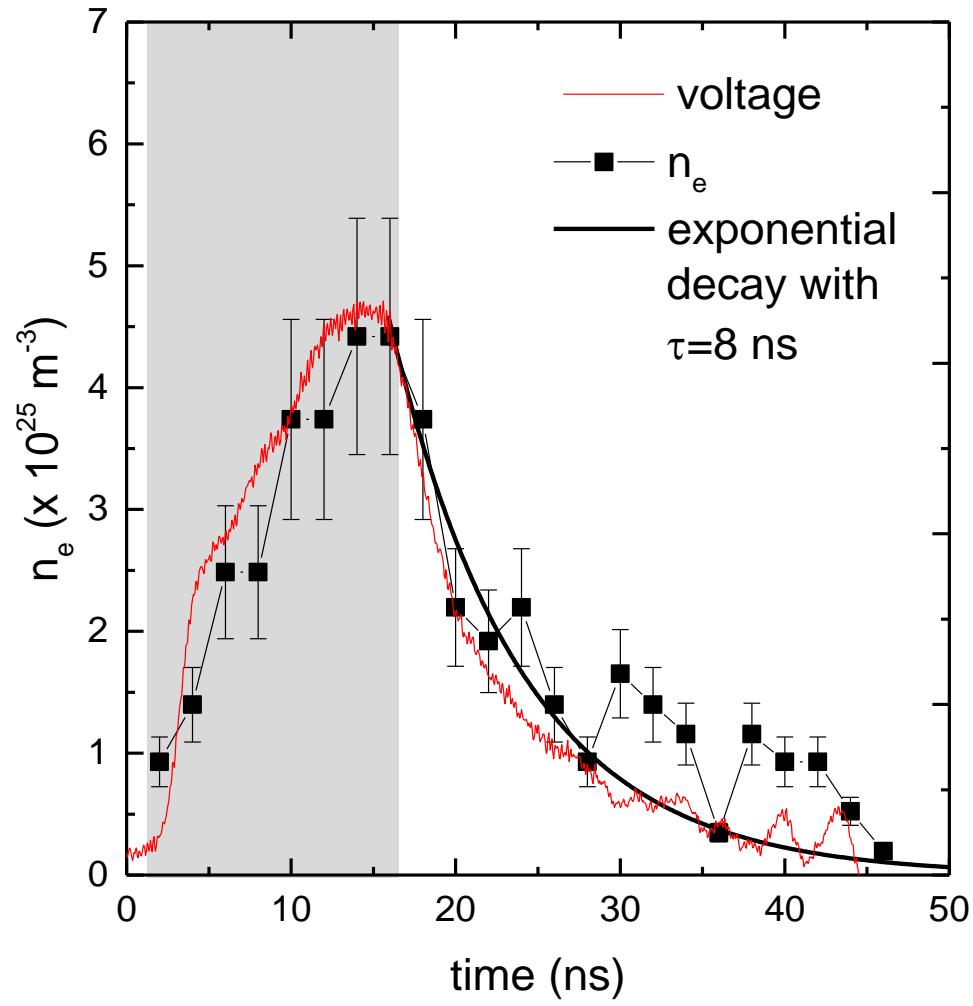
Line emission – Electron densities from Stark broadening of Hydrogen Balmer Series

electron density from Stark broadening

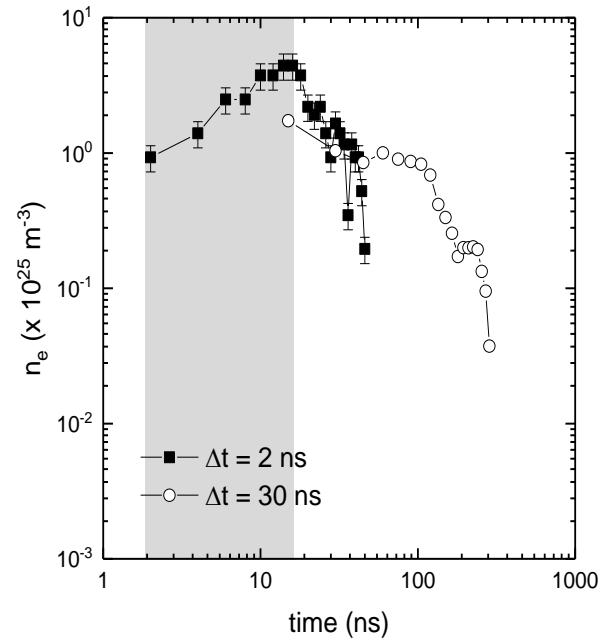
- n_e follows directly voltage
(no delay as in ns air plasmas)
ionization and recombination on ps
time scales at liquid densities)
- decay follows the voltage

$$n_e \sim 5 \times 10^{25} \text{ m}^{-3}$$

$$\text{ionization degree} \sim 10^{-3}$$



Comparison to n_e decay
in the re-ignited plasmas
During bubble expansion



Efficiency to create H₂O₂ with these plasmas

Measurement of the concentration with a colorimetric kit



Measurement:

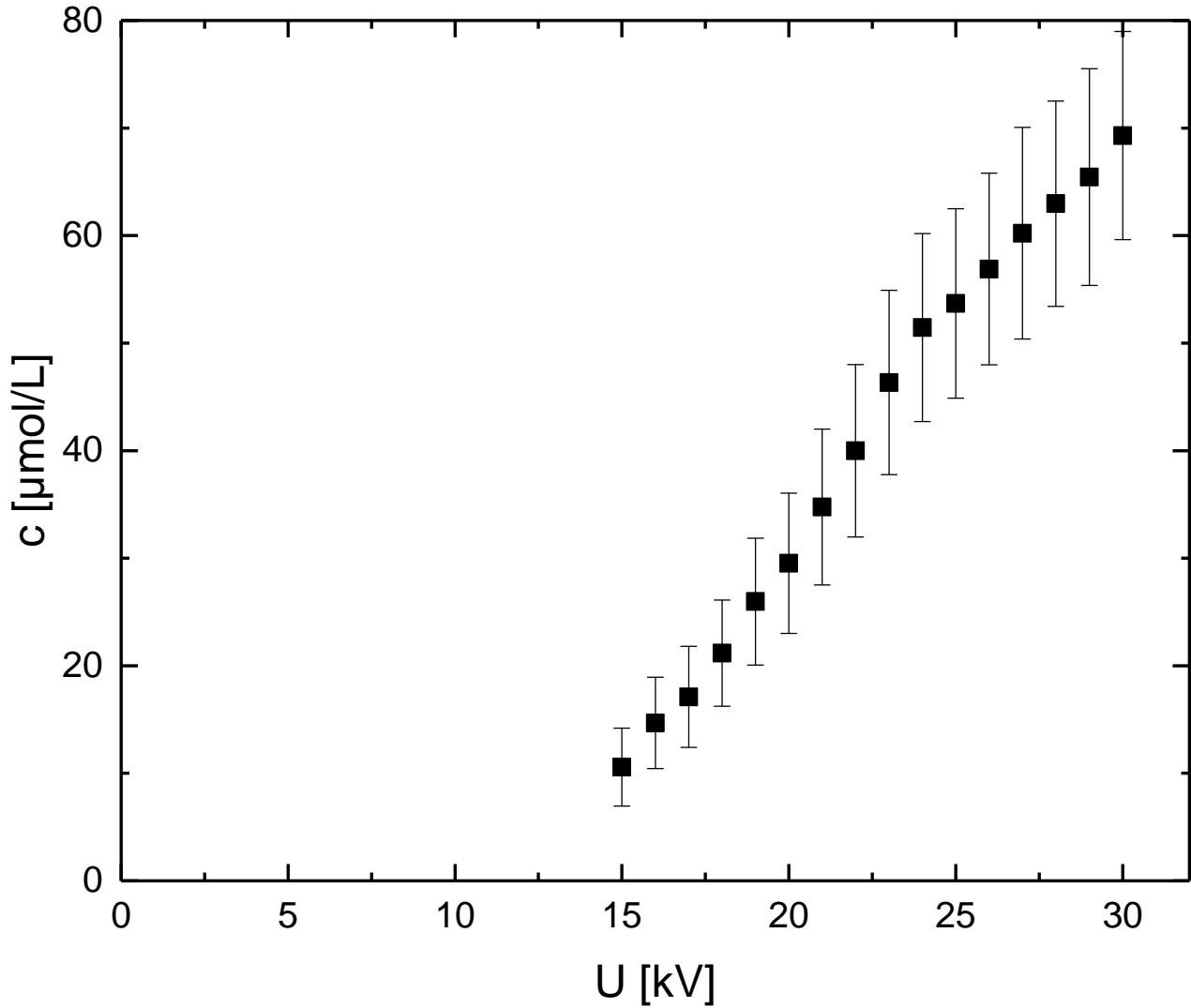
@ 20 kV, 30 µmol/l in 25 ml liquid is equivalent to 4×10^{17} species in total

Estimate:

Species in 25 µm radius spheres at density of liquid water $3 \times 10^{22} \text{ cm}^{-3}$: 2×10^{15} species

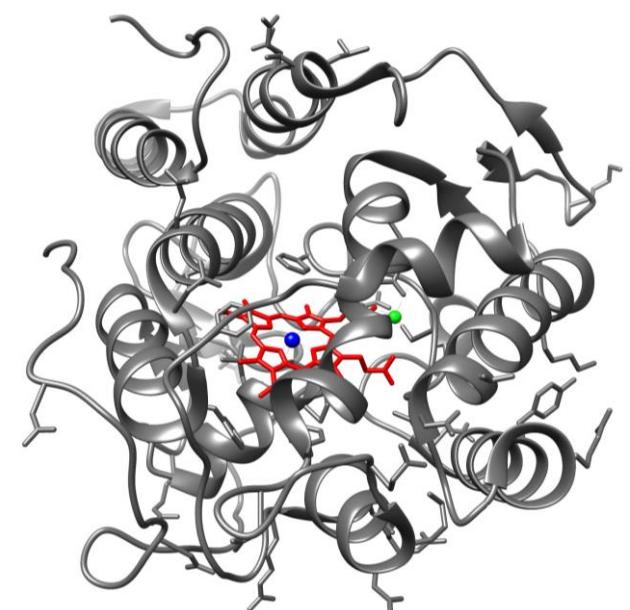
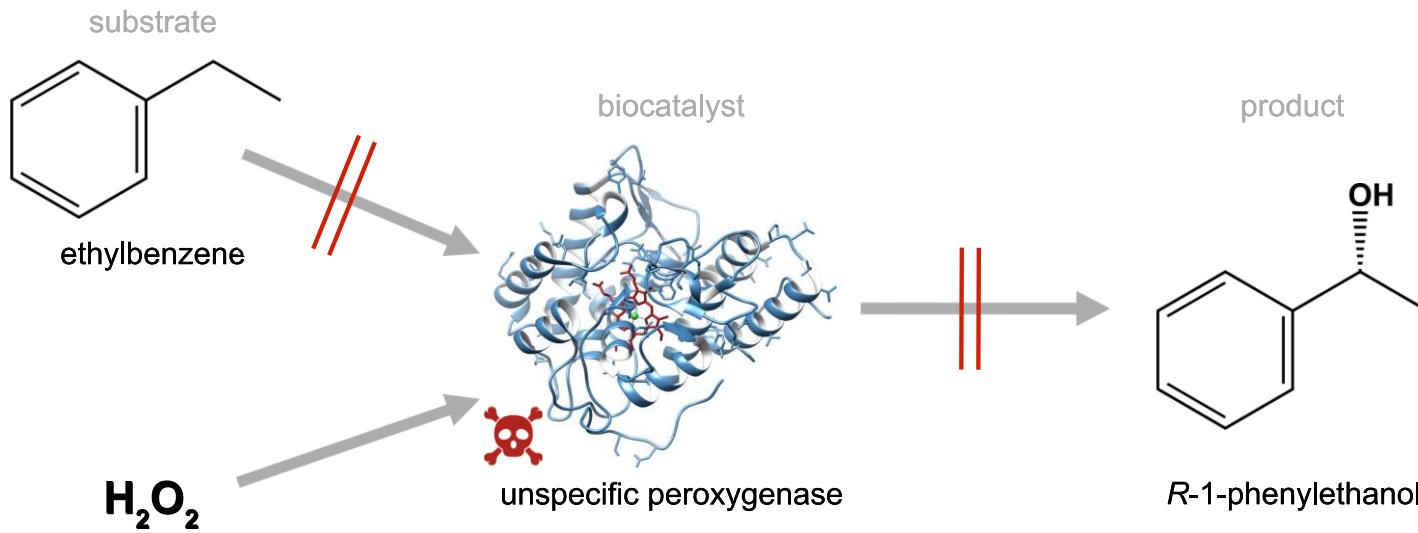
Total number of species treated in 10 min @ 15 Hz, 10 ns pulses: 3.8×10^{19} molecules are exposed to plasma

efficiency 1.1%



Biocatalysis

- use of H_2O_2 -dependent enzymes → oxidoreductases
- redox reaction → one substrate reduced, one oxidized
 - peroxyxygenases: $\text{H}_2\text{O}_2 + \text{AH} \rightarrow \text{A-OH} + \text{H}_2\text{O}$
- industrial application is limited → excess of H_2O_2 leads to suicide inactivation



unspecific peroxygenase of *Agrocybe aegerita* (AaeUPO) (PDB: 5OXU)